

GEORGIA INSTITUTE OF TECHNOLOGY  
OFFICE OF RESEARCH ADMINISTRATION

RESEARCH PROJECT INITIATION

Date: July 11, 1974

Project Title: Window and Curtain Wall Performance in High-rise Buildings

Project No: E-16-650/E-20-654

Principal Investigator S Dr. C. S. Martin (CE)  
Dr. J. I. Craig (AE)

Sponsor: National Science Foundation

Agreement Period: From 5/15/74 Until 10/31/75\* (Grant Expiration)

Type Agreement: Grant No. GK-43455X

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Reports Required:  
Annual Letter (if expected renewal is granted)  
Final Report

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\*Proposed project period (12 mos.) ends 5/14/75; all commitments to  
be met by grant expiration date unless formal extension is obtained in  
advance.  
Assigned to: Civil Engineering/Aerospace Engineering

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Date: 5/17/78

Project Title: "Window and Curtainwall Performance in Highrise Buildings"

Project No: E-16-650/E-20-654

Project Director: Dr. J. I. Craig/Dr. C. S. Martin

Sponsor: National Science Foundation

Effective Termination Date: 4/30/77

Clearance of Accounting Charges: 4/30/77

Grant/Contract Closeout Actions Remaining: NONE

- ☐ Final Invoice and Closing Documents
- ☐ Final Fiscal Report
- ☐ Final Report of Inventions
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- ☐ Other \_\_\_\_\_

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March 1975

Work on the subject grant to study the performance of windows and curtain walls in highrise buildings was initiated at the beginning of the summer academic quarter, 1974. The initial effort during the first year has been reduced somewhat from the level proposed. This was directly the result of, (a) unexpectedly increased teaching loads, and (b) delays in the completion of the particular highrise building to be studied. Work has proceeded along three lines and the progress to date in each area is summarized in the following sections.

Pressure Simulation Facility: The design of the full-scale wind pressure simulator has been completed. Detail and assembly drawings have been prepared, and final construction work is underway. Preliminary testing of the system is scheduled for the first week in May. The present facility has been extensively redesigned from that proposed in order to extend the performance. The basic design is of a positive displacement type where pressure loads are produced by controlled displacements of several large pistons forming part of the rear wall of a large test chamber (see attachment). The fixture will accept window test assemblies as large as 8 ft. by 15 ft. and is designed to produce pressure fluctuations of up to 50 psf, peak-to-peak, about mean pressures of up to 75 psf at cyclic rates as high as 20 Hz. The driving force is provided by small servo-controlled, 3000 psi hydraulic actuators connected to six large hinged pistons (doors) which cover the rear of the test fixture. The performance at extreme pressure and frequency is obtained by actually producing and controlling a resonance behavior of the pistons and test chamber. The apparatus will operate in a closed-loop fashion using a modified cascade control system and will accept both random and sinusoidal forcing functions.



The system will be initially operated with only two of the six actuators; however, provision has been made and hydraulic capacity provided to handle the remaining four actuators. Four 55 in. by 99 in. by  $\frac{1}{4}$  in. doubly glazed lights along with appropriate sections from the curtain wall used in the building have been acquired and are ready to be installed in the test fixture. The design and preliminary results for this facility will be presented in a Research Summary Report at the 2nd U.S. National Conference on Wind Engineering Research, Ft. Collins, June 1975.

Field Measurement Program: The field measurement study in the 30 story "Tower Place" building has been held up due to delays in the building construction program. It is anticipated that the curtain wall installation will be completed by the end of April so that instrumentation can be moved into the building during late spring. Work to date has centered in the laboratory on the design and construction of the response measurement instruments, particularly on the development of a remote sensing instrument to simultaneously from one side measure the movement of both lights in a doubly-glazed window. This latter instrument has been constructed and is presently being calibrated in the laboratory prior to its field use. The basic design uses a narrow collimated light beam from a small He-Ne laser reflected off of the inner and outer lights. Two solid state, linear position-sensitive photo-detectors are used to differentially analyze the movement of the reflected beam, and a simple analog computational circuit is used to derive the normal and angular displacement components. Details of the design and operation of this instrument will be presented shortly in a paper by the graduate research assistant participating in the study.

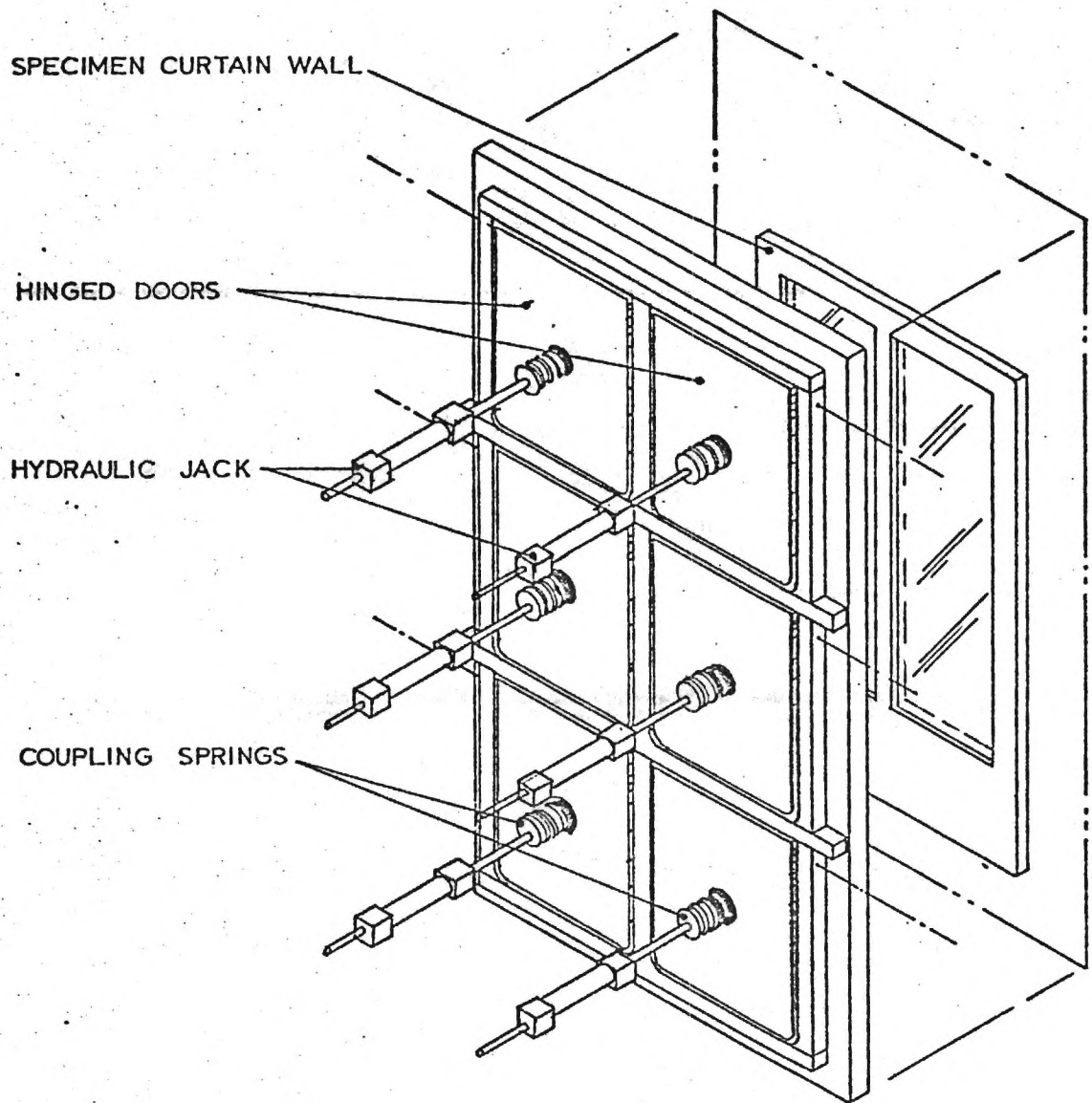
In a separate program funded by the building owners, a wind tunnel study of the building wind pressure loading was run in the Aerospace Engineering 9 ft. tunnel. All data from these tests have been recorded in computer-compatible



form and will be used for comparison with the full scale response measurements.

Analytical Studies: Analytical studies are being directed at modelling the dynamic response of both single and double-glazed lights. In the first case, the effort is being directed towards properly representing the damping and non-linearly elastic behavior in the boundaries (gaskets) by means of a hybrid finite element approach. In the second case, the boundary-coupled response of the two lights in a double-glazed assembly is being modelled and the effect of venting of the gas entrapped between the two plates is being studied.

Dr. L. Z. Emkin has been unable to participate as an associate on the program. His responsibilities, particularly in the analytical modelling, have been assumed as of January 1975 by Dr. Barry Goodno, recently appointed as Assistant Professor of Civil Engineering. A copy of his biosketch is attached.



CURTAIN WALL TEST FACILITY

## BIOGRAPHICAL SKETCH

Barry John Goodno

### Present Position and Address:

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School of Civil Engineering  
Georgia Institute of Technology  
Atlanta, Georgia 30332

Date and Place of Birth: July 19, 1947, Madison, Wisconsin

### Education:

Bachelor of Science, Civil Engineering, January, 1970  
University of Wisconsin, Madison, Wisconsin  
Master of Science, Structural Engineering, June, 1971  
Stanford University, Stanford, California  
Doctor of Philosophy, Civil Engineering, January, 1975  
Stanford University, Stanford, California

### Professional Society Membership:

American Society of Civil Engineers, (ASCE), Associate Member  
Earthquake Engineering Research Institute, (EERI), Member  
Seismological Society of America, (SSA), Member  
The Society of the Sigma Xi, Member  
Engineer-in-Training, (EIT), State of Wisconsin, (June, 1969)

### Honor Societies:

Phi Eta Sigma, Chi Epsilon, Tau Beta Pi, Phi Kappa Phi

### Employment History:

Assistant Professor, School of Civil Engineering  
Georgia Institute of Technology  
Atlanta, Georgia  
September 1974 to Present



Senior Structural Research Engineer  
John A. Blume and Associates, Engineers  
San Francisco, California  
April to July, 1973

National Science Foundation Teaching  
Fellow, Stanford University  
Spring Quarter 1972 and Winter Quarter 1973

Structural Engineer  
Arnold and O'Sheridan, Inc.  
Structural Engineering Consultants  
Madison, Wisconsin  
Summer 1969, 1970

Student Engineer Trainee  
Wisconsin Division of Highways  
Madison, Wisconsin  
Summer 1966, 1967, 1968

Publications:

Goodno, B. J., and Kost, E. G., Structural Design for Ground Motions: Reduction of Input Energy and Structural Response - Literature Survey, John A. Blume and Associates, Engineers, San Francisco, California, September 1973.

Goodno, B. J., Dynamic Analysis of Suspended-Floor Highrise Buildings Using Super-Elements, thesis presented to the Department of Civil Engineering, Stanford University, Stanford, California, in September, 1974 in partial fulfillment of the requirements for the degree of Doctor of Philosophy.

Current Fields of Interest:

Damping in Highrise Building Structures

Finite - Element Analysis of Shear Walls and Shear Cores in Highrise Buildings

Dynamic Analysis of Floor Systems to Reduce Human Discomfort in Building Structures

Performance of Windows and Cladding in Highrise Building Structures

E 76-650

Window and Curtainwall Performance  
in Highrise Buildings

Final Technical Report  
Submitted to the  
National Science Foundation  
by  
Georgia Institute of Technology  
School of Aerospace Engineering  
School of Civil Engineering

C. S. Martin, Co-principal Investigator  
J. I. Craig, Co-principal Investigator  
B. J. Goodno, Associate Investigator  
R. B. Deo, Post Doctoral Fellow

Grant ENG 73-04216  
Starting Date: May 1974  
Completion Date: March 1978

## Scientific Description of Research and Results

A detailed and comprehensive description of the research and the results obtained from this work are provided in the attached copy of the Final Report. This report covers the field study conducted at a highrise building in Atlanta, the analytical modelling of typical sheer glass cladding (including a full listing of all computer code developed), and the design, construction and use of a full scale laboratory window pressure test facility. The following paragraph provides a brief abstract of the work.

The present study is directed at more precisely characterizing the nature of the local cladding pressure loads in highrise buildings and more accurately modelling the structural response of the cladding to these loads. A combination of full scale field measurements of loads and response, detailed discrete element analytical modelling of the cladding itself, and controlled laboratory simulation and testing have been carried out. A comprehensive measurement methodology is developed and used in the field study to characterize the temporal and spectral features of the differential pressure loading and to assess the causality of these loads for the measured dynamic response of the cladding - in particular large insulating glass windows. The analytical work is centered around the development of a large and detailed discrete element model for the cladding, including its connection to the underlying structure. When the model parameters are adjusted or "calibrated" to laboratory measured values, it is shown that the predicted cladding performance agrees well with field and lab response measurements. Out-of-phase double panel glass plate modes are predicted and are measured. Finally, the design and use of a full-scale laboratory facility for simulating both the static and dynamic features of the cladding pressure loads are described. This facility is used for a detailed modal test of a large double panel window under typical mounting conditions and the results are compared to the present analytical predictions.

## List of Publications

1. Deo, R. "An Experimental Investigation of the Performance of Insulating Windows-Wind Loads and Dynamic Response," Ph.D. dissertation, Georgia Tech, June 1977.
2. Wey, J.D., "Properties and Behavior of Window-Frame Assemblies," M.S. thesis, Georgia Tech, September 1976.



3. Craig, J. I., and Palfery, J. G., "Full-Scale Wind Pressure Simulation for Windows and Curtain Wall Structures," ASCE-EMD Specialty Conference on Dynamic Response of Structures: Instrumentation, Testing Methods and System Identification, UCLA, Los Angeles, March 1976.
4. Craig, J. I., Palfery, J.G. and Martin, C. S. "Wind Pressure Simulation and Response Measurement for Windows," proceedings of Second U.S. National Conference on Wind Engineering Research, Colorado State University, Ft. Collins, Colorado, June 1975.
5. Deo, R. and Craig, J. I. "An Optical Technique for the Measurement of Time Dependent Displacements," Advances in Test Measurement, Vol. 15, Instrument Society of America, 1978, pp. 339342.
6. Deo, R. B. and Craig, J. I., "Wind Load Causality for Dynamic Response of Windows," proceedings of the 3rd U.S. National Conference on Wind Engineering Research, University of Florida, Feb. 1978, III - 12-1.
7. Deo, R. B. and Craig, J. I., "Full Scale Measurement of Cladding Pressure," proceedings of the 3rd U.S. National Conference on Wind Engineering Research, University of Florida, Feb. 1978, III -11-1.
8. Goodno, B. J., "Response of Glass Cladding in Highrise Buildings," proceedings of the 3rd U.S. National Conference on Wind Engineering Research, University of Florida, Feb. 1978, p. III-10-1.
9. Goodno, B. J., "Structural Performance of Insulating Windows," submitted for publication in the Journal of the Structural Division, ASCE.
10. Goodno, B. J., "Dynamic Response of Glass Curtain Wall Elements," submitted for publication in the Journal of the Structural Division, ASCE.
11. Craig, J. I., Goodno, B. J. and Deo, R. B., "Window and Curtain Wall Performance in Highrise Buildings," Final Report, NSF Grant ENG 73-04216, Georgia Institute of Technology, GITAER 78-100, SCEGIT 78-170, March 1978.

## Theses

One Masters and a Doctoral thesis were completed in support of and with funds provided by the grant. The titles are as follows:

Deo, R. B., "An Experimental Investigation of the Performance of Insulating Windows - - Wind Loads and Dynamic Response," Ph.D. Dissertation, Georgia Institute of Technology, School of Aerospace Engineering, June 1977.

Wey, J. D., "Properties and Behavior of Window - Frame Assemblies," M.S. Thesis, Georgia Institute of Technology, School of Civil Engineering, Sept. 1976.

## Inventions

A laser displacement meter, described in Reference #5, was developed to measure the movement of the outer pane of a doubly glazed window from a reference position inside the building. This instrument was not judged to be patentable but the present design is being studied and is now under consideration by others for use in various experimental studies.

## Scientific Collaborators at Georgia Institute of Technology

J. I. Craig, Associate Professor, Co-principal Investigator

B. J. Goodno, Assistant Professor, Associate Investigator

C. S. Martin, Professor, Co-principal Investigator

R. B. Deo, Post Doctoral Fellow

J. D. Wey, Graduate Research Assistant

J. G. Palfery, Research Engineer

## Comments

The reported work has resulted in the development of a laboratory facility for evaluating the performance of full scale window assemblies under both static and dynamic loads. The work has also promoted the development of a first-class experimental research capability in this area of structural dynamics at Georgia

Tech. Finally, these efforts have led, in part, to the preparation of a new 3 year research proposal to NSF/PFRA which has now been funded. The new grant, ENV 77-04269 is titled, "Influence of Nonstructural Cladding on Dynamic Properties and Response of Highrise Buildings," with co-principal investigators, B. J. Goodno, J. I. Craig, K. M. Will.

Signatures

\_\_\_\_\_  
C. S. Martin  
Co-principal Investigator

\_\_\_\_\_  
J. I. Craig  
Co-principal Investigator

\_\_\_\_\_  
B. J. Goodno  
Associate Investigator

Date: 5/17/78



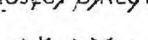
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1. INSTITUTION AND ADDRESS Georgia Institute of Technology Atlanta, Georgia		2. NSF PROGRAM Water Resources, Urban & Environmental Engr.	3. PRINCIPAL INVESTIGATOR(S) Dr. J. I. Craig Dr. C. S. Martin
4. AWARD NUMBER ENG73-04216 A01	5. DURATION (mos) 36	6. AWARD PERIOD from 5/15/74 to 4/30/77	7. AWARDEE ACCOUNT NUMBER E-16-650/E-20-654

8. PROJECT TITLE  
"Window and Curtain Wall Performance in Highrise Buildings."

9. SUMMARY (ATTACH LIST OF PUBLICATIONS TO FORM)

The present study is directed at more precisely characterizing the nature of the local cladding pressure loads in highrise buildings and more accurately modelling the structural response of the cladding to these loads. A combination of full scale field measurements of loads and response, detailed discrete element analytical modelling of the cladding itself, and controlled laboratory simulation and testing have been carried out. A comprehensive measurement methodology is developed and used in the field study to characterize the temporal and spectral features of the differential pressure loading and to assess the causality of these loads for the measured dynamic response of the cladding - in particular large insulating glass windows. The analytical work is centered around the development of a large and detailed discrete element model for the cladding, including its connection to the underlying structure. When the model parameters are adjusted or "calibrated" to laboratory measured values, it is shown that the predicted cladding performance agrees well with field and lab response measurements. Out-of-phase double panel glass plate modes are predicted and are measured. Finally, the design and use of a full-scale laboratory facility for simulating both the static and dynamic features of the cladding pressure loads are described. This facility is used for a detailed modal test of a large double panel window under typical mounting conditions and the results are compared to the present analytical predictions.

9. SIGNATURE OF PRINCIPAL INVESTIGATOR/ PROJECT DIRECTOR 	TYPED OR PRINTED NAME Dr. J. I. Craig	DATE 5/17/78
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**FINAL REPORT**

**WINDOW AND CURTAIN WALL PERFORMANCE  
IN HIGH-RISE BUILDINGS**

**By**

**J. I. Craig**

**B. J. Goodno**

**R. B. Deo**

**Prepared for**

**NATIONAL SCIENCE FOUNDATION**

**DIVISION OF ENGINEERING**

**GRANT ENG-73-04216**

**March 1978**

**GEORGIA INSTITUTE OF TECHNOLOGY**

**SCHOOL OF AEROSPACE ENGINEERING**

**SCHOOL OF CIVIL ENGINEERING**

**ATLANTA, GEORGIA 30332**

1978



WINDOW AND CURTAIN WALL PERFORMANCE  
IN HIGH-RISE BUILDINGS

J. I. Craig

B. J. Goodno

R. B. Deo

FINAL REPORT

Grant ENG-73-04216

Prepared for the

National Science Foundation

Division of Engineering

by

GEORGIA INSTITUTE OF TECHNOLOGY  
School of Aerospace Engineering  
School of Civil Engineering

March 1978



## ABSTRACT

The present study is directed at more precisely characterizing the nature of the local cladding pressure loads in highrise buildings and more accurately modelling the structural response of the cladding to these loads. A combination of full scale field measurements of loads and response, detailed discrete element analytical modelling of the cladding itself, and controlled laboratory simulation and testing have been carried out. A comprehensive measurement methodology is developed and used in the field study to characterize the temporal and spectral features of the differential pressure loading and to assess the causality of these loads for the measured dynamic response of the cladding-in particular large insulating windows. The analytical work is centered around the development of a large and detailed discrete element model for the cladding, including its connection to the underlying structure. When the model parameters are adjusted or "calibrated" to laboratory measured values, it is shown that the predicted cladding performance agrees well with field and lab measurements. Out-of-phase double panel glass plate modes are predicted and are measured. Finally, the report describes the design and use of a full scale laboratory facility for simulating both the static and dynamic features of the cladding pressure loads. This facility is used for a detailed modal test of a large double panel window under typical mounting conditions and the results are compared to the present analytical predictions.

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## 1.0 INTRODUCTION

### 1.1 Background

The past decade has witnessed considerable activity and progress in the analysis of and design for wind loads on tall buildings. Two important considerations, the effect of terrain roughness and the dynamic aspects of the wind loading itself, are now better understood because of a substantial research effort. Much has been accomplished in expanding our understanding of wind loads on the total structure, however our knowledge, relatively speaking, of the loading and response of windows, curtain walls and architectural details of the building cladding is yet in quite an infant state. To illustrate this it should be noted that design wind loads are usually given for the overall structure but not specifically for the glass panes which are incorporated in the cladding. The fact that such geometric features as corners have been shown to produce substantial changes in local flow conditions dictates that the design wind load for windows and cladding be different than that for the structure as a whole. Furthermore, from a dynamic standpoint the gust loading for a window would not be expected to be the same as that for the structure. In a first step, the American National Standards Institute code ANSI A58.1-1972 has recognized the problem and provided for increased design wind pressures in the vicinity of corners or setbacks.

At the same time, increased activity in building design, especially with regard to tall buildings of unusual design, is rapidly extending the engineering state-of-the-art. Unfortunately, engineering knowledge and design methodology in several areas appear to have been surpassed by

architectural innovation. This problem has been particularly evident in the design of building cladding systems which in recent years have become increasingly varied. Cladding in general is considered non-structural, but yet must form a first line of defense for the structure against environmental forces arising from a variety of sources: pressure loads due to wind or sonic boom, impact due to debris, solar thermal loads, seismic loading, and interactive structural loading due to building frame distortions. The complexity and interaction of these factors is shown schematically in Fig. 1.1-1 where a distinction has been made as to source and manner (static, dynamic) of loadings.

In traditional engineering design greater analytical attention has been given to major building components (foundation, frame) than to architectural details such as the cladding or window systems. The principal reason is, of course, that the failure of a major component would be catastrophic. However, the increasingly frequent failure of windows and cladding to perform as intended coupled with the fact that the cladding can now account for 10 - 20% of the initial building cost, has noticeably increased the economic risk of advanced design concepts. A detailed discussion of recent failures would be overly lengthy, but the instances documented by Refs. 1 - 10 clearly underline the fact that a reliable design methodology is lacking. The pattern of failures strongly suggests that wind forces constitute the major consideration for the design of cladding systems. The problem is compounded, however, by the fact that failures have been variously attributed to a number of factors including thermal stressing, inplane loading due to rocking and glazing interference, improper use of cleaning fluid, projectiles,



as well as dynamic effects such as glazing pull-out or interaction between two panes in a double-glazed window.

Clearly, many of the failures of windows and cladding assemblies, reported or not, can be generally attributed to improper design. The solution to these problems is not always so direct, however. While all important factors may not have been considered, it is also possible that some factors were previously unknown. The situation can only be made more acute by the healthy diversification of design and construction present in modern highrise buildings. The use of new materials and advanced construction techniques makes it imperative not only that all significant loads are defined but also that the various response patterns are properly identified and understood.

## 1.2 Design Considerations

Wind forces are recognized to constitute a major loading in the structural design of tall buildings and there is currently much interest and activity in refining knowledge of these loads. While treatment of the effect of wind loading on the overall structural performance has received wide attention, the localized effects of wind pressure, especially strongly fluctuating pressure, on build cladding has until recently received much less consideration.

In developing rational design procedures for cladding components and systems, much of the work to date has centered on the specification of static wind pressure loads. These figures are often augmented by factors reflecting the expected magnitudes of wind gusts. In some instances account is also taken of the pressure distribution caused by

wind flow around corners and other gross shape features of the structure (11).

Results from extensive testing programs by the two major glass manufacturers in the U. S. form the basis of most design procedures and codes. This data has been obtained exclusively from static testing although charts are available for various equivalent static loads (12). Extensive discussion of these procedures and codes, as well as consideration of wind load prediction and glass strength, formed a major portion of the national conference on building research in 1967 (13 - 15). Wind effects and wind loads on buildings and other structures have been the subject of much more work and several national and international conferences during the past decade have been largely devoted to this topic (16 - 21). While a majority of the reported work has dealt with overall or integrated wind loading, an increasing concern for the nature of localized loading is evident. A comprehensive treatment of these combined factors in current building design is beginning to emerge, and innovations such as buildings designed to minimize, not just withstand, wind loads (22), or specification of windows of a variable size with height for a tall building (23) have been reported.

A rational methodology for designing building cladding and windows in particular has lagged considerably behind, and only recently have attempts been made to characterize the nature of localized loadings (24 - 26). There have been only scattered attempts to ascertain the resultant localized structural response to these forces (27), although at least one experimental facility has been built to simulate the anticipated conditions (28). As far as response characteristics or

failure modes are concerned there is much conjecture but little hard evidence is available. It is observed that pressure loads may in certain circumstances cause window panes to actually be ejected from their glazing but it is often assumed that the force responsible for this is a quasi-static pressure load due largely to wind suction effects. The situation is often augmented by weakness in the mullion or glazing design or by improper installation of the light itself. Instances in which the light does not appear to fracture before leaving the glazing are often recorded and consequently it can be reasonably assumed that glass fracture by itself is not a major consideration in this connection. There has been, however, no real attention directed toward possible dynamic mechanisms which might play a role in this ejection type failure of lights. In fact, the only available reference to dynamic effects in commercial curtain wall designs that we have noted regards the possible psychological effects on building occupants of low frequency (3 - 7 Hz) panel vibrations (29, 30). As human response and discomfort is receiving more and more attention (31), psychological reaction to excessive glass movement must be considered. Regarding the range of natural frequencies of large windows, Pretlove and Bowles (32) have, from a large sampling of window response in Britain, found that a typical large window (88.3 in by 125.8 in) will have a natural frequency of 6 Hz. Although the deflection itself may not be structurally critical, excessive deflections can be of concern to occupants.

An interesting failure mode involving excessive deflection coupled with dynamic response has been suggested in windows that are doubly

glazed. The proposed failure mode, particularly for those windows incorrectly or incompletely installed involves a sharp transient vibration of the panes in which each moves out of phase with the other. If sufficient amplitude is developed, the panes might actually contact with the result that either or both may fail. While the ultimate failure is due to breakage, the initiation might be the result of dynamic response of the window.

In examining the overall design of curtain walls it appears reasonable to inquire as to the possible effects of the vibration behavior of lights as well as other component panels on the structural performance of the assembly. The current procedures for design of building windows principally involves the consideration of the maximum gust loading modified by local shape effects, the resistance of the glazing to rain penetration (and, possible freezing), and the thermal expansion or contraction behavior. The aggregate effect is then considered as a static design load and the glass pane is sized in thickness accordingly. The strength estimations for the panes are based on the statistical results of static tests to failure of numerous laboratory samples (33, 34). There is, of course, an inconsistency in this approach because glass, as it is well known, exhibits a pronounced static fatigue and consequently static failure stresses are often several times lower than those for dynamic loading. Wind loads by their nature represent a dynamic load component and therefore static strength figures may be unduly conservative. This effect is considered to some extent in the specification of a safety factor of 2.5 for this utilization. In contrast, factors of from 5 to 8 are common for static loading that prevails in aquariums or shelving.

The dynamic processes should be present to some degree in curtain wall windows of all sizes but the effects may become more pronounced in larger windows (greater than 20 sq ft). In these cases the natural frequencies are lower and the vibrations are often of higher amplitude. Since pressure loading increases as the square of the dimensions, the restraint forces per unit length along the boundaries are increased in direct proportion to the edge dimensions. This requires correspondingly heavier glazing and magnifies dynamic effects here. Tempered glass because of its strength advantage over annealed plate glass is sometimes used in large windows to conserve weight and hence cost. Unfortunately, the elastic properties (modulus) are largely unaffected by tempering and hence thinner lights possible with this material have lower natural frequencies which aggravate the boundary restraint problem noted above and which also may produce adverse psychological effects.

Finally, while the previous design considerations have implicitly assumed that the cladding system is essentially nonstructural, it is becoming evident from the results of more refined structural analyses and experimental data that architectural details such as partitions, light walls, etc., may provide measureable structural stiffening and damping action (35 - 39). Building cladding in many instances may also provide similar benefits although it is seldom if ever accounted for. Usually the problem is just the opposite: how to minimize the interactive structural effects and transmit only the direct wind pressure or seismic loads to the structure. On the other hand, cladding is generally composed of thin plate-like elements which may have considerable in-plane load carrying ability and shear strength



that might with proper design be utilized to control structural stiffness or energy dissipation. Because of the lack of adequate knowledge as to precisely how loads are transmitted into and through the cladding, it is not currently feasible to directly utilize this component as more than an environmental barrier.

The formulation of a rational design procedure for modern window and curtain wall construction as can be seen from the brief discussion above involves consideration of a large number of factors. There are numerous sources of dynamic loading on a window and the relative magnitudes of the factors among themselves as well as compared to static loads have not been established. At the same time, the dynamic response of windows and claddings is not well understood. Clearly, there is currently the need for a fuller understanding and more complete knowledge of these topics. This is especially appropriate because of the more diverse design and sophisticated analysis of tall buildings, involving, for example, detailed computer structural analysis and model wind tunnel testing.

### 1.3 Statement of Work

The present report describes the results of a two-year program of research aimed at providing some of the answers to the questions raised concerning the localized character of wind loads on windows and cladding and the types of response to these loads. The overall problem is complex and multi-faceted and in certain respects the present work represents an initial step. The study has been broken into several more clearly defined areas, and a description of each along with the overall objectives are given in this section. Results of the research in each area follow in succeeding chapters.

Figure 1.3-1 diagrams the overall organization of the study and shows the grouping into three areas: (i) field study at a 29-story lightweight clad highrise building; (ii) analytical modelling of a representative window assembly and associated cladding for both static and transient dynamic behavior; and (iii) development of a laboratory facility for full-scale dynamic load simulation and testing. The work in each area proceeded in parallel, but was accompanied by extensive cross-flow of data, results, and insights throughout the course of the study. The objectives of the research are outlined in the figure but can be more generally summarized as:

- (i) characterize the spectral and temporal features of the localized cladding pressure loading, especially in relation to building geometry and flow features;
- (ii) identify the major window response mechanisms, both experimentally and analytically, and assess the degree of causality between measured forcing function and response (e. g., identify the degree of response attributable to a particular forcing function);
- (iii) develop an analytical model that can satisfactorily represent both the static and dynamic behavior of a representative window-cladding assembly;
- (iv) develop a capability and methodology for performing full-scale dynamic testing of window or curtain wall assemblies.

The diagrams in Figs. 1.3-2 to -4 show a schematic outline of each of the research areas and define the basic approaches towards accomplishing the objectives. The field study (Fig. 1.3-2) involved at the outset a parallel effort to obtain wind tunnel model data and to develop the instrumentation and measurement techniques for full-scale testing. The model data was used for preliminary identification of "critical" regions or other areas of interest, while the instrument development phase was aimed at constructing the specialized transducers

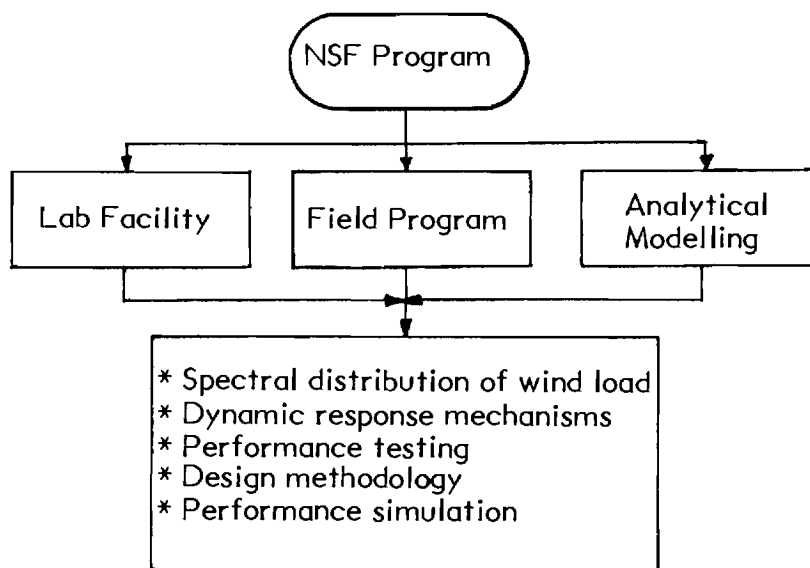


Figure 1.3-1 Program Organization

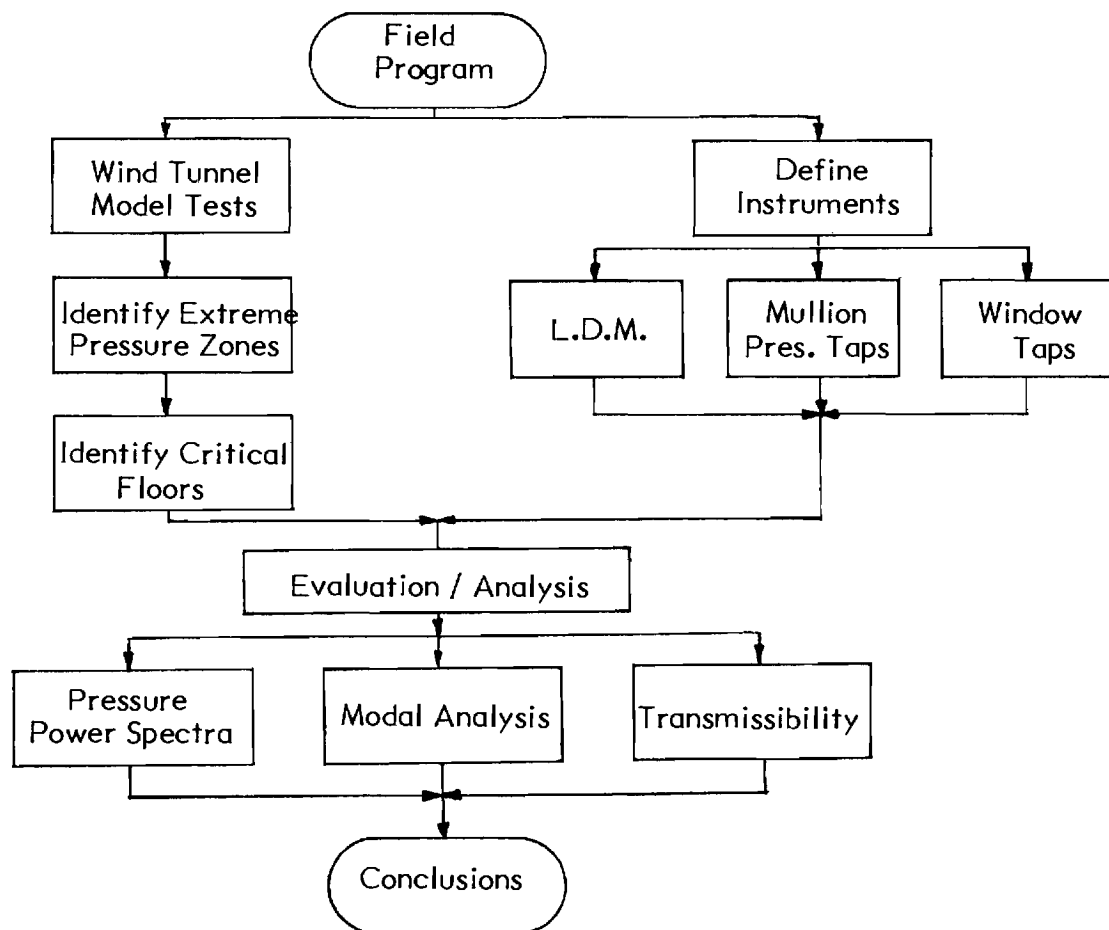


Figure 1.3-2 Field Study Program

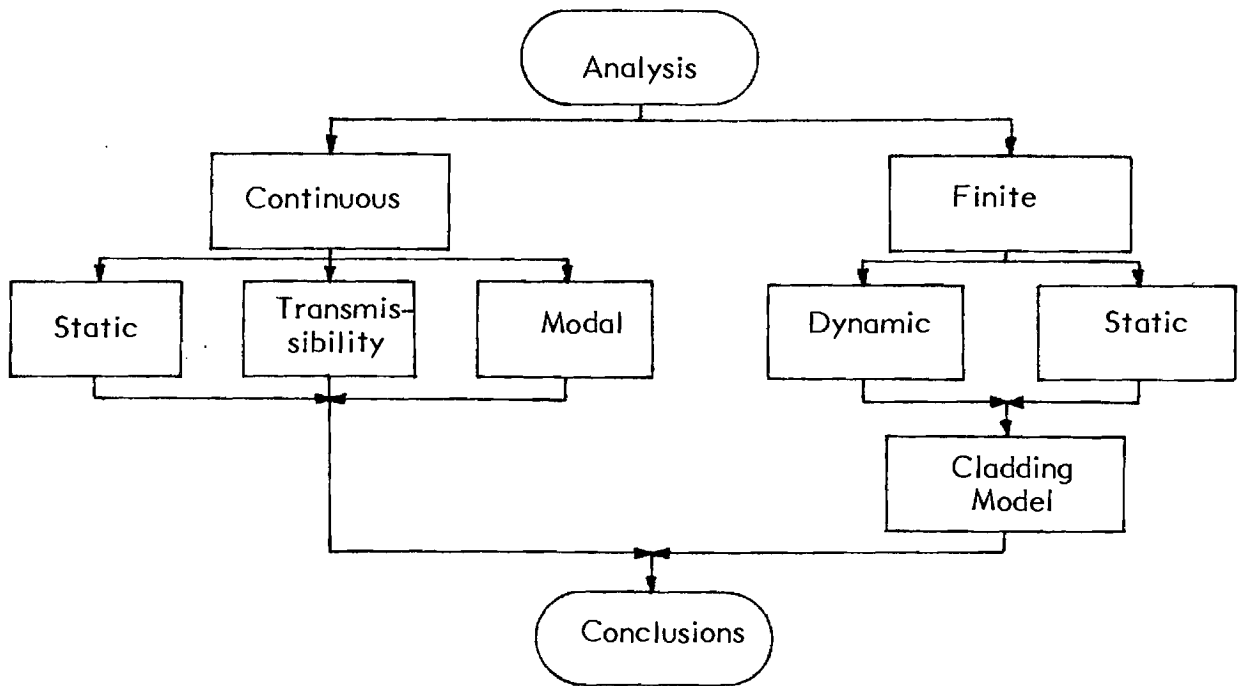


Figure 1.3-3 Analysis and Modelling Program

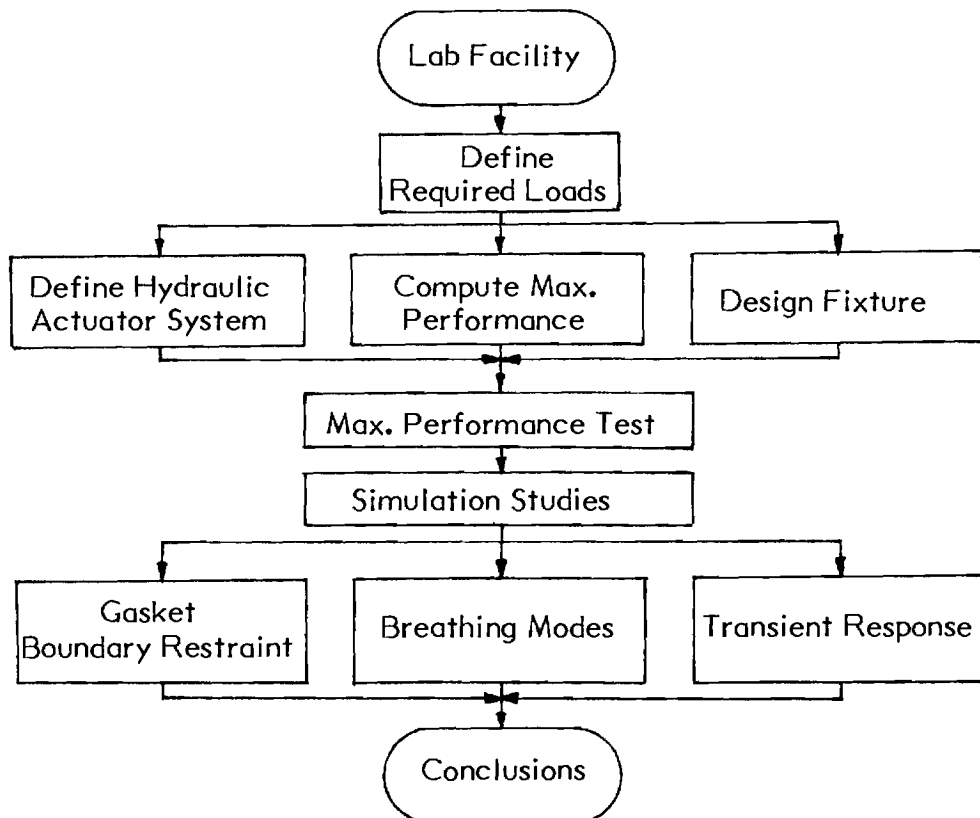


Figure 1.3-4 Laboratory Study Program

and techniques for the subsequent testing. This was followed by development of the measurement analysis techniques for evaluation of the data and finally by the actual testing (40).

The analytical study, Fig. 1.3-3, is intimately related to the other two areas, and for simplicity, can be broken into continuous and finite analyses. The continuous modelling was designed to provide the qualitative, if not quantitative, characterization of both the static and dynamic response to directly support the field and full-scale lab testing. Particular emphasis was on describing the dynamics of a double plate window assembly. The finite modelling, on the other hand, was designed to provide a detailed quantitative description of the window and cladding system that could be "calibrated" against the observations from full-scale and lab measurements. The ultimate objective for utilization of a model of this type is as a design-oriented subsystem model for incorporating cladding effects into more general building models. The basic results are a series of dimensionless parametric performance curves for representative window-framing systems along with a detailed listing of the finite element coding developed for the study.

The laboratory study, Fig. 1.3-4, was aimed at defining the requirements for a full-scale dynamic pressure simulation facility and then building the apparatus. Following a series of performance demonstration tests, the equipment was used to study under laboratory conditions the static and dynamic response characteristics of a window-framing system similar to that employed at the building used for the field study.



## 2.0 FIELD STUDY

### 2.1 Background

The nature of wind forces acting on windows is determined not only by the wind characteristics, but also by the wind-structure interaction phenomena. Even though wind forces on buildings have been of concern in the past, the fluid mechanics related to buildings in atmospheric boundary layer winds is not well understood. Unlike the aerodynamics of streamlined bodies which is highly developed for aeronautical applications, the aerodynamics of bluff buildings in turbulent, nonhomogeneous boundary layers is complex and is not well described by either analytical or numerical techniques. Consequently, the available wind force information for buildings and structures has been obtained primarily from measurements on small scale models and in limited instances on actual structures.

Boundary layer wind tunnel measurements on tall buildings have shown that the wind-structure interaction gives rise to fluctuating local wind pressures. It has been recognized that, whereas the temporal and spatial fluctuations of the local wind pressures from the mean may not affect the design of the primary structure, their effect on portions of the building such as windows and cladding panels cannot be ignored. Model tests have also shown that the most severe pressure fluctuations are associated with the following mechanisms:

- a. Separation and reattachment of flow around buildings (41).
- b. Local vortex formation and the subsequent shedding of these vortices.

Flow separation around buildings can induce locally high suction (outward acting) pressures which often exceed the dynamic pressure of the mean wind. Probability density estimates of pressure fluctuations on tall building models indicate (25) that in regions of separation, negative values of the peak pressure coefficient,  $C_{p_{\max}}$ , can attain magnitudes of 8 - 10 times the rms pressure coefficient,  $C_{p_{\text{rms}}}$ . This is strongly affected by the geometrical features of the building and the most severe pressure fluctuations have been found to occur at reentrant corners where flow separating at one edge of the corner reattached at the opposite edge (41).

Vortex flows may be developed by a particular combination of building shape and wind direction, and are responsible for periodicities in the cladding loads. Model studies of the Bank of America World Headquarters Building (42), show that fluted walls in conjunction with setbacks at various elevations are the main reasons for the appearance of various vortex flows over the entire wall surface. The shedding of these vortices is one of the primary aerodynamic phenomena which develop periodic forces on the building as well as its cladding.

The frequency content of both types of pressure fluctuations is of concern in defining wind loads for window glass and cladding panels, since sufficient energy at frequencies in the vicinity of panel resonance can result in dynamic amplification of deflections and stresses. However, the available spectral data from either model or full-scale tests is rather sparse. Upwind surface pressure fluctuation spectra for a pair of square section towers have been measured on a model (43) and show that most of the pressure fluctuation energy is

concentrated in the low frequency range extending to about 20 Hz. These spectra have been found to be related closely to turbulence spectra of the simulated oncoming wind.

Full-scale tests on existing structures have provided greater insight into the nature and severity of surface pressure fluctuations on tall buildings. Some of the pertinent full-scale tests are surveyed in Table 2.1, partially reproduced from Reference 19. Newberry et al. (44) from their measurements on an 18-story rectangular building conclude that the most severe suction pressures occur under the incidence of glancing winds and large negative excursions are apparent at locations around corners. Dalgliesh (45) in comparing model and full-scale pressure coefficients, points out that large values of  $C_{p_{rms}}$ , and hence large pressure fluctuations from the mean, occur when the wind is nearly parallel to the building surface on which the tap is located. The data also demonstrate a reduction in wind pressures as the ground level is approached. Full-scale measurements by Dalgliesh (46) have provided valuable information on another aspect of the wind loading mechanism, namely: gust effects. Gust action usually is considered as random turbulent fluctuations superimposed on the mean wind and is generally defined by its time domain characteristics. Statistical analyses of extreme pressure data were used to define static pressure magnification or gust effect factors for various exposure conditions and expected recurrence intervals.

Results of frequency decomposition of surface pressure fluctuations on actual structures have been presented in Refs. 47 - 50. Pressure spectra in Refs. 47 and 48 have been obtained over a range of low

Table 2.1 A Partial Survey of Full Scale Investigations of Wind Effects on Tall Buildings

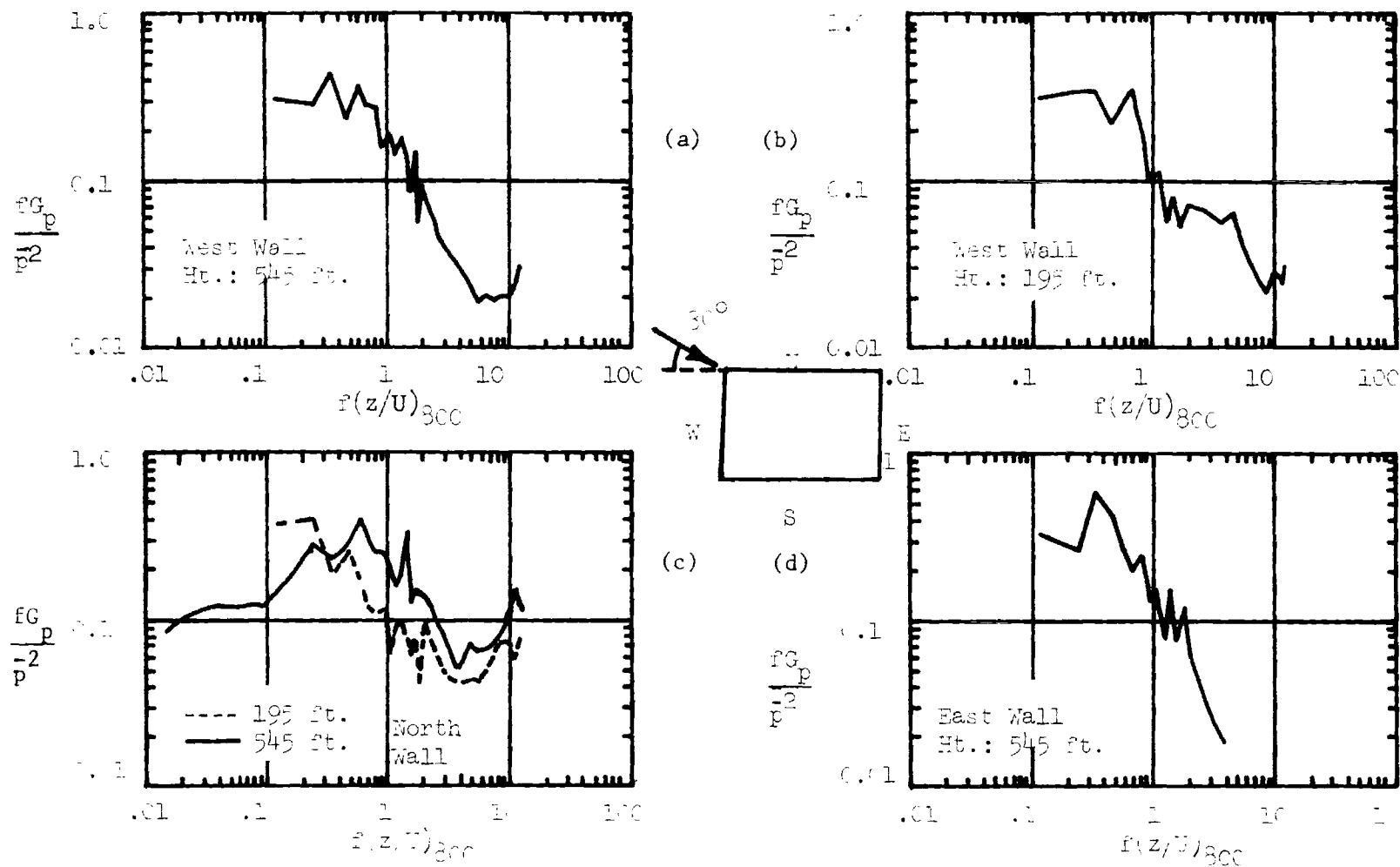
Investigators	Ref.	Building		Experimental Procedure				Analysis		
		Type	Ht. in ft.	Pressure		Wind		Pressure		
				Ext.	Int.	V	Dir.	Mean	Histo- gram	PSD
C. W. Newberry K. J. Eaton J. R. Mayne	44	R. C.	267	48	3	1	1	x	x	x
W. A. Dalgliesh W. von Tobel H. S. Ward	45, 46	St. Fr.	784	34	1	2	2	x	x	x
W. A. Dalgliesh W. von Tobel H. S. Ward	47	St. Fr.	430	12	0	1	1	x		x
Mario Takeuchi Gengo Matsui Satora Kazama Ryoichi Nagai	48	Comp. St./C	251	20	0	1	1	x	x	x
W. A. Dalgliesh W. von Tobel H. S. Ward	49	St. Fr.	640	12	0	1	1	x	x	x
S. Miyoshi Masahiko Ida T. Terayama	50	St. Fr.	482	6	6	2	2			x

\*Abbreviations: St. = Steel; Fr. = Frame; R.C. = Reinforced Concrete;  
Comp. = Composite Steel & Concrete

frequencies pertinent to the vibrations of the entire building. Since the fundamental frequencies of window glass and cladding panels are typically in the range of 5 - 50 Hz, the results of Refs. 49 and 50 are significant in that the spectral bandwidth exceeds 10 Hz. Standen et al. (49) have obtained power spectra of pressures at various locations on the surface of a rectangular platform high-rise building, Fig. 2.1-1. rapid roll-off of power density with frequency is evident in most cases. However, the increased power at higher frequencies in (c) with the pressure taps subject to glancing winds shows that some locations on the surface may be susceptible to dynamic loads.

Measurements of Miyoshi, et al. (50), on a tall rectangular building appear to be the first data which provide response information in terms of glass panel strains in addition to pressure spectra. The pressure spectra, measured under typhoon conditions, are reproduced in Fig. 2.1-2. A similar feature to that of Fig. 2.1-1c is apparent in curve,  $P_E$ , where the pressure fluctuations for glancing flow contain significant energy at higher frequencies as compared to the sharp roll-off observed in the other spectra.

In the past few years the problem of determining wind loads for cladding design has centered on the assumption of static or quasi-static pressures, and, therefore, on the specification of static loads for cladding panel design against outright fracture. Current building codes, as for example the ANSI Standard, allow for large pressure fluctuations and gusts by augmenting the static loads with gust and terrain factors. The basis for a static pressure assumption is the



Symbols:  $f$  = frequency,  $(z/U)_{800}$  = Height to velocity ratio at 800 ft.,  $G_p$  = PSD  
 $\bar{p}^2$  = Mean square value of the pressure,  $U = 50$  ft./sec.

Figure 2.1-1 Power Spectra From Full-Scale Tests of Standen et al. [49]

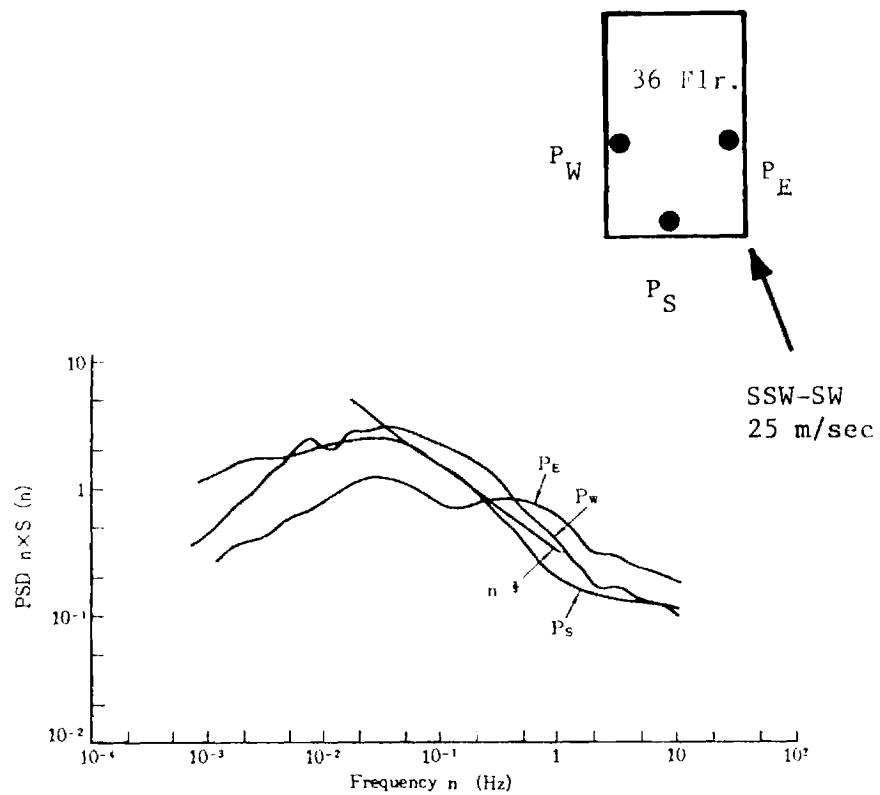


Figure 2.1-2 Power Spectra From Full-Scale Tests of Miyoshi et al. [50].



preliminary pressure fluctuation data for buildings with simple geometries which generally show that the wind pressure energy is concentrated over frequencies considerably lower than the resonance frequencies typical of cladding construction. It has been acknowledged, however, that the nature and severity of local dynamic wind loads on buildings with complicated shapes such as those with fluted walls, reverse setbacks, and surface irregularities may not follow this pattern (51,52).

The results of Refs. 49 and 50 discussed above, hint at the existence of "critical response locations" where intense high frequency pressure fluctuations would necessitate consideration of dynamic loading mechanisms and wind tunnel tests have illustrated that these critical response locations should depend on the nature of the wind-structure interaction. Relatively large windows and cladding panels in these locations would be susceptible to vibration and flutter, and damage in the form of glass panel breakage, cladding panel fatigue failure, and panel anchor failure. Therefore, design of window glass for static loads, in such locations would be inadequate. In addition, design to minimize breakage alone may not in itself be a solution to the window failure problem.

Wind loads for buildings under design can presently be obtained only by measurements on models in wind tunnels. However, to establish a systematic set of simulation criteria and to ensure the predictive quality of model test results, measurements on actual structures and model/full-scale comparisons are essential. Full-scale tests are also of value in understanding the wind loading mechanisms and the

characteristic response. However, since full-scale experiments pertain to particular structures with unique properties and exposure conditions, an ensemble of such measurements need to be examined prior to rational definition of wind loads and formulation of design procedures.

This chapter reports the results of a full-scale study of the characteristics of the localized window and cladding pressure loadings. This approach to the problem of window design has been adopted to understand the dynamic response characteristics of real window installations and the nature of dynamic wind loads, particularly the heretofore unreported dynamic effects under so-called "normal" wind conditions.

The study concerns itself with the problem of dynamic wind load and response identification for doubly glazed windows. Mathematical modelling of the window response, both static and dynamic, is the subject of Chapter 3, and modelling of the wind structure in the strictest sense is not within the scope of this work. The principal objective was to assess the wind loads and window dynamic response with regard to:

- (a) Capability of the wind loads to excite window vibrations and the dependence of the wind loads on the building profile.
- (b) The degree of dependence of the window response on wind loads from amongst other excitation sources.
- (c) Vibration parameters of the window system from measured data.

## 2.2 Design Considerations

In formulating rational design procedures for the elements of a glass curtain wall, the first step is to translate the wind environment of buildings into a set of numerically-defined loads. The kinds of failures that these loads might cause must then be studied in light

of the different response mechanisms. This process leads to the formulation of a set of failure criteria. Once the loads have been defined and the failure criteria established, the design problem reduces to one of predicting the response of the window system, and evaluating this response in terms of the damage criteria.

In a practical application this reduces to the situation: Given a design wind velocity, one should be able to predict the response of the system in the form of peak stresses, accelerations deflections, etc. This requires mathematical models for the prediction of the forces of interaction of the main structure and the turbulent, gusty wind. Wind-structure interaction forces can be aerodynamic as well as aeroelastic, but insofar as window systems in building structures are concerned, only aerodynamic forces are important. Given the current state-of-the-art, the greatest uncertainty is in defining the wind environment of buildings and translating this into cladding loads.

The American National Standards Institute ANSI 58.1-1972 standard (11) contains a method for determining wind loads applicable to the design of parts and portions of a building. Design wind pressures for window glass are assumed to be static and are defined by:

$$P_1 = .00256 V_{des}^2 K_z \cdot G(C_p - C_{pi}) \quad (2.2-1)$$

where

$P_1$  is the design pressure in  $\text{lbs/ft}^2$  (Psf);

$V_{des}$  is the fastest mile wind speed with a 50-year recurrence period, in MPH;

$K_z$  is the velocity pressure coefficient which depends on the exposure and the height  $z$  above ground;

G is the gust response coefficient;

$C_p$  is the external pressure coefficient; and

$C_{p_i}$  is the internal pressure coefficient.

Minimum values of  $C_p - C_{p_i}$  for locations susceptible to large suctions, for example, near corners, are specified. The standard was founded on the available wind research and was prepared for use in conjunction with common building shapes (i.e., rectangular), and is similar to the approach employed by the National Building Code of Canada (53).

Current window design procedures, as for example in the Southern Standard Building Code (54), assume the design pressure as a uniformly-applied static design load. With this load, a factor of safety representing a value which has appeared to give an acceptable probability of breakage in extensive static strength tests conducted by window glass manufacturers (12), is used.

Recent efforts at improving upon the static load specifications have been directed at modifying the gust factor, G. In Ref. 51 it is suggested that G be expressed as:

$$G_1 = 1 + g \cdot \frac{C_{p'}'}{\bar{C}_p} \quad (2.2-2)$$

where g is a gust effect factor;

$C_{p'}$  is rms value of the pressure coefficient (about the mean);  
and

$\bar{C}_p$  is mean pressure coefficient.

The gust effect factor,  $g$ , from Eq. (2.2-2) has values lying between 3 and 5 for windward faces, as determined experimentally. Allen and Dalglish (26) have proposed a gust factor  $G_2$  which incorporates effects of dynamic amplification. These dynamic effects are included in the gust factor as a "gust material" factor,  $K_m$ , with the reasoning that the effects are related only to the turbulent component of the wind. The gust factor is then expressed as:

$$G_2 = 1 + g \cdot K_m \cdot \frac{C_p}{C_p} \quad (2.2-3)$$

Window design requires consideration of the mechanical properties of glass, especially its brittle behavior. The statistical variation in glass strength, due to its sensitivity to surface flaws, is accounted for by designing on the basis of probability of breakage. Design procedures for window glass considering its static fatigue\* property have been proposed, where gust factors are replaced by pressure coefficients averaged over a period determined from the strength-load duration curves (55, 56). Consideration of edge supports, clearances and glazing details has been emphasized in a design method (29) based on static pressures defined by the ANSI standard.

The widespread adoption of static design loads is a direct consequence of the available experimental data. The experimental data base

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\* In the accepted definition "static fatigue" implies decrease in breaking stress with increased duration of load.

has also provided valuable insight into the wind-structure interaction phenomena. However, the model test results are beset by a lack of confidence in their predictive quality. Full-scale tests have been limited in their scope by the emphasis on measurements for validation of wind tunnel test techniques, and have used buildings with simple shapes. The local nature of failure of windows, especially those near the corners of a building, has been recognized, but assessment of factors influential in specification of wind loads for curtain walls is as yet incomplete. It is not unreasonable to suspect that the severe pressure fluctuations in the separation region at a corner of the building could excite the window dynamically as well. Buildings with the following shape characteristics are prime candidates for localized regions of relatively high frequency loading:

- (a) Setbacks or reverse setbacks along the height;
- (b) Fluted walls presenting a cascade of triangular outsets and reentrant corners to the flow;
- (c) Rough surface texture of the building facade as dictated by the architectural details.

These "critical response locations" would also depend upon the directional properties of the wind.

In regions where dynamic wind loads are severe enough to offset the alleviating effect of glass static fatigue, dynamic response of the glass panels becomes an important consideration. Ejection type failures, where relative motion between the glass and its frame is suspected, accentuate the need to define the response of the entire window assembly. A review of curtain wall damage caused by a hurricane (57) shows that all the components of a window assembly should be included in a design philosophy.

## 2.3 Site Specification

### 2.3(a) Test Site Description

The structure selected is a 29-story steel frame medium high-rise building of modern design located some six miles north of the downtown Atlanta area on the crest of a ridge and one-half mile from the nearest comparable building to the northeast. The lightweight exterior curtain wall makes extensive use of aluminum extrusions in its design. The cladding consists of large (100 in x 57 in) doubly-glazed windows incorporating a heat-reflecting exterior light and smaller singly-glazed spandrel panels. A view of the building just after completion is shown in Fig. 2.3-1, and details of the mullion and glazing are shown in Fig. 2.3-2.

The 400-ft-tall structure is essentially prismatic with a plan-form that includes twelve exterior right corners at the base. The geometry is modified for the upper twelve floors to provide twenty and thirty-six corner floor plans, Fig. 2.3-3. In the aerodynamic sense, the building is most interesting with profiles to the wind which change with height. The architectural features contribute to a striking design, and provide the variable geometry so important to the present investigation.

The prevailing wind direction for the Atlanta area is the northwest, and the exposure of the building to prevailing winds can be described as a heavily-forested, relatively flat region of residential construction. The terrain in other directions with respect to the building can be similarly described with surrounding buildings 3 - 4 stories high. The exposure conditions determine the vertical velocity



Figure 2.3-1 Field Site in Atlanta; View from the Southeast  
Relative to the Building Axes.



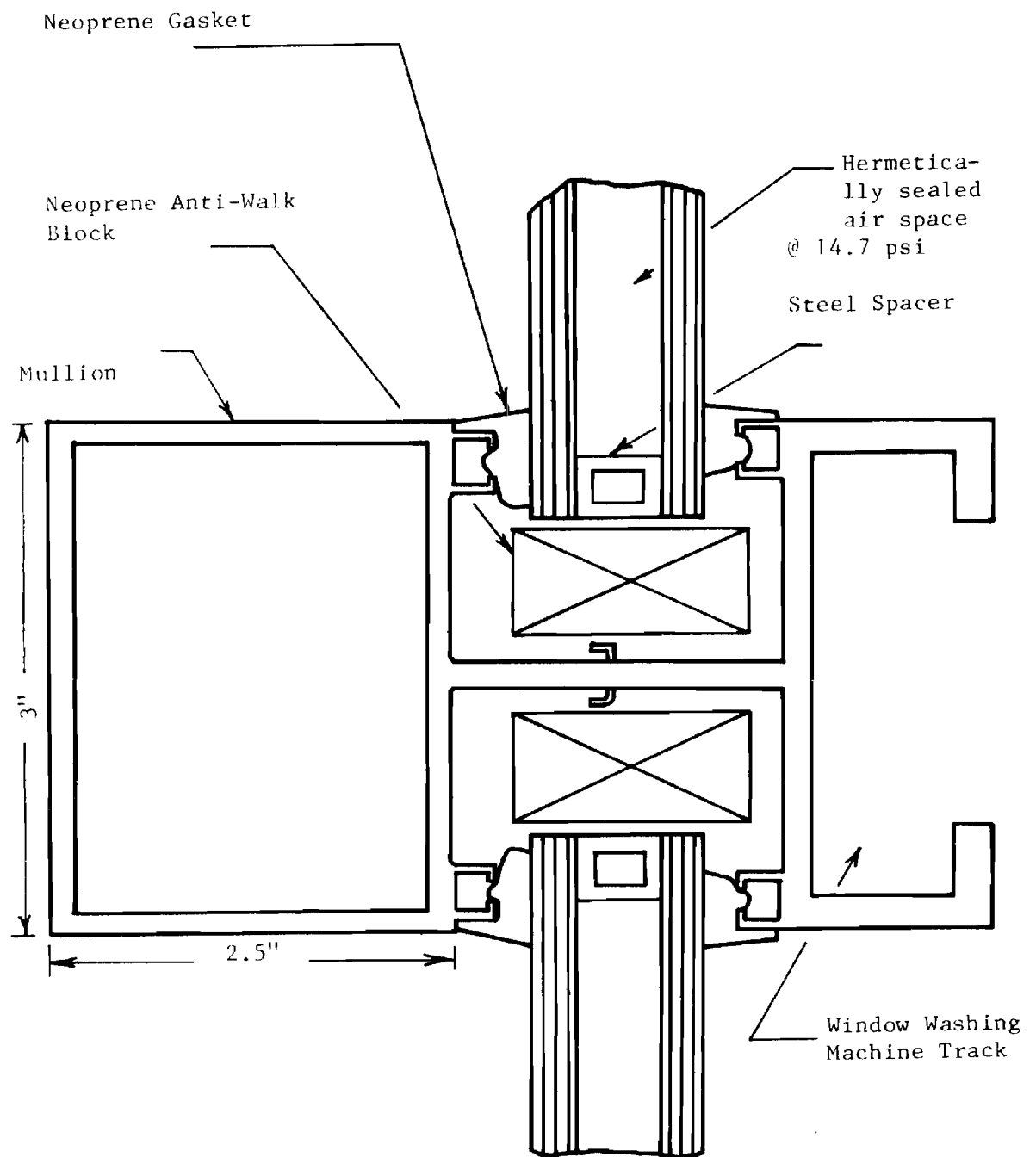


Figure 2.3-2. Mullion and Glazing Detail.

profile of the mean wind boundary layer which can be approximated by the relation

$$\frac{v}{v_o} = \left[ \frac{z}{z_o} \right]^{1/n} \quad (2.3-1)$$

where  $v_o$  is the measured wind speed at height,  $z_o$ , and  $v$  is the velocity at a height,  $z$ . The exponent,  $n$ , depends on the aerodynamic roughness of the surrounding terrain and in accordance with Ref. (103) in following relationships apply:

<u>Terrain</u>	<u>1/n</u>
Open Sea	0.1
Flat Open Country	0.16
Wooded Forest	0.28
Urban Area	0.40

The velocity profile for the wooded area surrounding the building can be approximated by:

$$\frac{v}{v_o} = \left[ \frac{z}{z_o} \right]^{0.28} \quad (2.3-2)$$

### 2.3(b) Wind Tunnel Study

At the outset of the research program, a scale model wind tunnel investigation of the flow patterns around and the wind-induced pressures on the building was conducted in the 9-ft. low speed wind tunnel facility at Georgia Tech. The purpose of the study was to obtain an overall qualitative view of the flow characteristics around the building for the given terrain and to determine flow directions and response locations that resulted in the most severe pressure loadings. Details of the study have been reported in a GIT-DGSA report (58).

The tests were run using a special flat floor in the tunnel with the 1/160 scale model placed at the center of a 4-ft-diameter turntable. The desired mean wind profile was developed by arranging a series of 1/4-in. wire screens in the throat of the upstream contraction cone and measuring the velocities with a pitot-static rake. Due to the limited upstream space available in the test section, no attempt was made to accurately model the turbulence characteristics. Due to 2-axis symmetry only half the model needed instrumentation, and a total of 88 flush pressure taps were located in horizontal rows at 4 vertical elevations and in 8 positions on the roof. In order to minimize tubing length and thus improve dynamic response, the pressure taps were connected to two 48-port "Scanivalves" and a variable capacitance electronic manometer located inside the model. The stepping valves and pressure transducer were operated under minicomputer control for rapid port-by-port sampling in any specified sequence.

All tests were run at an indicated airspeed of 50 mph as measured at an elevation equal to the building height and the pressure measurements reduced to dimensionless pressure coefficient form. Flow directions of from  $0^\circ$  to  $180^\circ$  were used with increments of  $10^\circ$  or  $5^\circ$  as required. Both tuft and oil flow studies were used to explore an unusual vortex pattern over the roof of certain incidence angles.

The results generally revealed no anomalous behavior or unusually large pressure coefficients. As might be anticipated, the maximum (negative) coefficient was measured on the flat side of the model at the higher elevation (Figure 2.3-3(c)) for a glancing flow. Under the relatively low turbulence level achieved with the screens alone, the

maximum coefficient magnitude was about 1.4 which should, on the basis of additional oscilloscope measurements, be increased to about 2.3 to account for the averaging effect of the integrating voltmeter used. The most extreme pressures occurred over the upper half of the model wither on the flat faces for glancing flows or within the set-back regions for  $\pm 45^\circ$  (to building axes) flow directions. Based on these very simple and basically qualitative measurements, locations in the set-back and flat regions at 3 elevations were identified for testing in the full-scale study.

#### 2.4 Evaluation Methodology

In much of the full-scale test work reported to date the primary emphasis has been placed on the logistics and techniques for acquiring the data. Relatively little study has been given to the test methodology, that is, to the manner in which the data will be handled from acquisition, through processing, and finally to evaluation, in light of current analytical understanding. In full-scale tests, and to a large extent in boundary layer simulation tunnel testing, one must deal with a system in which the loads, or inputs, cannot be defined in a deterministic manner. Consequently, statistical techniques must often be employed. Unfortunately, much of the work has dealt with basic time domain statistics such as means, moments, probability distributions, etc., while the application of the frequency domain has been either overlooked or misapplied.

In the present study, the wind load-cladding response problem is modelled as an input-output system with the dynamic wind loads constituting the input and the measured responses as the output. The loading

and response problem has been purposely cast in this form to allow the application of many of the powerful time series analysis techniques. The following subsections describe the evaluation methodology in general and its unique application to the present loading-response identification problem.

#### 2.4(a) System Description

The window-cladding problem can be represented by the generalized model shown in Fig. 2.4-1. The present concern is with the inverse problem in engineering where the input to the system and the resulting response are measured and are to be used to describe the system dynamics as well as the input characteristics.

The input, i.e., the wind loading, itself can be visualized as the out-put of another system describing the wind flow interaction with the structure. In the present study this wind-structure interaction system identification is not considered. In this context the term load identification implies a study of the measured wind loading with respect to a set of parameters describing the wind characteristics and the building profile.

As is the case in the majority of real-world large systems, the window as a structural system is expected to have nonlinear and time dependent properties. These properties, for example, could be due to the fact that a window unit is an isolated part of an entire curtain wall, and that glass strength is sensitive to surface flaws which are affected by service load duration. The identification problem in this case can be formidable (59). In this investigation, the nonlinearities of the system and its time dependent character are neglected as a first

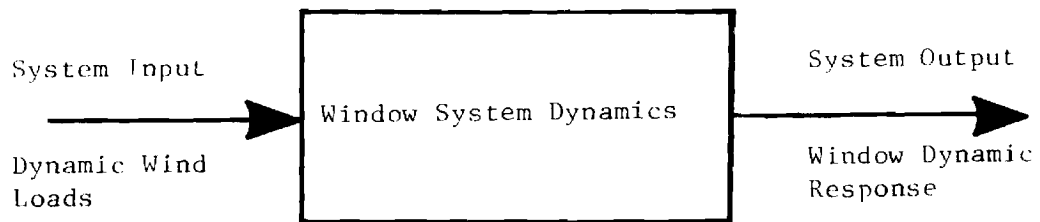


Figure 2.4.1. The Window System as an Input/Output Model.

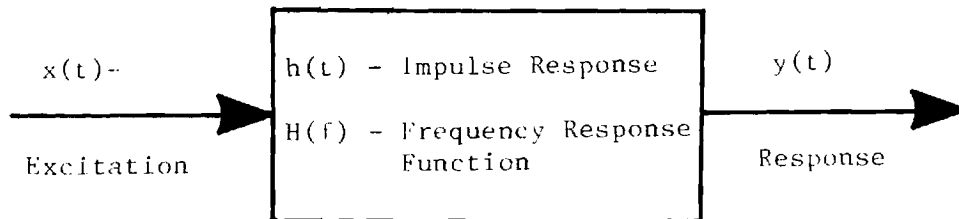


Figure 2.4-2 Single Input/Single Output Linear System.

approximation. The system model is thus simplified so that it can be described by linear constant coefficient differential equations.

Linear system theory is the mainstay of the present data analyses procedures. In this approach the window assembly is assumed to be the linear system and the dynamic wind loads considered as the band-limited random input. The theoretical concepts are not new; however, their present application to large structures is broader in scope, specifically, with respect to the use of the coherence function in assessing the linear dependence of the response on the wind loads.

A linear, time-invariant system also called a constant parameter linear system (CPLS) is one which can be described by linear differential equations with constant coefficients. The dynamic characteristics of a constant parameter linear system can be described by its impulse response function,  $h(\tau)$ . The input-output relations for such a system, represented by the block diagram of Fig. 2.4-2, can be described by the convolution integral:

$$y(t) = \int_0^{\infty} h(\tau) x(t - \tau) d\tau \quad (2.4-1)$$

where a physically realizable system,  $h(\tau) = 0$ , if  $\tau < 0$ . The system can also be described by a transfer function,  $H(s)$ , which is the Laplace transform of  $h(\tau)$ :

$$H(s) = \int_0^{\infty} h(\tau) e^{-s\tau} d\tau \quad s = a + jb \quad (2.4-2)$$

An alternate representation of the transfer function is the frequency response function where  $s = j\omega = j2\pi f$ . The frequency response function is the Fourier transform of  $h(\tau)$ :

$$H(f) = \int_0^{\infty} h(\tau) e^{-j2\pi f\tau} d\tau \quad (2.4-3)$$

Taking Fourier transforms of both sides of Eq. (2.4-1) yields the input-output relation in the frequency domain:

$$Y(f) = H(f) X(f) \quad (2.4-4)$$

where  $X(f)$  and  $Y(f)$  are Fourier transforms of  $x(t)$  and  $y(t)$ . Equation (2.4-4) can also be obtained as a direct result of the Convolution Theorem which states that convolution in the time domain is equivalent to multiplication in the frequency domain and vice versa.

The frequency response function,  $H(f)$ , is a complex valued quantity and can be written in polar notation as:

$$H(f) = |H(f)| e^{-j\phi(f)} \quad (2.4-5)$$

It follows from Eq. (2.4-3) that:

$$H(-f) = H^*(f)$$

where  $*$  denotes complex conjugate, and therefore:

$$|H(-f)| = |H(f)|$$

$$\phi(-f) = \phi(f)$$

If the excitation,  $x(t)$ , is a sample function of a stationary random process  $\{x(t)\}$ , then the output,  $y(t)$  is also a sample function of a stationary random process,  $\{y(t)\}$ , as will be shown. The input-output relations for this case must be described in terms of the expected values of the input and the output. The expected value of any real single-valued continuous function,  $g(x)$ , of the random variable,  $x$ , is given by:

$$E[g(x)] = \int_{-\infty}^{\infty} g(x) p(x) dx \quad (2.4-6)$$



where  $p(X)$  is the probability density function for  $x$ . For stationary random processes,  $\{x(t)\}$  and  $\{y(t)\}$ , the correlation functions are defined as follows:

$$R_{xx}(\tau) = E[x(t) x(t+\tau)] \quad (2.4-7)$$

$$R_{yy}(\tau) = E[y(t) y(t+\tau)]$$

$$R_{xy}(\tau) = E[x(t) y(t+\tau)]$$

The first two of Eq. (2.4-7) define the autocorrelation functions of  $\{x(t)\}$  and  $\{y(t)\}$  respectively, while the third defines the cross correlation function between  $\{x(t)\}$  and  $\{y(t)\}$ .

For the single input-single output linear system of Fig. 2.4-2 the output autocorrelation function can be obtained from the input autocorrelation function as follows:

$$y(t) y(t+\tau) = \int_0^\infty \int_0^\infty h(t_1) h(t_2) x(t-t_1) x(t+\tau-t_2) dt_1 dt_2$$

and taking the expected value,  $E[\ ]$  (30):

$$R_{yy}(\tau) = \int_0^\infty \int_0^\infty h(t_1) h(t_2) R_{xx}(\tau+t_1-t_2) dt_1 dt_2 \quad (2.4-8)$$

Thus the response of a CPLS to a stationary random input is a stationary random process. Similarly, the stationary cross-correlation function

$R_{xy}(\tau)$  between the input,  $x(t)$ , and the output,  $y(t+\tau)$ , is:

$$R_{xy}(\tau) = \int_{-\infty}^\infty h(t) R_{xx}(\tau-t) dt \quad (2.4-9)$$

The transformation of Equations 2.4-8 and 2.4-9 to the frequency domain yields input-output relations in terms of auto and cross-power spectral densities.

$$S_{yy}(f) = |H(f)|^2 S_{xx}(f) \quad (2.4-10)$$

$$S_{xy}(f) = H(f) S_{xx}(f)$$

where  $S_{xx}$ ,  $S_{yy}$  and  $S_{xy}$  are the theoretical two-sided spectral density functions defined over  $-\infty \leq f \leq \infty$  as:

$$S_{ij}(f) = \int_{-\infty}^{\infty} R_{ij}(\tau) e^{-j2\pi f\tau} d\tau \quad (\text{Wiener-Khintchine Relations})$$

From the symmetry properties of stationary correlation functions, it is seen that  $S_{xx}$  and  $S_{yy}$  are real, non-negative, and symmetric functions of  $f$ , whereas  $S_{yx}$  is complex valued.

In practical situations the negative frequencies have no real significance, and it is more convenient to define a practical power density function  $G_{ij}(f)$ , which exists for positive frequencies only:

$$\begin{aligned} G_{ij}(f) &= 2S_{ij}(f) \text{ for } f \geq 0 \\ &= 0 \text{ otherwise.} \end{aligned}$$

Eq. (2.4-10), therefore, can also be written as

$$\begin{aligned} G_{yy}(f) &= |H(f)|^2 G_{xx}(f) \\ G_{xy}(f) &= H(f) G_{xx}(f) \end{aligned} \quad (2.4-11)$$

The cross power spectral density function,  $G_{xy}(f)$ , may be shown to satisfy the inequality (60):

$$|G_{xy}(f)|^2 \leq G_{xx}(f) G_{yy}(f) \quad (2.4-12)$$

Also, from Eq. (2.4-12), a generalized cross-spectral density function, known as the coherence function may be defined,

$$\gamma_{xy}^2(f) = \frac{|G_{xy}(f)|^2}{G_{xx}(f)G_{yy}(f)} = \frac{|S_{xy}(f)|^2}{S_{xx}(f)S_{yy}(f)} \quad (2.4-13)$$

where for random processes

$$0 \leq \gamma_{xy}^2(f) \leq 1$$

Eqs. (2.4-11) and (2.4-13) provide the basis for the physical applica-

tion of linear system theory to the problem of dynamic loading and response of window systems. In particular, the coherence function can be used as a powerful tool for assessing causality in systems and its application in the present study is a point of major interest.

The coherence function, also known as mean squared coherence or the coherence coefficient, was introduced by Wiener (61) in 1930. The terminology was motivated by the interference effects of coherent, and incoherent light sources; however, the "coefficient of coherency" was derived as a natural extension of the theory of harmonic analysis of time series. Wiener also pointed out the analog between the coefficient of coherence and the correlation coefficient,  $\rho_{xy}$ , between two random variables,  $x$  and  $y$ , defined as:

$$\rho_{xy} = \frac{C_{xy}}{\sigma_x \sigma_y}$$

where  $C_{xy}$  is the covariance of  $x$  and  $y$  and  $\sigma_x$  and  $\sigma_y$  the standard deviations of  $x$  and  $y$ , respectively. It was remarked that the correlation coefficient is a tool for the statistical analysis of frequency series, where time does not enter as a parameter; whereas, the coefficient of coherence is a tool in the analysis of time series.

For two stationary time series,  $x(t)$  and  $y(t)$ , the coherence function, Eq. (2.4-13), as a consequence of the Schwarz inequality, Eq. (2.4-12), is restricted to values between 0 and 1. For deterministic signals the coherence function can only take the values 0 or 1. As applied to stationary stochastic processes, the coherence function possesses the properties outlined below, which suggest several of its applications.

Consider the linear system configuration shown in Fig. 2.4-3 (62). It is assumed that  $s(t)$ ,  $n_1(t)$ , and  $n_2(t)$  are sample functions of stationary, uncorrelated random processes. By the application of Eqs. (2.4-11) and the basic relations for multiple input linear systems, (Ref. 60, p. 147), it can be shown that:

$$\gamma_{y_1 y_2}^2 = \gamma_{r_1 r_2}^2 = \frac{G_{xx}^2(f) |H_a(f) H_b^*(f)|^2}{G_{xx}(f) |H_a(f)|^2 + G_{n_1 n_1}(f) \quad G_{xx}(f) |H_b(f)|^2 + G_{n_2 n_2}(f)} \quad (2.4-14)$$

The properties of the coherence function can now be summarized from Eqs. (2.4-13) and (2.4-14) as follows:

1. The coherence function is a real valued function conveniently normalized to lie between zero and unity.

2. Since  $\gamma_{y_1 y_2}^2(f)$  is independent of  $H_1(f)$  and  $H_2(f)$  in Eq. (2.4-14), the coherence function between two time series is not affected by the arbitrary linear processing of these signals. In other words, the coherence function is invariant under linear filtering operations.

3. With  $n_1(t) = n_2(t) = 0$ ,  $\gamma_{y_1 y_2}^2(f) = 1$ . Consequently the coherence function of two signals derived from the linear operations on the same signal is unity. Equivalently, the coherence of two linearly related signals is unity. It can be shown (63) that the coherence function is less than unity for any nonlinearities in the system, and the linear-to-nonlinear ratio (in the absence of noise)

is given by  $\gamma_{y_1 y_2}$

$$\frac{\gamma_{y_1 y_2}}{1 - \gamma_{y_1 y_2}}$$

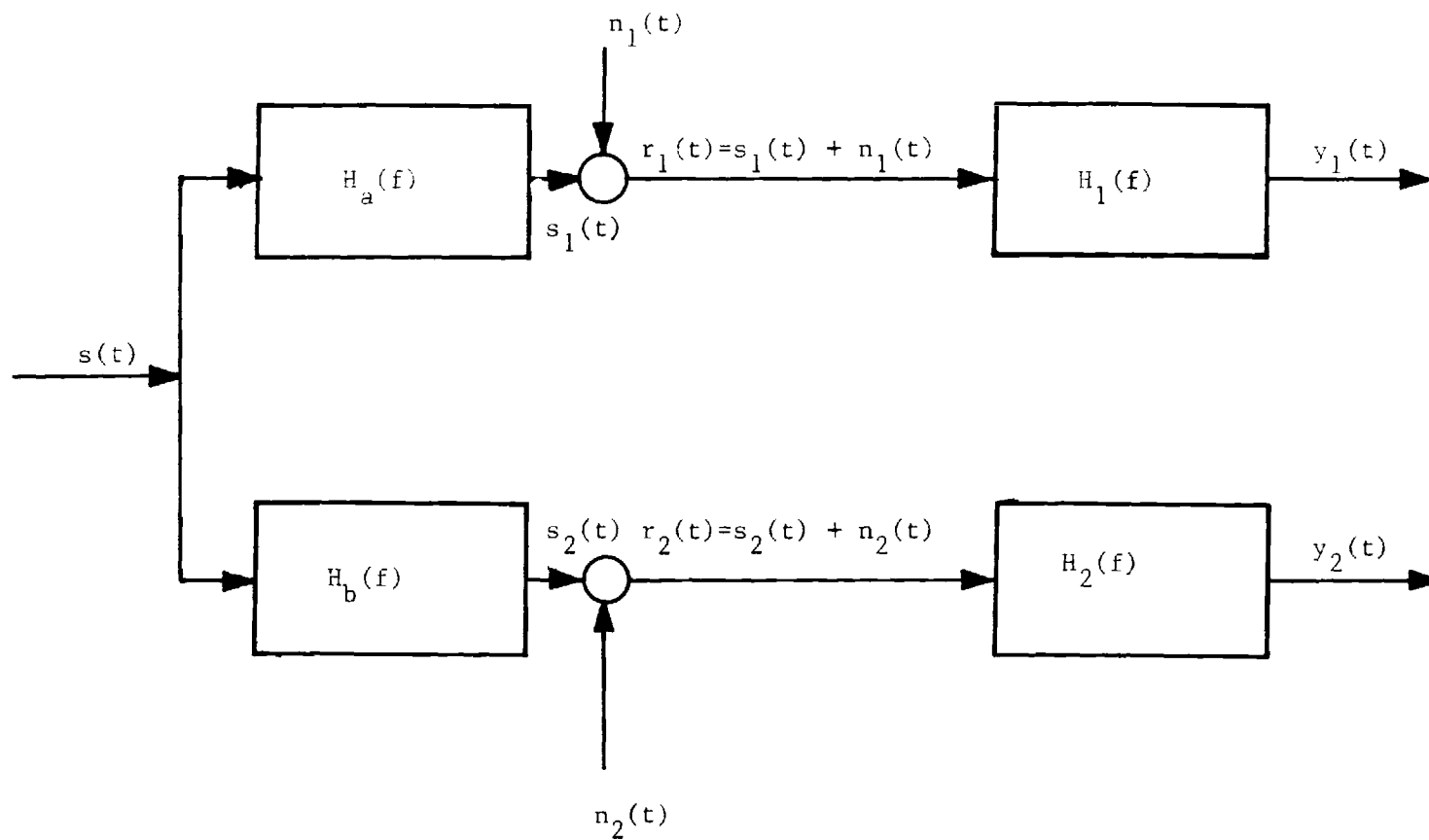


Figure 2.4-3 Configuration for Examining Properties of the Coherence Function.

4. For a linear system, the presence of noise, i.e.  $n_1(t)$  and  $n_2(t)$ , tends to reduce the true value of the coherence function. In such a case

$$\frac{\gamma_{y_1 y_2}}{1 - \gamma_{y_1 y_2}}$$

is a measure of the signal to noise ratio.

5. An interesting property of the coherence function is obtained from its relation to the error spectrum resulting from a linear prediction,  $y(t)$ , of  $y(t)$  from  $x(t)$ , described in (60):

$$y(t) = \int_0^{\infty} h_0(\tau) x(t-\tau) d\tau$$

where  $h_0(\tau)$  is a weighting function to be determined and  $x(t)$ ,  $y(t)$  are sample functions of a stationary stochastic process. When the linear filter is chosen to minimize the mean square value of the error,  $e(t)$ , i.e., the area under the error spectrum, where,

$$e(t) = E[\{y(t) - \hat{y}(t)\}]^2$$

the error spectrum takes the form (31, 34):

$$G_{ee}(f) = [1 - \gamma_{xy}^2(f)] G_{yy}(f)$$

Therefore, the error is small when the coherence function is near unity, indicating that the  $y(t)$  series can be predicted very well from the  $s(t)$  series data. On the other hand when  $\gamma_{xy}^2(f)$  is near zero, the error spectrum very nearly coincides with that of  $y(t)$ , and the  $x(t)$  data contribute almost nothing to the predictability of  $y(t)$ . This property relates closely with the intuitive meaning of coherence, i.e.

two time series can be described as "coherent" at a frequency,  $f$ , when  $\gamma_{xy}^2(f) \approx 1$ . When  $\gamma_{xy}^2(f) \approx 0$ , they are referred to as "incoherent". Alternatively,  $\gamma_{xy}^2(f)$  is the proportion of  $G_{yy}(f)$  contained in the linear component of  $y(t)$ , and  $1 - \gamma_{xy}^2(f)$  is the proportion of  $G_{yy}(f)$  contained in the error, or, in the absence of noise, in the nonlinear component of  $y(t)$ .

Properties (3) and (5) indicate that the coherence function can be used as a measure of linearity. Furthermore, property (5) enables the use of the coherence function in establishing the degree of linear dependence of a system output  $y(t)$ , on the input  $x(t)$ . Property (4) indicates that the coherence function can be used in measuring the signal to noise ratio. It is because of these properties that the coherence function has found widespread application in communication systems in estimating the signal/noise ratio and determination of time delay.

The coherence function has numerous other applications, as for example in control systems for system identification (65), and in geophysics (66) to assess the signal to noise gain of an array of sensors. The earliest applications of the coherence function to the structural dynamics of mechanical systems appear to be those of Kaneshige (67) and Barnoski (68). The coherence function, however, was used only to assess the validity of the estimated frequency response functions. A novel application of the coherence function in acoustic noise measurements (69) is based on the capability of  $\gamma_{xy}^2(f)$  to determine if two signals that produce a spectral line are correlated. This property enables the detection of noise sources

in a mechanical system.

The above application suggests use of the coherence function in the identification of loading sources that produce vibrations in a structural system. Recalling properties 3 and 5, it is seen that  $\gamma_{xy}^2(f)$  can be used to assess the degree of linear dependence of the dynamic response of a structure on a particular loading source. With this point of view, the coherence function can be utilized to assess the causality in a complex system. For example, in the present study the coherence function can be used to evaluate the causality of the wind loads in exciting the dynamic response of window units or other elements of the curtain wall.

Theoretically, a similar assessment can, of course, be made by cross-correlation techniques. In fact, cross-correlation techniques have been applied to isolation of noise sources (70), where the technique was used for the separation of noise at a given point A into components according to the source, the transit time from the source to point A, and the frequency.

The choice of which approach to adopt for a particular problem can only be made when the overall objectives of the study and the methodology to be used have been established. If the problem description is based largely on a frequency domain formulation, for instance, using modal analysis or eigenvalue techniques, a transfer or frequency response function characterization is most appropriate, and use of the coherence function would provide more readily usable information. On the other hand, if a time domain formulation of the problem was made, for example, to study transient response behavior or transit



time effects, then a correlation function approach would likely be preferable. Generally, the type of approach used in any situation will depend on the computational efficiency and the type of information desired, and the choice of using either the coherence or correlation function for source identification purposes must be made accordingly.

#### 2.4(b) Application to Wind Loads and Window Response

General: The window assembly shown in Fig. 2.4-4 which together with smaller spandrel panels formed the exterior curtain wall of the building, is divided into two linear subsystems for the present application, Fig. 2.4-5(a).

Wind loads act on the double plate system formed by the two glass panes and through boundary connections are transmitted to the supports (shear forces in the case of simply-supported boundary conditions). The two systems are assumed to behave linearly under ambient wind loads.\* The supports are assumed to be rigid for the evaluation of the transmissibility function defined below.

The input to the interconnected double plate system is the non-deterministic surface pressure due to the wind. The displacement response of the two plates, Fig. 2.4-4, was considered as the output. The input and the response parameters in the transmissibility model of the supporting structure depend on the definition of transmissibility used. From Fig. 2.4-5 the force transmissibility is defined as the magnitude of the ratio of the total force  $F_1$  transmitted to the plate perimeter to the driving force  $F_0$  (71). The force  $F_1$  comprises a

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\*"Ambient Wind Loads" excludes unusually strong meteorological phenomena such as tornadoes and hurricanes. Observations in strong squalls ( $v > 65$  mph) indicated that static glass pane deflections comparable to glass thickness occur and linearity may not hold.

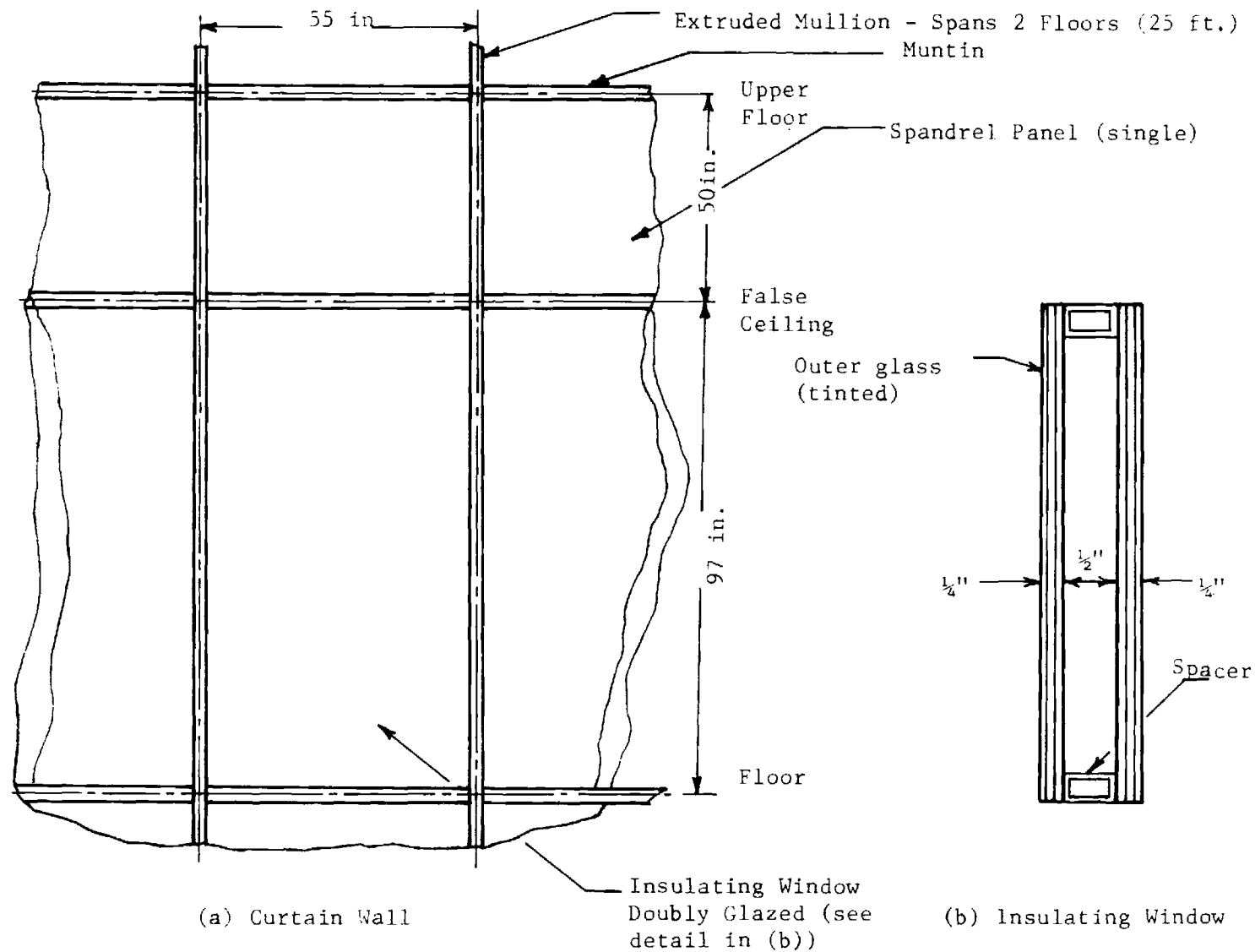
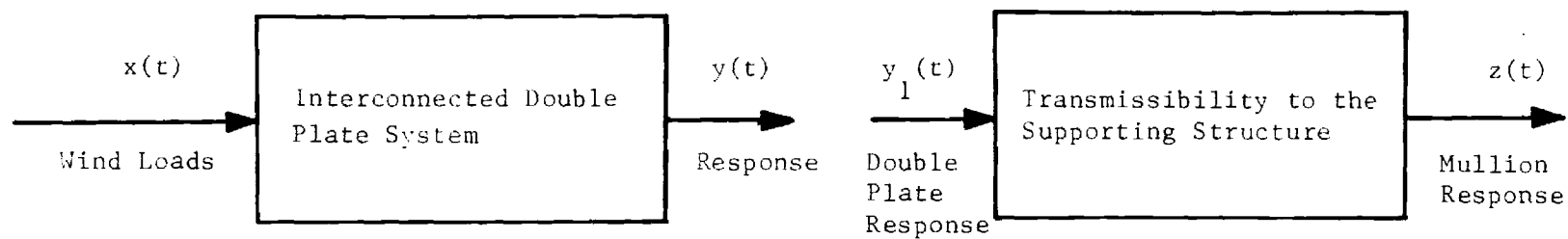
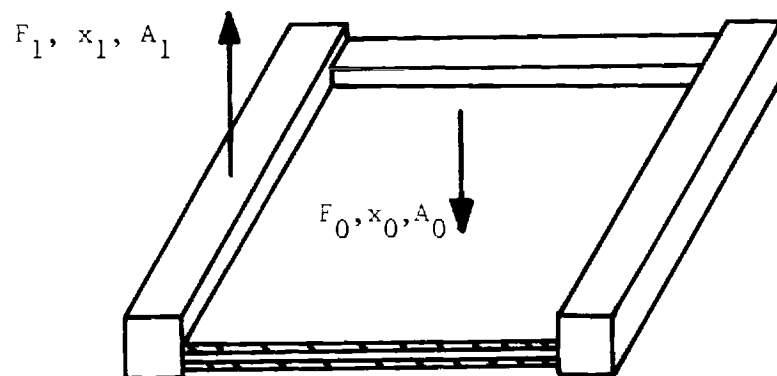


Figure 2.4-4. Components of a Curtain Wall and a Doubly Glazed Window Unit.



(a) Subdivision of The Window Assembly Into Linear Systems.



(b) Definition of Transmissibility.

Figure 2.4-5 Window System Models

distributed force around the plate boundaries plus four discrete forces one at each plate corner. The displacement,  $x_o$ , of the midpoint of the plate to the displacement,  $x_1$ , of its boundary (71). As a result of a general transmissibility theorem (72) the two definitions are identical in the frequency domain. Furthermore, at any given frequency the displacement transmissibility is identical with the corresponding acceleration ratios:

$$|T(f)| = \left| \frac{F_1}{F_o} \right| = \left| \frac{x_o}{x_1} \right| = \left| \frac{A_o}{A_1} \right| \quad (2.4-15)$$

Since the rigid body accelerations can be measured with relative ease, in an experimental study, the acceleration ratio is generally measured.

The structural dynamic behavior of the two linear systems shown in Fig. 2.4-5 can be completely characterized by either their impulse response functions,  $h(\tau)$  and  $s(\tau)$ , in the time domain or by their transfer functions,  $H(f)$  and  $T(f)$ , in the frequency domain. The use of the transfer function, i.e., frequency domain characterization, is convenient in an experimental study where, as in the present case, the system characteristics are not known. Among the reasons for the choice are:

(a) the difficulties associated with interpreting the time domain response of a linear structure where several modes have been excited. This situation is aggravated if the modes are closely spaced.

(b) the ease with which background noise in the measured input and response can be minimized in frequency domain analyses.

(c) use of the Fast Fourier Transform (FFT) algorithm permits rapid analysis of voluminous data on a digital processor thereby

providing greater computational efficiency in the estimation of frequency domain functions, e.g., power spectra.

The disadvantage of the frequency-domain viewpoint is that it is not directly applicable to nonlinear systems. In the present application, however, the system behavior is assumed to be linear and only the steady state characteristics are of interest. Therefore, frequency domain representation of the CPLS is used, with Eqs. (2.4-11) and (2.4-13) providing the basis. The most severe restrictions in the use of this formulation are that the system must be linear, time-invariant, and the input must be a stationary random process.

The Input-Wind Loads: The natural wind turbulence together with the roughness effects on the earth's surface features cause the wind pressures on a building facade to be random in nature. Cladding loads, therefore, must be considered as a random process and described in terms of statistical functions. It appears reasonable to assume that over a short period of time the wind loads can be classified as stationary.\* The frequency decomposition and distribution characteristics of the process can then be described by its power spectral density function (PSD) and amplitude probability density function (pdf).

The wind-structure interaction modifies the free stream wind properties, i.e., the PSD's and the pdf's. As a result, the surface pressure frequency content and amplitude distribution depend on the following four variables:

1. Surrounding terrain conditions or exposure.
2. Mean wind speed, denoted by  $\bar{v}$ .

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\*The classification of cladding loads as stationary in light of the experimental data is discussed in section 2.5.

3. Mean wind direction,  $\bar{\theta}$ .

4. Building profile, assumed to be described by the parameters

$\phi$  and  $z$ , where  $\phi$  denotes the surface orientation at a height  $z$ .

These features of the wind flow past a bluff body may combine to produce the following effects related to cladding loads in the time and frequency domain:

1. Time domain - The descriptor in this case is the amplitude pdf.

(a) Variation in the magnitude of rms pressure fluctuations as indicated by the variance of the pdf.

(b) Variation in gust factors,  $g$ , as defined by Eq. (2.2-2).

(c) Negative peak  $C_p$  excursions as determined from the extreme values in the pdf.

2. Frequency domain - The descriptor in this case is the power spectral density function, and effects of interest in addition to the overall bandwidth relative to the window natural frequencies are:

(a) Periodic forces due to vortex shedding.

(b) Narrow band pressure pulsations.

(c) Modification of the slope of the spectrum from the Kolmogorov inertial subrange spectrum which is proportional to  $f^{-5/3}$ .

(d) Variation in levels of the spectrum. This variation is related to the severity of dynamic loading.

To assess the dependence of these effects on the four variables mentioned above it is necessary to examine the roles each one of these plays in the present investigation.

(1) Terrain Conditions - The terrain surrounding the building used as the test site shows no characteristic variation and, therefore,

can be eliminated from the list of independent variables.

(2) Wind Speed - Wind speed is converted to the nondimensional Reynolds number,  $Re = \frac{\bar{v}d}{\nu}$ , where  $d$  is a characteristic dimension. The Reynolds number does not explicitly appear as a variable since it affects the surface pressures only with a transition from one Reynolds number range to another.

(3) Wind Direction - The wind direction,  $\theta$ , is an explicit variable in that it governs the relative pressure distribution on the building facade.

(4) Building Profile - The building geometry is best described in terms of a coordinate system similar to cylindrical. The coordinates,  $\phi$  and  $z$ , can be visualized as a polar angle and the height, respectively. For the present study, the angular interpretation for  $\phi$  is discarded and instead  $\phi$  is given discrete labels corresponding to the sites on each floor at which measurements were conducted.

If  $x(t)$  is a random variable describing the instantaneous surface pressure, then a realistic functional form of the power spectral density can be hypothesized as:

$$G_{xx}^R(\xi) = F_x(\phi_z, z, \xi, \theta; R_e) \quad (2.4-16)$$

where,

$$G_{xx}^R(\xi) = \frac{G_{xx}(\xi)}{[\frac{1}{2}\rho \bar{v}^2]^2} \quad \text{in units of 1/Hz is the reduced PSD,}$$

and

$\xi$  is the reduced frequency given by  $\frac{f \cdot d}{v}$ .

Local pressures due to wind may not necessarily follow the Gaussian distribution. Wind tunnel measurements (25) show that for the wind-

ward face of a structure, the local pressures follow the same distribution as expected for velocity fluctuations in the approach flow (usually Gaussian for open exposure). For the leeward side, however, the positive local pressures show a generally Gaussian character whereas the large negative pressures are non-Gaussian. It has been pointed out by Davenport (73) that the strong negative tails are exponential in form. Since the pdf depends on the nature of the wind-structure interaction, the mean and the variance of the pdf of  $C_p$ 's can be expected to have a functional form as follows:

$$\begin{aligned} \text{Mean: } \bar{C}_p &= C_1[\theta, \phi_z, z; R_e] = \frac{\int_{-\infty}^{\infty} \Delta p \cdot p(\Delta p) \cdot d\Delta p}{\frac{1}{2} \rho \bar{v}^2} \\ \text{Standard Deviation: } C'_p &= C_2[\theta, \phi_z, z; R_e] = \frac{\int_{-\infty}^{\infty} (\Delta p - \bar{\Delta p})^2 \cdot p(\Delta p) \cdot d\Delta p}{\frac{1}{2} \rho \bar{v}^2} \end{aligned} \quad (2.4-17)$$

where,  $p(\Delta p)$  is the probability density function of the surface pressure,  $\Delta p$ , and  $\bar{v}$  is a mean reference velocity. Experimental investigation of the hypothesis of critical response locations, therefore, reduces to a parametric study of the functional forms of Eqs. (2.4-16) and (2.4-17). Load identification consists of examining the effects of the independent variables on the estimated PSD's and pdf's.

Causality of Wind Loads: There are several possible sources of the dynamic loading acting on a window system as mentioned in 2.1. Therefore, the degree of linear dependence of the response on the direct wind loads must be known in order to establish the causality of wind loads in exciting the window response. This can be accomplished by means of the coherence function, Eqs. (2.4-13) and (2.4-14). Although



its meaning in communication theory has been known for several decades, the use of a coherence function for identification of noise and vibration sources has only recently received attention (69, 74).

The theoretical coherence function is basically a measure of the linearity between two processes. In practical applications, however, measurement noise becomes an important consideration and the real-valued coherence function may be simultaneously affected by nonlinearity in the system and measurement noise. When the system between the input and the output is linear,  $\gamma_{xy}^2(f)$  assumes the value 1. Low values of the measured coherence function are indicative of:

- (i) Extraneous noise present in the measurements;
- (ii) The system relating  $x(t)$  and  $y(t)$  is not linear;
- (iii)  $y(t)$  is an output due not only to  $x(t)$  but also to other unspecified inputs (partial coherence problem).

System linearity for multiple inputs can be estimated from the multiple coherence function. For several correlated inputs the system linearity is determined from the partial coherence functions. In the present application the ordinary coherence is used.

The Response: The frequency response function of the double plate system,  $H(f)$ , can be represented in a matrix form as:

$$H(f) = \begin{bmatrix} H_{11}(f) & H_{12}(f) \\ H_{21}(f) & H_{22}(f) \end{bmatrix} \quad (2.4-18)$$

where  $H_{ij}(f)$  is the ratio of the displacement response at plate  $j$  to an input force at plate  $i$ . The force to displacement transfer function is called the compliance transfer function. The elements of the transfer function matrix of Eq. (2.4-18) may themselves be matrices.

The frequency response function estimated from the measured input and output can be used to determine the vibration properties of the double plate system, namely

- (a) Natural Frequency
- (b) Damping
- (c) Mode Shapes
- (d) Modal Mass
- (e) Modal Stiffness

An excellent survey of system identification techniques is given in Ref. 59. The particular procedure used in the identification of the modal parameters of the double plate system is detailed in 2.5.

The transmissibility function is used to evaluate the effect of the glazing retaining the double glass plate in the supports. Comparison of theoretical and measured acceleration transmissibilities provides an estimate of the reduction in support acceleration and thus forces caused by the neoprene gasket.

#### 2.4(c) Data Processing

The large quantities of inherently nondeterministic data that result from an experimental study required that data processing techniques be fast and fairly streamlined. In recent years, the advent of high performance low cost minicomputers, and the Fast Fourier Transform algorithm have made digital processing of time series of data attractive. The various frequency domain quantities in the transfer function description of linear system dynamics, can be rapidly computed using digital time series analysis procedures. Digital techniques also offer the advantages of ease of manipulating the processed results,

and greater dynamic range. These considerations, and the availability of a digital system to be described later, dictated use of digital computation of the various auto and cross-power spectra from the experimental data.

The mathematical development in 2.4(a) of the frequency-domain input-output relations for a linear system dealt with exact values of the frequency-domain quantities. This means, in the case of random time series, that the spectra are based on time samples of data that are both infinite and continuous. In practice, of course, the sample lengths are finite and in the case of digital analysis are further limited to discrete data values. The effect of truncating the data is that magnitude and phase information in the frequency domain is obtained only at discrete frequencies. The effect of sampling at discrete time increments is to limit the highest frequency at which valid information is obtained. In addition, to determine the descriptive properties of random time series from finite data, estimators, as for example the sample mean and the sample variance, for the various spectral quantities must be used. This means that only estimates of the properties of a random process can be obtained. Thus, it is necessary to recognize the uncertainty or statistical variation in the estimates and know the techniques to minimize this uncertainty.

Digitization of random analog data and estimation of their descriptive properties from finite samples, therefore, requires consideration of the following:

- (a) Sampling at discrete time increments
- (b) Effects of truncation of the random time series

(c) The Estimators used

(d) Statistical errors in the estimates

Practical implications of these aspects of digital time series analysis in computing the various frequency-domain quantities are briefly outlined in the following paragraphs.

Digitization of Analog Data: Digitizing analog data of duration,  $T$ , by sampling in the time domain with an interval,  $\Delta t$ , requires that the frequency spectrum be band-limited to a frequency,  $\frac{1}{2\Delta t}$ . This is a consequence of Shannon's sampling theorem which states that it requires slightly more than two samples per period to uniquely define a sinusoid (75). In terms of the time domain and frequency domain interrelationship this can be written as

$$F_{\max} < \frac{1}{2\Delta t} \quad (2.4-19)$$

where  $F_{\max}$  is the highest frequency present in the spectrum. The sampling rate,  $F_s = \frac{1}{\Delta t}$ , determines the Nyquist folding frequency,  $f_N = 2F_s$ . Frequencies in the spectrum which are greater than  $f_N$  are seen as replicas of frequencies in the base band,  $(0 - f_N)$ , e. g.,  $f_N + f_i$  will be seen as  $f_N - f_i$  (76). This folding back of frequencies is known as aliasing, and the "image" frequencies cannot be distinguished from the true base band spectrum. In practice aliasing can be avoided by low pass filtering the data with a sharp cut-off frequency slightly less than  $f_N$ . The total number of samples is given as:  $N = T\Delta t$ . In digital computation,  $N$  is frequently referred to as the block size.

Sampling the data over only a finite duration,  $T$ , results in discretization of the frequency domain, i. e., the Fourier transform

of the sample can be calculated only at a finite number of frequencies spaced  $\Delta f = \frac{1}{T}$  apart. The frequency resolution,  $\Delta f$ , is related to the Nyquist frequency,  $f_N$ , and the total number of discrete values,  $N$ , in the time domain as follows:

$$\Delta f = \frac{f_N}{N/2}$$

This is because  $N$  discrete values in the time domain allow computation of frequency and phase information in the frequency domain at  $N/2$  points. Theoretically, discretization of the data in the frequency domain at intervals of  $\Delta f = \frac{1}{T}$  is equivalent to assuming the data to be periodic, i. e., the sampled segment repeats itself with a period,  $T$ .

Truncation of a random time series is mathematically equivalent to multiplication of the data by the boxcar function shown in Fig. 2.4-6, and expressed as:

$$\begin{aligned} u(t) &= 1 & 0 \leq t \leq T \\ &= 0 & \text{otherwise} \end{aligned}$$

The data are then said to have been viewed through a rectangular window of width,  $T$ . The Fourier transform of the rectangular window, also known as its line shape function, is:

$$U(f) = \frac{\text{Sin} \pi f T}{\pi f}$$

Computation of Fourier transform of a finite data sample, therefore, is equivalent to taking the Fourier transform of the product of an infinite data record and the time window. The result in the frequency domain is the convolution of the transform of the record of infinite

length with the transform of the time window. For example, if  $G_{xx}(f)$  denotes the true power spectrum of a stationary time series  $x(t)$  and  $\tilde{G}_{xx}(f)$  the net result of convolution with the time window transform, then

$$\tilde{G}_{xx}(f) = \int_0^{\infty} G_{xx}(f_o) U_T(f - f_o) df_o$$

Thus, the true spectrum is "smeared" by the window line shape. This convolution is effected before quantization along the frequency axis. Thus, if the original data are periodic such that there is an integer number of periods in the time window, the line shape does not appear in  $\tilde{G}_{xx}(f)$  because frequency samples occur only at zero crossings of  $U_T(f)$ . However, if there is an odd number of half cycles in the time window, frequency samples occur at the peaks of  $U_T(f)$ , and a single frequency waveform may transform into a broad band of frequencies. This phenomenon is known as frequency leakage and the cure is to reshape the time window to produce a line shape with smaller side lobes (See Figure 2.4-6.). This technique will broaden the main lobe and hence reduce frequency resolution, but it reduces the interference of one frequency with its neighbors.

The selection of a time weighting function to reshape the rectangular window transform and thus suppress leakage reduces to a study of the various windows and their line shapes. A number of weighting functions, also called windows, and their line shapes are described in Refs. 60, 75, 76, and 77. For the present application the Hanning window of Fig. 2.4-7 described by

$$w_H(t) = 0.5 \left( \cos \frac{2\pi t}{T} - 1 \right)$$

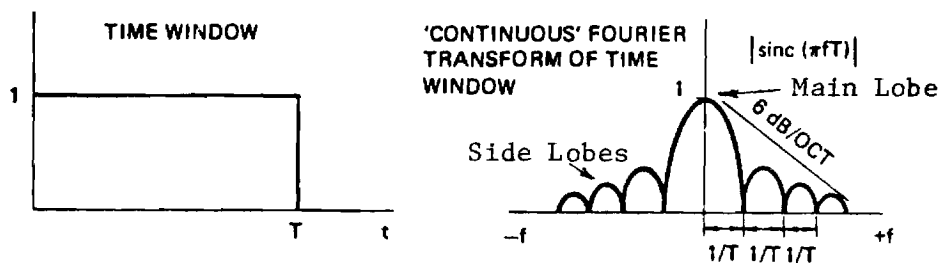


Figure 2.4-6. Boxcar Function and Its Fourier Transform.

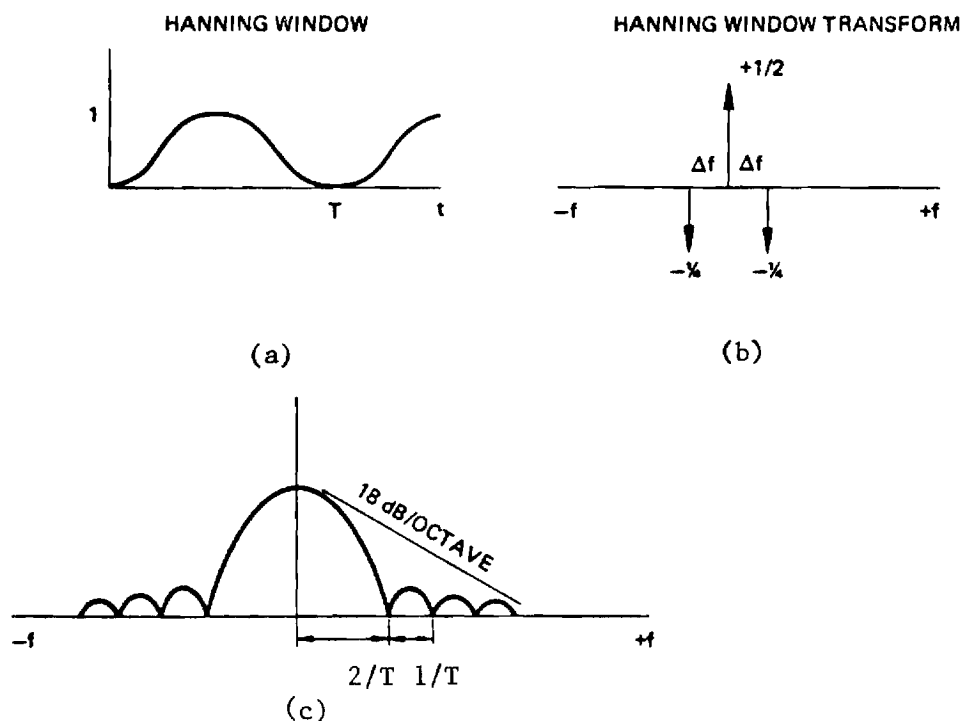


Figure 2.4-7. Hanning Window (a), its Fourier Transform (b), and its Convolution With The Boxcar Function(c).

$$W_H(f) = 0.5\delta(0) - 0.25\delta(\Delta f)$$

was used. The Hanning window is attractive due to its relatively narrow main lobe together with greatly reduced amplitude of the side lobes (78).

Digital Computation Requirements: The descriptive properties of a random process,  $x(t)$ , cannot be precisely determined from sample data. Therefore, suitably defined estimators have to be used in computing the spectral density, transfer, and coherence functions from finite samples of data. The raw estimates (denoted by superior carets) thus obtained require smoothing operations to reduce the statistical uncertainty in the estimates. The smoothed quantities are denoted by superior tildes.

The accuracy of parameter estimates based on sample values can be described by a mean square error defined as (60):

$$\text{Mean Square Error} = E[(\hat{\phi} - \phi)]^2 \quad (2.4-20)$$

where  $\hat{\phi}$  is an estimator of  $\phi$  and  $E$  is the expectation operator. The root mean square can also be expressed as

$$\text{rms error} = \sqrt{\sigma^2[\hat{\phi}] + b^2[\hat{\phi}]} \quad (2.4-21)$$

or, the normalized random error

$$\epsilon = \sqrt{\epsilon_r^2 + \epsilon_b^2} = \frac{\sqrt{\sigma^2[\hat{\phi}] + b^2[\hat{\phi}]}}{\phi}$$

where  $b\hat{\phi}$  is the bias error in the estimate and  $\sigma[\hat{\phi}]$  is the standard deviation of the estimate and  $\epsilon_r = \frac{\sigma[\hat{\phi}]}{\phi}$  is called the normalized random error and is a measure of the variance of the estimate.

$\epsilon = \frac{b[\hat{\phi}]}{\phi}$ , on the other hand, is a measure of the bias of the



estimate and is called the normalized bias error. An estimator with a small random error is said to have high stability and that with a small bias error is said to reproduce the true function with high fidelity.

Estimates of power spectral density functions from digital data can be obtained by one of the following two methods (60):

1. Blackman - Tukey Method: The power spectral density function is computed by first estimating the autocorrelation function and then Fourier transforming it, i. e., by using the Wiener-Khintchine relations.
2. Cooley - Tukey Method: This method is based on computing the discrete finite Fourier transform, via the FFT, of the original data.

Consideration of the computational speed advantage of the direct FFT make the second method more attractive; in fact, the autocorrelation function can often be most easily computed using the inverse FFT and the Blackman-Tukey method in reverse. The Cooley-Tukey method was used in the present study to compute the various spectra. In this method, raw estimates of the spectral density functions are computed as follows:

$$\hat{G}_{xx}(k\Delta f) = \frac{2}{T} |X(k\Delta f)|^2 = \frac{2}{T} X(k\Delta f) \cdot X^*(k\Delta f) \quad (2.4-22)$$

$$\hat{G}_{xy}(k\Delta f) = \frac{2}{T} X^*(k\Delta f) Y(k\Delta f)$$

where  $X(k\Delta f)$  and  $Y(k\Delta f)$  are discrete Fourier transforms of stationary, Gaussian, zero-mean random processes,  $\{x(t)\}$  and  $\{y(t)\}$ , respectively. The index,  $k$ , is an integer and  $*$  denotes the complex conjugate.

The estimator,  $\hat{G}$ , of Eq. (2.4-22) can be shown to be a  $\chi^2_2$  random variable and has a substantial random error (60). In practice the

random error of an estimate produced by Eq. (2.4-22) is reduced by smoothing over an ensemble (averaging) or by smoothing over frequency. The averaging technique was used in the present application and consists of dividing a record of length,  $T_S$ , into  $n_d$  disjoint segments such that  $T_S = n_d T$  and  $T = N \Delta t$ , where  $N$  is the number of values in each data segment. Individual estimates are computed for each of the  $n_d$  disjoint segments and the smoothed estimate,  $\tilde{G}$ , obtained by averaging these. The normalized (with respect to the expected value) random error for smoothed estimates is given by

$$\epsilon_r[\tilde{G}(f)] = \sqrt{\frac{1}{n_d}} = \sqrt{\frac{1}{\Delta f \cdot T_S}} \quad (2.4-23)$$

The bias error,  $b\{G(f)\}$ , is proportional to  $\Delta f^2 = \left(\frac{1}{T}\right)^2$ . Hence,  $T$  should be made as large as possible to minimize bias, but at the same time,  $n_d = \frac{T_S}{T}$  should be large to reduce the random error. Thus, reduction of bias and random errors imposes conflicting requirements on the data analysis parameters.

The overall level or mean square value estimates can be readily obtained from the power spectra by using Parseval's theorem which can be stated as follows:

Let  $\Psi_x^2 = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T x^2(t) dt$  be the true mean square value of  $x(t)$ , then  $\Psi_x^2$  can also be expressed as:

$$\Psi_x^2 = \int_0^\infty G_{xx}(f) df$$

The implication of Parseval's theorem is that the mean square value of a signal is the area under its power spectrum. Alternatively, the area under the power spectrum in a frequency band,  $(f, f + \Delta f)$ , is

the mean square value of the signal in that frequency band. This provides the basis for analog computation of the power spectral density function by filtering and squaring.

The mean square value can be estimated from the spectral density as

$$\hat{\psi}_x^2 = 2 \sum_{k=1}^{N/2} \tilde{G}_{xx}(k\Delta f) \quad (2.4-24)$$

The errors in the estimate,  $\hat{\psi}_x$ , depend on the errors in the smoothed estimates,  $\tilde{G}_{xx}$ . It can be shown (69) that relatively large errors in power spectrum estimates will have relatively small errors for  $\hat{\psi}_x^2$  when  $N$  is sufficiently large.

Coherence function estimates can be obtained from two sample time history records of  $\{s(t)\}$  and  $\{y(t)\}$ , assuming the sample records to exist over a common time interval  $T$ , as:

$$\hat{\gamma}_{xy}^2(f) = \frac{|\tilde{G}_{xy}(f)|^2}{\tilde{G}_{xx}(f)\tilde{G}_{yy}(f)} \quad (2.4-25)$$

The statistical accuracy of the estimate  $\hat{\gamma}_{xy}^2(f)$  depends on the estimation procedure of Eq. (2.4-25) as well as the accuracy of the power spectral density estimates. Empirical studies of Ref. 79 show that estimates of the coherence function in the range of  $0.34 \leq \gamma_{xy}^2 \leq 0.95$  based upon spectral density estimates with  $n_d \geq 20$  can be evaluated in terms of the transformation

$$\hat{w}(f) = \tan^{-1} \hat{\gamma}_{xy}(f)$$

In terms of  $n_d$ , the number of disjoint independent segments, the bias and variance errors are given by

$$b[\hat{w}(f)] \approx \frac{1}{2n_d - 2}$$

$$\sigma[\hat{w}(f)] \approx \frac{1}{2n_d - 2}$$

The frequency response function estimate for a single input-single output CPLS with transfer function  $H(f)$  input  $x(t)$  and output  $y(t)$  is obtained as

$$\hat{H}(f) = \frac{\hat{G}_{xy}(f)}{\hat{G}_{xx}(f)} = |\hat{H}(f)| e^{-j\hat{\phi}(f)} \quad (2.4-26)$$

The bias and random errors in this estimate are shown in Ref. 60 to depend on the number of degrees of freedom,  $n = 2n_d$ , of the spectral density function estimates and the coherence function estimate,  $\hat{\gamma}_{xy}^2(f)$ . The bias error is generally less than the random error for any given combination of  $n$  and  $\hat{\gamma}_{xy}^2(f)$ . The random error in  $\hat{H}(f)$  approaches zero as either  $n \rightarrow \infty$  or  $\hat{\gamma}_{xy}^2(f) \rightarrow 1$ . In interpreting practical measurements, if  $\hat{\gamma}_{xy}^2(f) \approx \gamma_{xy}^2(f) = 0$ , it follows that no physical system exists between the measurement points and no frequency response function can be defined. On the other hand, if  $\gamma_{xy}^2(f)$  has a nonzero value, no matter how small, then a frequency response function can be estimated with any desired degree of accuracy given sufficient data to make  $n$  appropriately large. Thus,  $\hat{\gamma}_{xy}^2(f)$  is an indicator of the validity of the frequency response function.

The Digital Processor: For computation of power spectra, transfer and coherence functions and amplitude histograms by digital techniques, a minicomputer based Fourier Analyzer (HP5451B) was used. The main features of the Fourier Analyzer are a 2-channel, 12-bit analog-to-digital converter which accepts analog data to 25 kHz, and a dedicated

minicomputer (32k core). Time-domain and frequency domain computations are performed in accordance with instructions from a keyboard. Simultaneous two-channel operations are possible and the results can be recorded on a x-y plotter or punched on paper tape. "Quick Look" data inspection is possible on a CRT display, an integral part of the unit.

#### 2.4(d) Experimental Techniques

The experimental plan called for the measurement, at various locations throughout the building, of cladding loads and the response of the two glass plates and the mullion to these loads. In addition, the reference wind velocity and direction were also required. The inaccessibility of the outdoor light from within the building required that a method for measuring its response without direct contact be devised, and a novel transducer was designed specifically for this purpose. Conventional transducers with suitable performance specifications were chosen for the measurement of the other variables. The instrumentation and its integration into a portable system is described below.

Wind Loading: The pressure loading on a window was assumed at the outset to be uniform over the area. The assumption, although at first seemingly crude, is realistic since the window area is relatively small compared to the building facade and, therefore, sudden spatial variations of pressure over this area are not expected. With this assumption, the wind loading on a window can be measured from the difference in the external and internal pressure provided the internal pressure is known.

Internal pressure variation due to the normal building activity, such as opening and closing of doors, operation of elevators and mechanical services of the building does not permit its use as a reference pressure. The ideal reference pressure would give a zero output for a transducer when the wind speed was zero. In comparing full-scale measurements with wind tunnel tests, various methods of referencing the wind pressure have been proposed, including most recently, the use of absolute pressure transducers (80). However, since the cladding loads are simply the differential pressure between the inside and the outside, no reference pressure is required in the present measurements. The differential pressure transducers used were variable capacitance Datametrics Type 501-11 Barocels with an internal volume of 0.15 cu. in. Datametrics Type 1015 signal conditioners were used with the Barocels. The auto-ranging feature of the Type 1015 was used to ensure that the pressure signals were not clipped in strong squalls, and was useful in providing adequate dynamic range.

With the uniform pressure assumption, a small number of point pressure taps can provide a measure of the loads on a window. Pressure taps through the mullion were planned initially but the channel-shaped mullion exterior appeared in early tests to cause upflow interference effects. The pressure taps were ultimately located on the muntion (horizontal top support) of the window, flush with the exterior surface. One side of the Barocel was connected to the tap and the other was vented inside. The external pressure tap and its location are shown in Fig. 2.4-8.

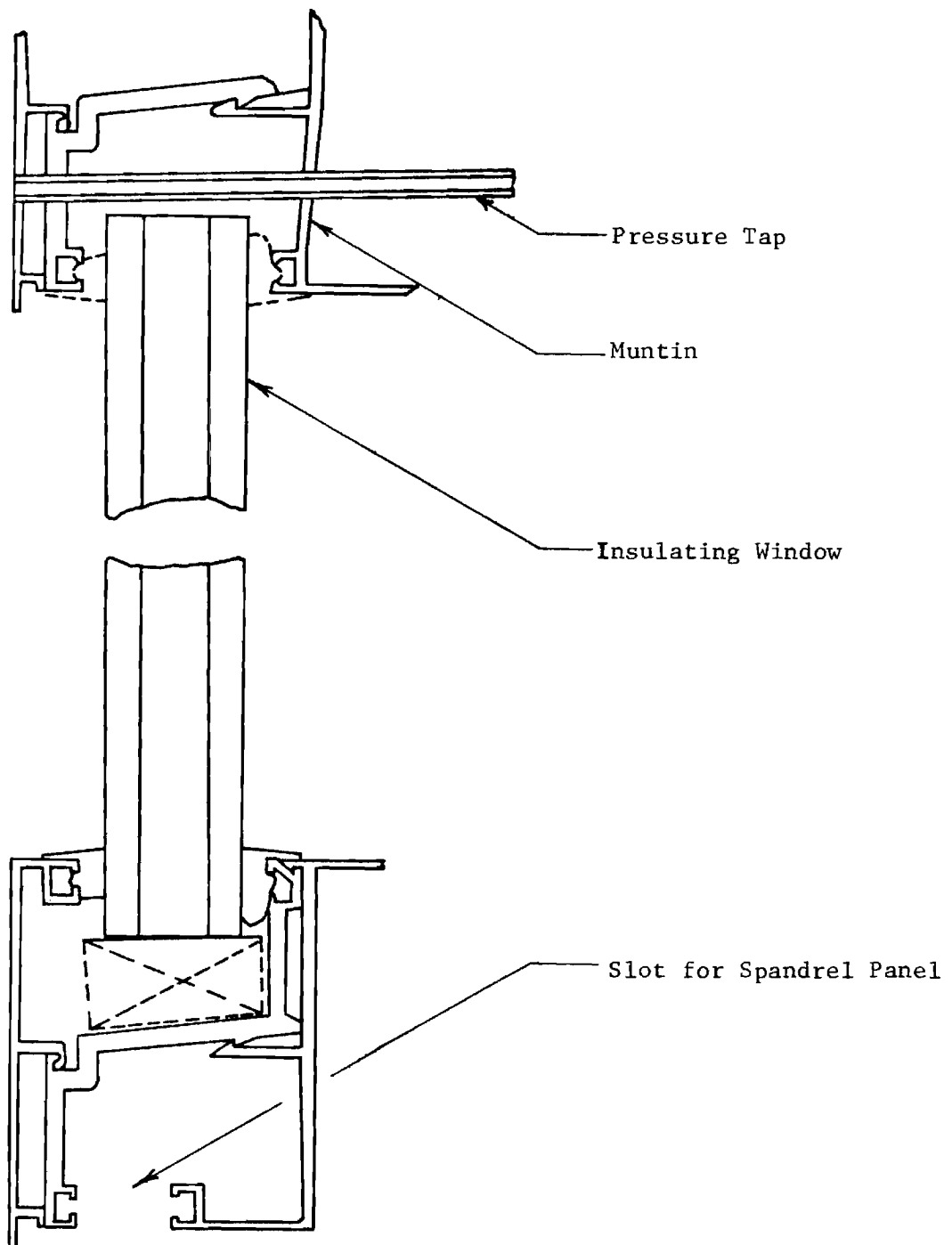


Figure 2.4-8. Pressure Tap and its Location.

The use of the sensor in dynamic pressure measurements required that the length of tubing used for the external pressure tap be selected to ensure adequate dynamic response. A suitable length was chosen after considering the resonance effects and damping in the system, and hardware limitations. Details of the response estimate are given in Ref. 40. A tubing length of 18 in. gave a flat response bandwidth up to 38 Hz with phase linearity up to 22 Hz. The bandwidth was considered suitable for pressure measurements (auto spectra) and the range of phase linearity adequate to span at least two structural response modes.

Reference Wind Velocity and Direction: The reference wind speed and direction were measured with an MRI Model 1022 Wind Sensor. The speed sensor has a range of 0.5 - 80 mph. To obtain reasonably undisturbed measurements of wind speed and direction, the anemometer was mounted on an "open" antenna tower at a height of approximately 455 ft above the ground level. Due to the existence of an atmospheric boundary layer compounded by interference effects due to the building, the velocity measured is not truly that of the mean wind. If the building interference effects are ignored, however, velocities at any other height can be obtained by using Eq. (2.3-2) applicable to the terrain conditions in the vicinity of the building. The wind speed and direction were recorded on strip chart recorders and the records used to compute the mean wind speed and direction over a certain length of time.

Response Measurement: The primary response parameters for the window glass were single point transverse surface displacements,  $w_i$



and  $w_e$ , of the inner and outer panes, respectively, and mullion acceleration,  $a_m$ .

The inner glass displacement,  $w_i$ , was measured by a linear variable differential transformer (LVDT) mounted in a fixture referenced to the building structure. The LVDT core was flexure-mounted so as to enable it to accurately follow the dynamic glass displacement. The kHz carrier frequency of the LVDT permitted adequate dynamic response for the present application.

The mullion acceleration was measured using an MB accelerometer. The piezoelectric accelerometer could not be used for frequencies below 5 Hz; however, for transmissibility measurements in the vicinity of resonance ( $\sim 10$  Hz, 1st mode) the dynamic response posed no serious limitations.

Measurement of  $w_e$  for a double glazed window is complicated by the inaccessibility of the outer pane to measurement by conventional transducers located inside the building. Therefore, a method for measuring  $w_e$  without direct contact was required. Some common contactless transducer concepts are based on capacitive, inductive or optical effects. Capacitive and inductive methods generally would require treatment of the outer glass. Optical techniques, on the other hand, were judged most promising and a variety of available optical displacement transducers were evaluated with respect to their performance and portability. Most of these required relatively complex processing of the measured signals, while others were found to be unsuitable because of a limited range or were affected by the presence of background illumination. As a result, a unique optical

transducer was designed which provided simplicity of operation, sufficient range and accuracy, and a capability of integration into a portable measurement system.

The sensor operates on geometric optics and utilizes a He-Ne laser beam for a collimated light source. The geometric movement of a beam reflected off the outer glass pane is sensed by an array of linear semi-conductor detectors. The detector output after conditioning by a set of preamplifiers and analog computing circuits produces voltages proportional to the motion and intensity of the beam. The design of the geometrical configuration of the laser-detector system and other details such as elimination of the effect of rotation in the vertical plane, and special signal conditioning electronics are given in Ref. 104. The transducer is capable of measuring both transverse displacements and changes in surface slope of the window. For the particular configuration used in the measurement of  $w_e$ , the transverse displacement sensitivity was 6.25 V/in. The usable range of displacements was 0.35 in. before effects of detector differential nonlinearity and violation of small angle assumptions caused the output to deviate from the linear relation.

Physical construction of the transducer is best described by Fig. 2.4-9 where the laser and the detector enclosure box are shown mounted on a 4" x 4" steel beam supported by a torsionally rigid tripod. The laser displacement meter (LDM) support also served as the fixture for the LVDT which is partially visible. The Barocel sensor seen clamped to the beam was located on the top horizontal support during an actual run. Retractable wheels on the tripod

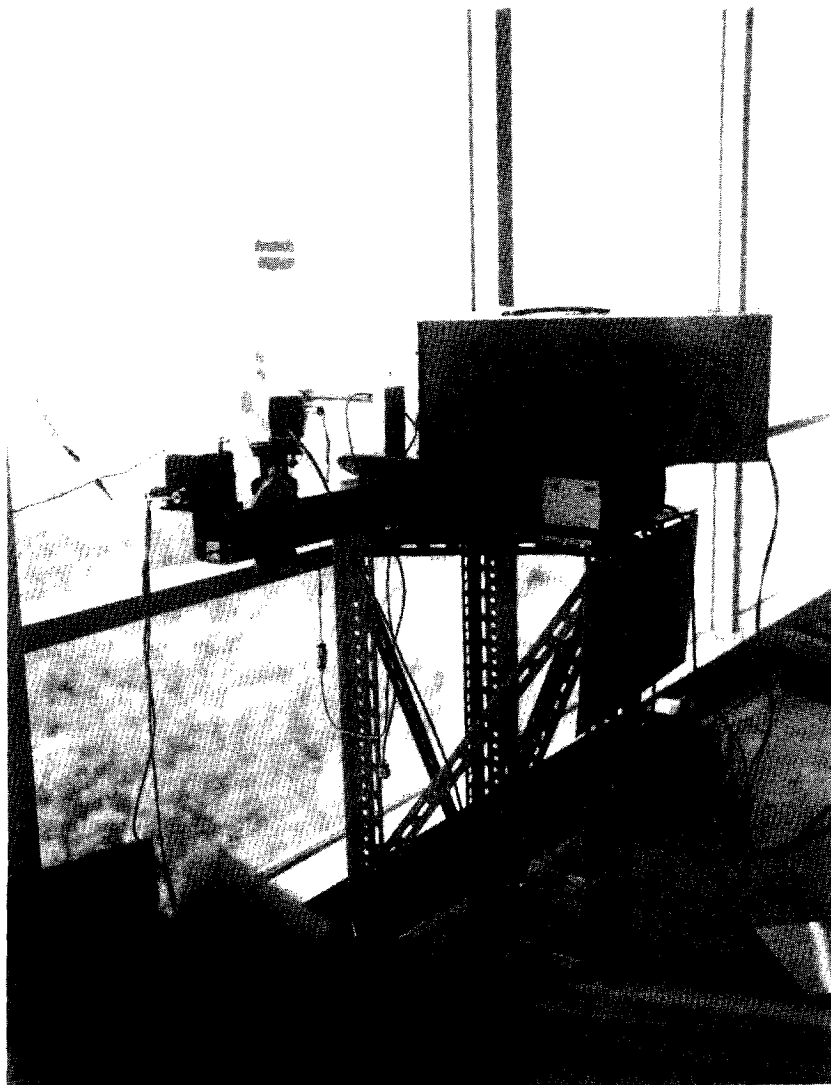


Figure 2.4-9 The Experimental Set-up.

permitted easy transport of the LDM/LVDT combination to various locations throughout the building. The figure also shows the displacement transducers positioned (by means of aluminum "spacer" rods to ensure calibration) opposite a window to record displacement histories.

The LVDT and LDM were positioned to measure the transverse displacements at corresponding points on the two panes. The measurement point was typically at the centroid of the horizontal axis and approximately 5 in. below the centroid of the vertical axis. This arrangement is suitable for detecting motion in any symmetric modes of the window glass. To measure the displacement response of the plates in antisymmetric modes a position closer to the supports is desirable and could be used by repositioning the platform. A schematic of the entire instrumentation is shown in Fig. 2.4-10.

Data Recording: The large quantities of data inherent in full-scale investigations require that careful consideration be given to a data acquisition system. Portability, unattended operation for several hours and compatibility of the system with the transducers is essential. While the use of a total digital recording system would have facilitated digital data analysis, the relatively high cost and lack of portability of such a system led to the choice of analog recording. For the present experiments the analog signals from all the transducers except the remotely located anemometer were recorded on a portable FM tape recorder (HP 3968A). Since the planned data analyses required only a mean wind speed and direction for a set of sample records from the response and load transducers,

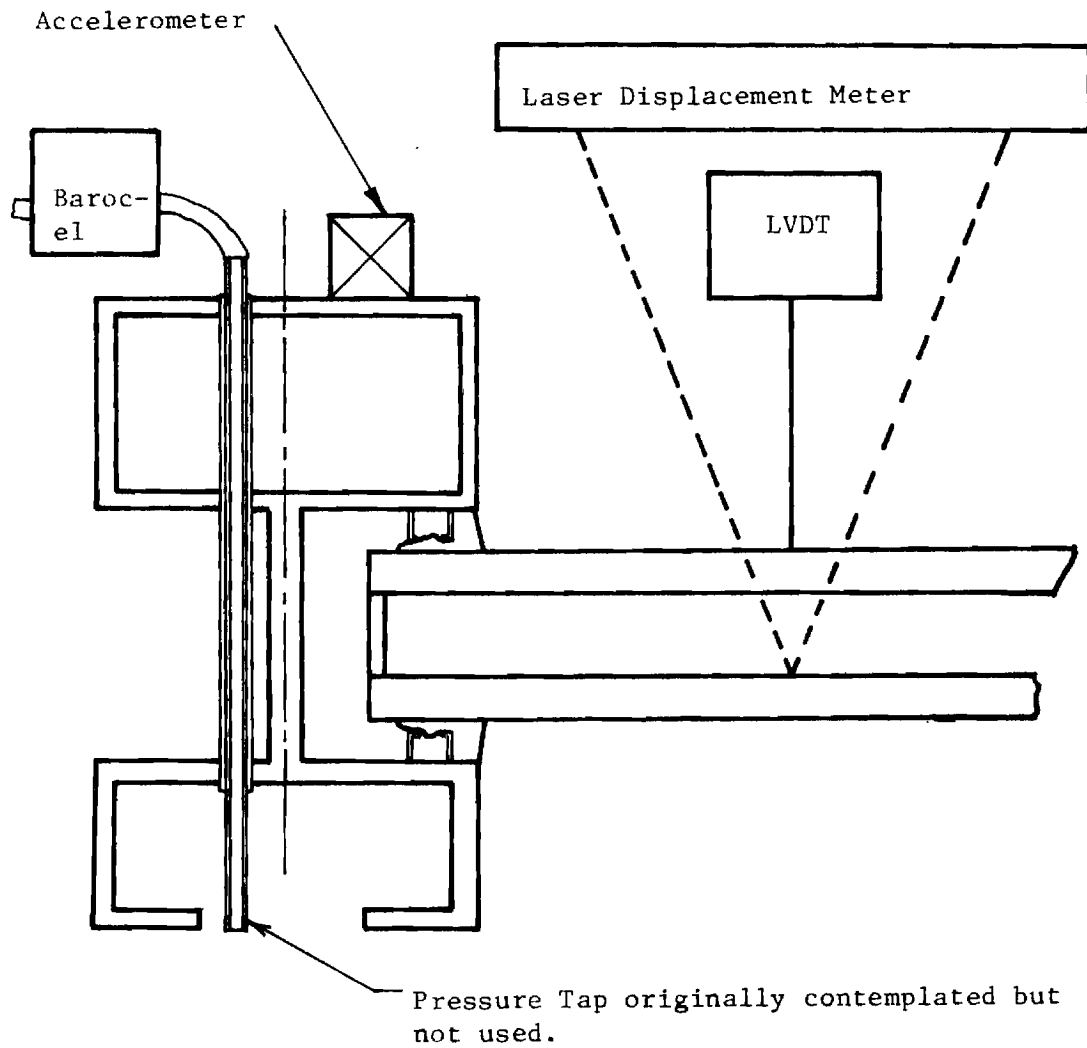


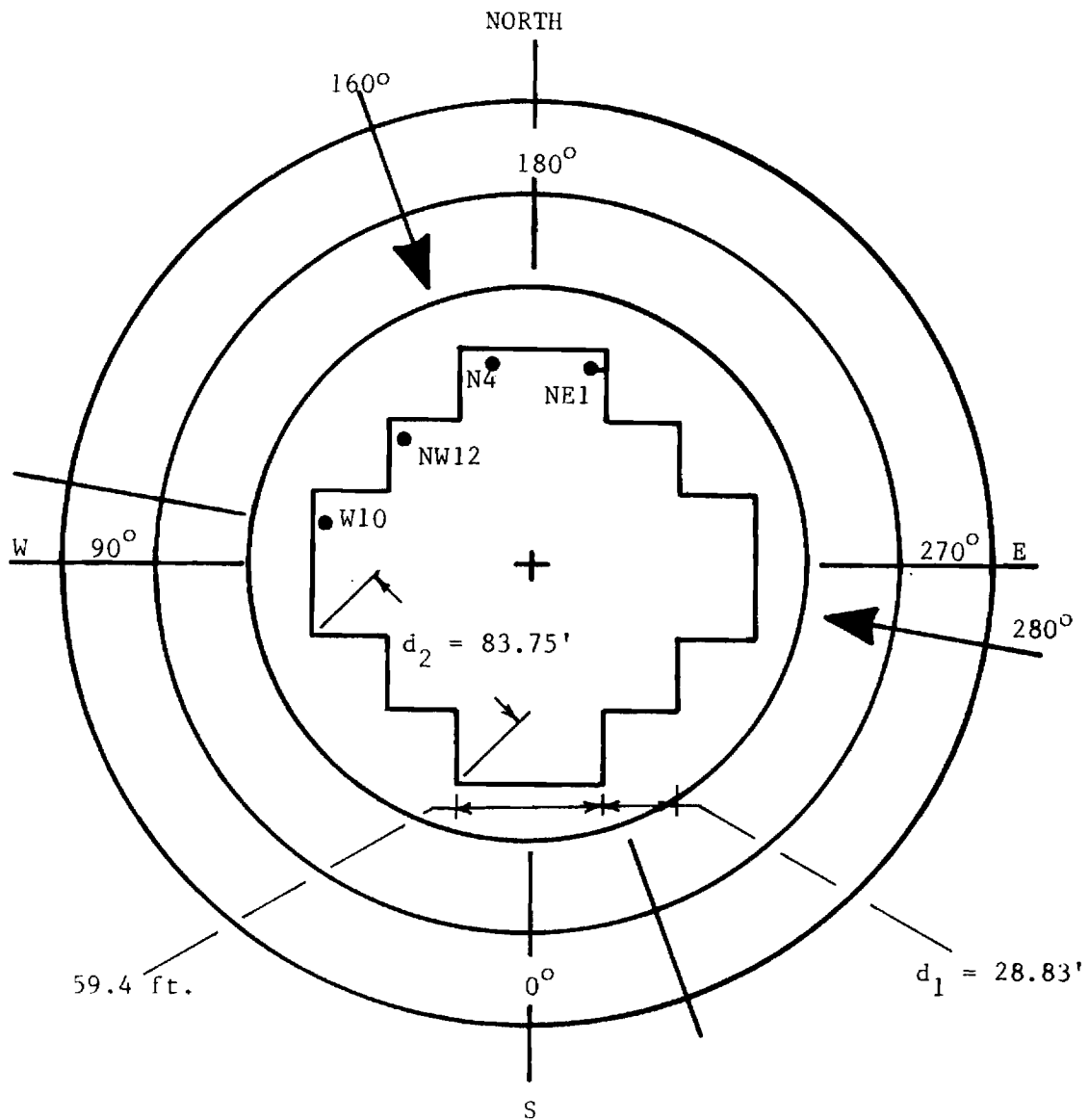
Figure 2.4-10. Schematic of the Instrumentation.

the use of strip chart recorders for the anemometers posed no serious problems. Data recording was continuous for a particular session (up to 16 hours) and the records were sorted for relatively large pressure fluctuations prior to analysis.

Test Methodology: Weather forecasts were used to select days on which the experiments were run. The test sites for the measurements were selected on the basis of available wind tunnel pressure distribution data from an earlier simulation (58). Windows were selected in areas where large negative pressure coefficients occurred in the simulation with the prevailing wind direction as the flow direction. In order to assess the variation of wind loads with height, tests were planned to include three different floors of the building, viz. 9th, 21st, and 28th. The window nomenclature used and the wind directions for which data were recorded are shown in Fig. 2.4-11.

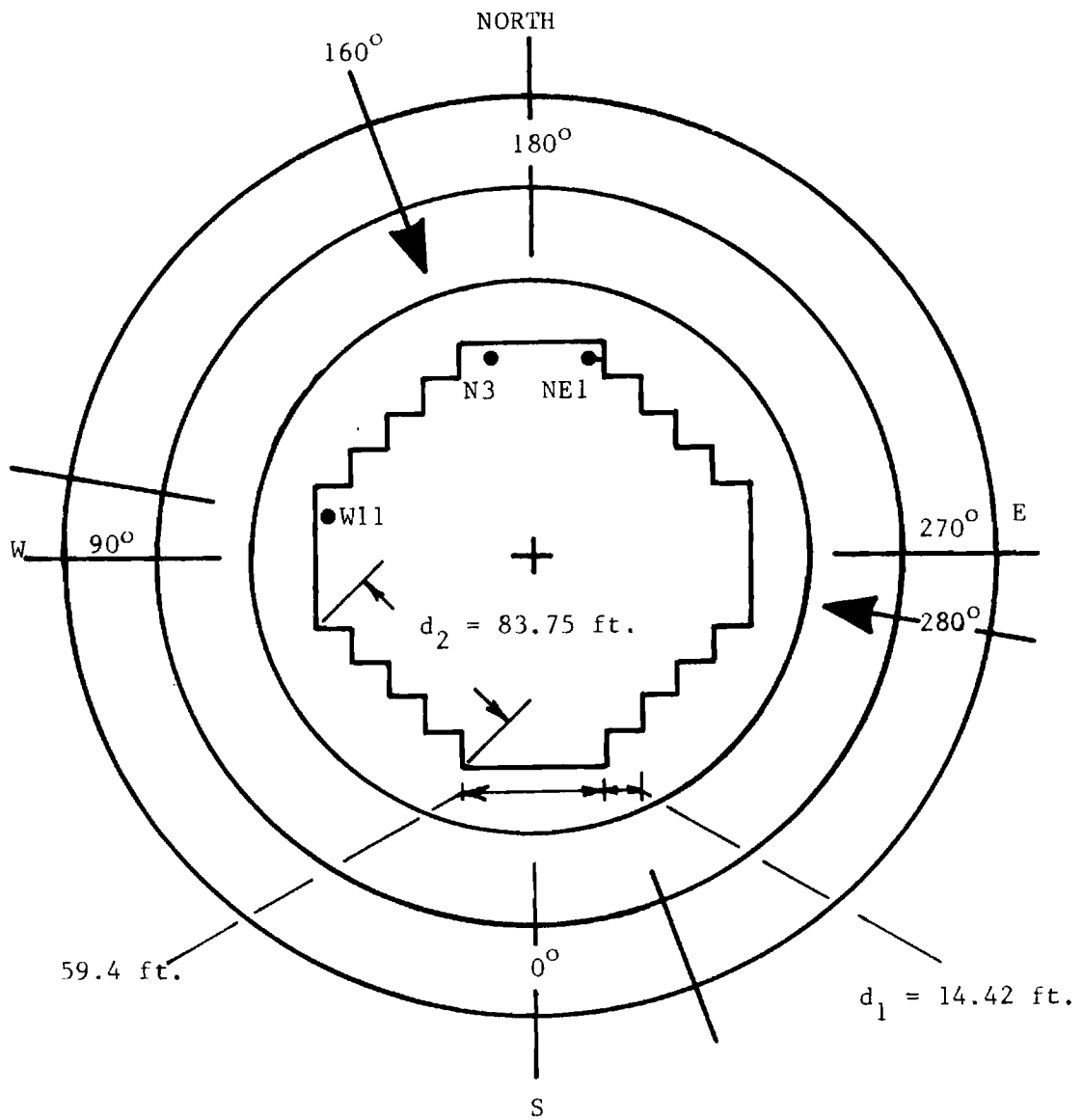
On days when reasonable wind velocities were forecast (about 15 - 20 mph) the portable instrumentation system was set up at a suitable window, depending on the wind direction, and the data recorded. For unexpected storm systems moving into the area data acquisition was not possible due to the time required to get to the test site and make the measurement system operational. A summary of the acquired data is shown in Table 2.4. Recording sessions were usually unattended and on the average were conducted over periods of approximately six hours.

The major problem encountered in the data acquisition process was one of dynamic range of the transducers. The selection of transducer signal conditioner range was dictated by the extremes of adequate



(a) 9<sup>th</sup> Floor

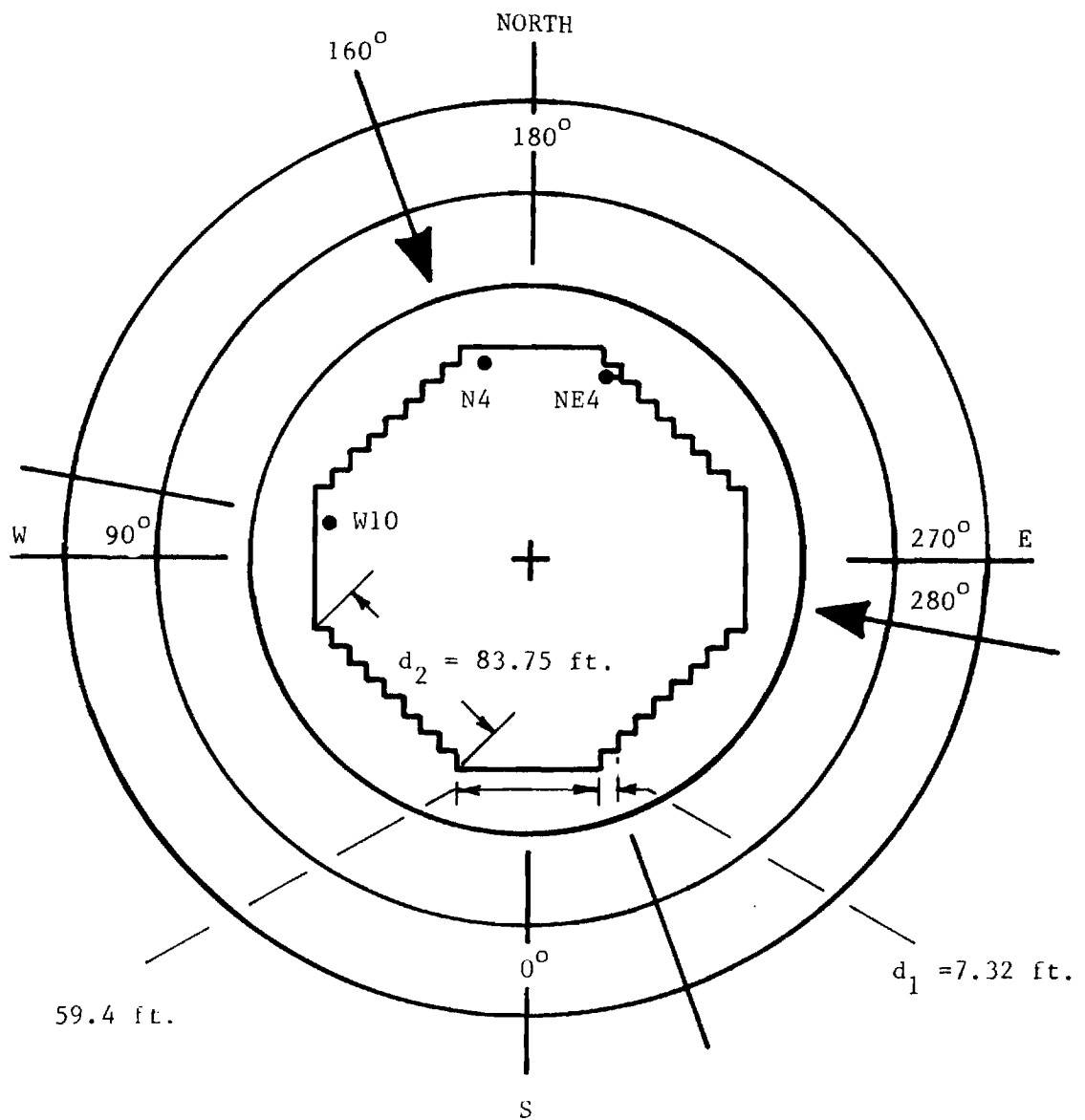
Figure 2.4-11. Window Nomenclature; Directions With Respect to Building Axis.



(b) 21<sup>st</sup> Floor

Figure 2.4-11 (contd.). Window Nomenclature; Directions With Respect To Building Axis.





(c) 28<sup>th</sup> Floor

Figure 2.4-11 (contd). Window Nomenclature; Directions With Respect To Building Axis.

Table 2.4 Summary of the Data Collected

#	Date (1976)	Duration Start/# hrs	Window	fps $\bar{v}$	$0^\circ$ S $\bar{\theta}$	Transducers Used				
						$\Delta p$	Wext	Wint	Am	$\Delta p$ null
1	6-2	3:52pm/3	F28/W10	23	108 $^\circ$	x	x	x		x
2	6-5	2:35pm/4	F28/NE3	23	290 $^\circ$	x	x	x		x
3	6-6	2:52pm/6	F28/N10	18	270 $^\circ$	x	x	x	x	
4	6-7	5:43pm/3	F28/N4	15	281 $^\circ$	x	x	x	x	
5	6-15	2:45pm/6	F21/NE1	12	324 $^\circ$	x	x	x	x	
6	6-16	3:50pm/8	F21/W11	15	98 $^\circ$	x	x	x	x	
7	6-19	3:58pm/4	F21/NW19	12	72 $^\circ$	x	x	x	x	
8	7-16	12:30pm/5	F21/N3		194 $^\circ$	x	x	x		
9	7-16	8:09pm/6	F21/N3	15	108 $^\circ$	x	x	x		
10	8-2	2:55pm/3	F21/N3			x	x	x		
11	8-3	2:29pm/4	F21/NE1			x	x	x		
12	8-18	4:05pm/4	F21/N3,NE1	23	280 $^\circ$	x	x	x		
13	8-19	1:42pm/5	F9/N4,NE1,NW12	30	270 $^\circ$	x				
14	8-20	2:13pm/8	F9/N10,N4,NE1	15	270 $^\circ$	x	x	x		
15	9-10	1:38pm/5	F9/W10,N4,NE1	30	160 $^\circ$	x	x	x		
16	9-17	5:42pm/3	F21/NE1,N3,W10	18	160 $^\circ$	x				
17	9-21	6:20pm/8	F28/W10,N4,NE4	23	160 $^\circ$	x				
18	10-21	10:30pm/8	F9/N4,NW12,NW4	30	160 $^\circ$	x	x	x	x	

signal-to-noise ratio and avoidance of clipping of the high amplitudes. The range selected in the first hour or so of a run, therefore, in some cases was inadequate for the strong gusts that occurred later in the recording period. To overcome this problem, an autoranging signal conditioner that automatically switched to a higher range when a preset threshold was crossed (accompanied by a switching pulse to indicate the change) was used. In a typical run the starting range and that at the end of the run were recorded. The switching transients in the data were then used to associate the correct scales with the recorded segments.

## 2.5 Experimental Results

Having defined the site and established the ground work for the data evaluation, it is now possible to present the results of the field study. In the sections that follow, the characteristic features of the localized pressure loads and their relationship to the overall structural geometry will be presented, the causality between the pressure loads and the window assemblies themselves will be determined. In making comparisons with expected behavior, reference will be made to the theoretical descriptions in Chapter 3. Two basic themes underlie the results that follow. The first is, of course, the problem of parametrically analyzing the pressure loads and defining critical combinations of independent variables that could give rise to extreme dynamic loads at a particular site. The second is the uniqueness of applying the particular evaluation methodology to the present problem. While most of the procedures have been previously developed in the area of information theory, their application to problems in mechanics

is quite new.

## 2.5(a) Parametric Analysis of Wind Loads

In Section 2.4(b) functional forms for the wind load characteristics, namely the reduced power spectral density, the mean, and the rms pressure coefficients, were hypothesized. The differential pressure measurements were conducted to examine the effects of changes in the independent variables. It can be anticipated that determination of the exact functional forms of Eqs. (2.4-16) and (2.4-17) by field tests on an actual structure is a practical impossibility. Rather, in the present case, data from sites preselected on the basis on wind tunnel tests and from the viewpoint of susceptibility to dynamic loads are analyzed. The objective is to assess the severity of dynamic loading and describe in a qualitative fashion the dependence of the load properties at these sites on the various profile and flow parameters.

Before proceeding with the analyses of the data certain preprocessing operations were required. The large volume of data (close to 300 hours, real time) necessitated selection of segments from the analog records where, for example, significant pressure fluctuations were observed. Error considerations in the estimation of the spectra required that the spectral analysis parameters be suitably chosen. In the following sections, the data sorting procedures used and design of the spectral analysis are described.

Data Sorting: The differential pressure and the response in a particular session were recorded over spans of up to 12 hours. During these intervals, periods of calm where the wind speeds are so low as to produce insignificant cladding loads were observed in the

recorded data, as expected, and unless the winds were due to a transient weather system, the periods of calm could be of long duration. Therefore, selection of particular segments of pressure and response signals of interest was required. This sorting process could, of course, have been obviated if the recorder were actuated when, for example, the reference wind velocity exceeded certain threshold. There are, however, several drawbacks to this mode of data acquisition that ruled it out for the present study:

1. The recorder takes time (approximately 10 secs.) to reach the correct recording speed, during which a significant part of the data may be lost.

2. The triggering process may introduce nonstationarity in the mean square value.

3. Due to the 3-dimensional nature of the flow around the building, the wind speed measured at the remotely located anemometer may not indicate the local activity.

Instead, weather forecasts were used to select periods and the recording equipment was operated continuously during that time. The data segments to be analyzed were selected on the basis of the wind speed records and visual inspection of the pressure data. Portions of the analog records with relatively large pressure fluctuations were isolated, and the corresponding wind speed and direction records were then inspected in order to choose those segments where  $\bar{v}$  and  $\theta$  showed a fairly constant mean value over a period of about 15 - 30 min. Records from a given session with  $\bar{v}$  greater than 20 fps were analyzed.

In order to assess the dependence of the wind loads on the independent variables of Eqs. (2.4-16) and (2.4-17), the analyses were carried out in accordance with the following:

1. Variation of wind loading with height  $z$ . Corresponding window locations were chosen on each of the three floors, 9th, 21st, and 28th. For each of these locations, records corresponding to a constant wind direction  $\bar{\theta} = \bar{\theta}_0 \pm 10^\circ$  were processed to obtain the PSD and histograms.

2. Dependence of the wind loading on the wind direction  $\bar{\theta}$ . The variables  $\phi_z$  and  $z$  were kept constant by selecting pressure records, for a particular window on a given floor, corresponding to wind directions differing by at least  $30^\circ$ .

3. Variation of wind loading along the building profile. To assess this effect, the variables  $\bar{\theta}$  and  $z$  had to be kept constant and the local pressure characteristics examined as  $\phi_z$  was varied. Thus for a preselected wind direction (accurate to within  $10^\circ$ ) the pressure spectra and histograms for a number of windows on a given floor were obtained. The variation of wind direction from the selected mean for each of these windows was less than  $20^\circ$  with the difference due to the data being recorded in different sessions.

In all, about 96 hours of usable data were processed. The grouping of the sorted data is summarized in Table 2.3.

Spectral Analysis: The formal mathematical basis for power spectrum estimation techniques can be found in the more general topic of estimation theory. However, estimation of power spectra with high stability and fidelity as it is currently practiced has a very

Table 2.3 Data Sorting on the Basis of Wind Direction

Window Location	$\bar{\theta}$	$\bar{v}$ fps	Record #	$\bar{\theta}$	$\bar{v}$ fps	Record #	$\bar{\theta}$	$\bar{v}$	Record #	$\bar{\theta}$	$\bar{v}$	Record #
F28 W10	160°	24	17				108°	24	1	72°	24	1
N4	160°	24	17	280°	15	4						
NE3	160°	24	17									
NE3				280°	24	2						
N10				280°	18	3						
NW4												
F21 W11	160°	18	16				81°	15	6			
N3	160°	18	16	280°	24	12	119°	21	8			
NE1	160°	18	16	280°	24	12	324°	12	5			
NW2												
NW19												
F9 W10	160°	30	15									
N4	160°	30	15	270°	30	13						
NE1	160°	30	15	270°	30	13						
NW4	160°	30	18									
NW12	160°	30	18	270°	30	13						

strong empirical basis. A general agreement on the best combination of analysis parameters is rare prior to data collection. After data acquisition, however, the analysis parameters can be properly chosen by pilot analyses.

Before conducting a pilot spectral analysis of the raw data certain preprocessing decisions must be made that involve:

1. Sampling considerations
2. Trend removal
3. Stationarity check

Sampling considerations require the maximum frequency,  $F_{\max}$ , present in the data be known and subsequent digitization be in accordance with Shanon's sampling theorem. The maximum frequency in the pressure data was limited to 50 Hz by an analog low-pass filter with a roll-off rate of 48 dB/octave, to minimize aliasing errors; the actual frequency range of interest was about 40 Hz since it was sufficient to include the first few response modes.

A trend in the data is defined as any frequency component whose period is longer than the record length. This means that the lowest frequency in the data should be greater than the frequency resolution,  $\Delta f$ . A low frequency trend is illustrated by curve (a) of Fig. 2.5-1. The trend is generally removed using linear or polynomial regression, but for the present pressure data, considering the physics of the phenomena (significant energy at low frequencies), high-pass analog filtering with a cut-off frequency of approximately  $2\Delta f$  was used instead. The advantages were threefold:



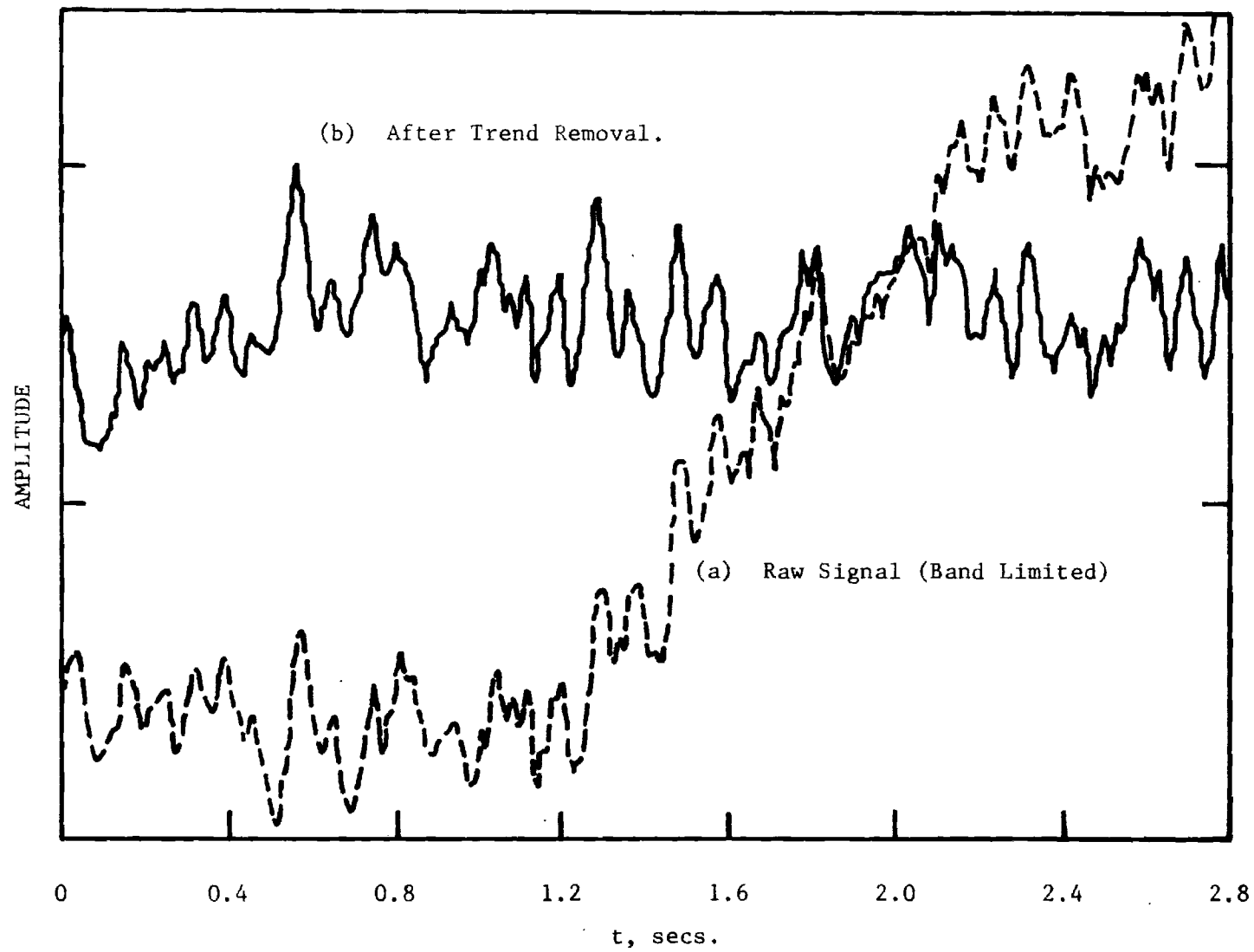


Figure 2.5-1. Trend Removal.

(a) Reduced quantization noise. To keep the quantization noise level low the signal must occupy as much of the input range of the A/D converter as possible (78). High-pass filtering eliminates the large low frequency excursions and permits higher A/D gain for digitization of the small fluctuations about the mean.

(b) Increased dynamic range. Dynamic range can be defined as the ratio of the highest and lowest spectral component that the processor can detect from the roundoff and quantization noise, and for the present processor is about 90 dB (double precision spectra, 32-bit word size). Since a high value of the dc component can often exceed this range, effectively drowning out small values at higher frequencies, high pass filtering thus provides increased dynamic range.

(c) Removal of the low frequency trend (curve (b) Fig. 2.5-1). Trend removal by high-pass filtering required that the pressure spectrum be analyzed in two parts: a high frequency and a low frequency band. The dividing frequency was chosen after analysis parameters had been finalized.

The wind pressures were not expected to be stationary over the entire recording period (~6 hours) due to diurnal changes and gust variations. However, stationarity over a short period (an hour or so) was checked for representative data in accordance with Ref. 60, and over these periods the data were found to be stationary.

Estimation of the power spectra from the stationary random data requires consideration of the following errors:

1. Frequency leakage
2. Bias errors in the estimate
3. Variance errors in the estimate

With reference to section 2.4, Eqs. (2.4-23) to (2.4-26), these can be minimized to acceptable levels by selection of:

- (a) a suitable time window,  $w(t)$
- (b) frequency resolution,  $\Delta f$ , or block size,  $N$
- (c) number of ensemble averages,  $n_d$

To a large extent, the choice depends on the nature of the data and reduces to a design problem. The single Hanning time window was chosen to reduce leakage for reasons outlined in section 2.4. Figure 2.5-2 illustrates the effect on the bias in the estimate of varying the block size,  $N$ , or equivalently,  $\Delta f$ . As  $N$  is increased the bias is reduced (the difference between spectra is diminished) at the expense of the variance.

The bias for  $N = 512$  and  $1024$  shows no significant change, however, since the computation time for  $N = 1024$  is substantially increased,  $N = 512$  was chosen. The variance can be reduced by increasing the number,  $n_d$ , of disjoint records that are averaged:

$$\frac{\sigma}{\mu} = \frac{1}{\sqrt{n_d}}$$

The mean power spectrum can thus be estimated within 10% of its true value with 95% confidence (Gaussian distribution) if:

$$0.01 = \frac{1.96\sigma}{\mu} = \frac{1.96}{\sqrt{n_d}}$$

which yields  $n_d \approx 384$ . The reduction in variance is apparent from Fig. 2.5-3 and is independent of  $N$  or  $\Delta f$ .

The analysis parameters for the high-band portion of the power spectrum can be summarized as follows:

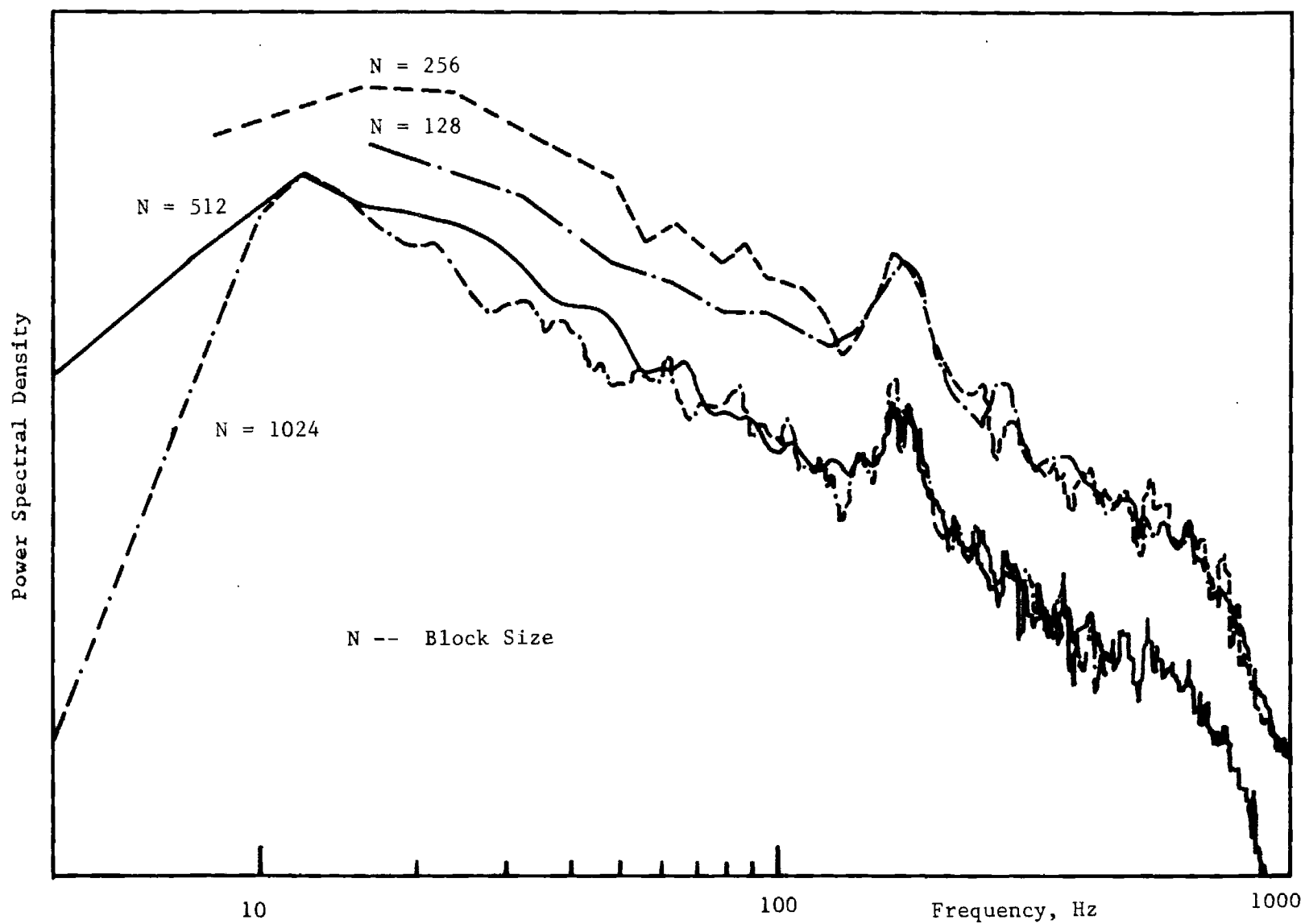


Figure 2.5-2. Effect of Block Size on Spectral Estimates.

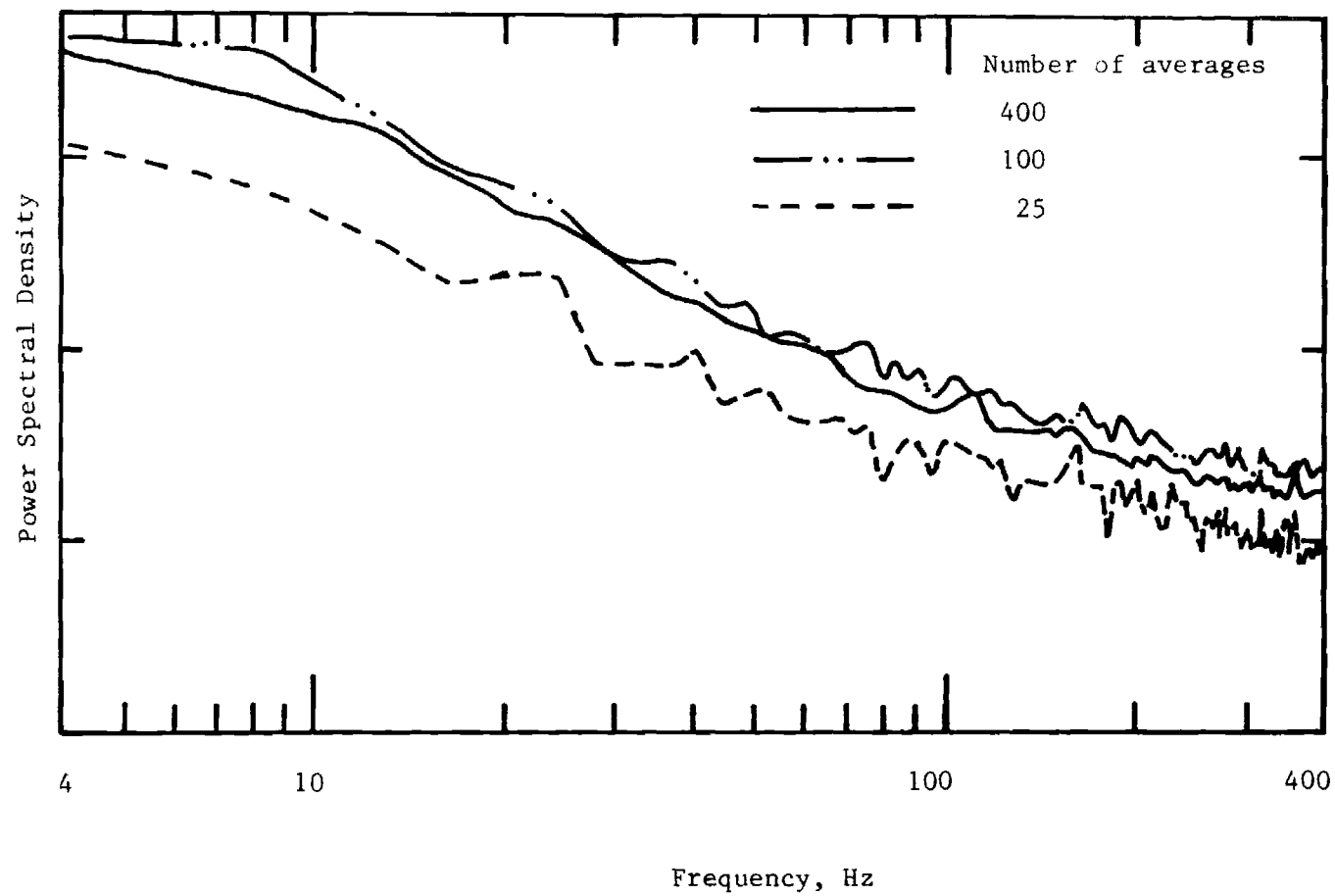


Figure 2.5-3 Effect of Averaging on Spectral Estimates.

1. Maximum Frequency = 50 Hz      Block Size = 512  
     Anti-Aliasing Filter = 3 Hz      High Pass Filter = 0.025 Hz  
                                          Number of Averages = 30

The lowest frequency that can be resolved above is about 0.75 Hz.

The parameters used for the low frequency part of the spectrum were:

2. Maximum frequency = 3.125 Hz      High Pass Filter = 0.025 Hz  
     Anti-Aliasing Filter = 3 Hz      Number of Averages = 30  
     Block size = 512

The lowest frequency that could be resolved was 0.025 Hz, corresponding to a period of 40 secs., which is sufficient to include short duration gusts.

Discussion of Results: In order to describe the wind pressure loading as a random process, the data were processed in accordance with the above specifications to obtain the following descriptors:

1. High frequency band power spectral densities (PSD) (0.8 - 40 Hz)
2. Low frequency band power spectral densities (PSD) (0.03 - 3 Hz)
3. Normalized pressure amplitude histograms

After sorting, the available data permit a qualitative assessment of the wind loading with respect to the following values of the independent variables:

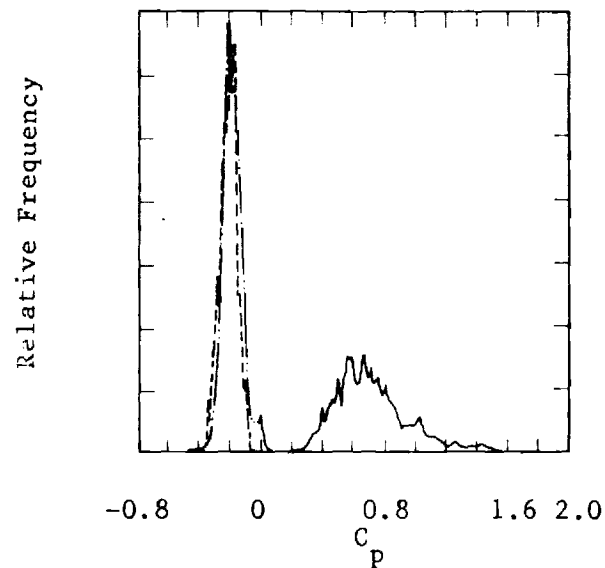
1. Wind Direction,  $\bar{\theta}$ :      160°, 280° ( $\pm 10^\circ$ )
2. Height, z:                      9th Floor (120 ft.)  
                                          21st Floor (270 ft.)  
                                          28th Floor (357.5 ft.)  
                                          (Total height: 376.15 ft.)
3. Window Locations,  $\phi$ :      N4, W10, NE1

Wind Incidence,  $\bar{\theta} = 160^\circ$ : Figures 2.5-4a, b, and c show the spectra and normalized histograms for window N4 with the height,  $z$ , as a parameter and  $\bar{\theta} = 160^\circ$ . The respective floor planforms and wind incidence for the three floors are shown in Fig. 2.4-11.

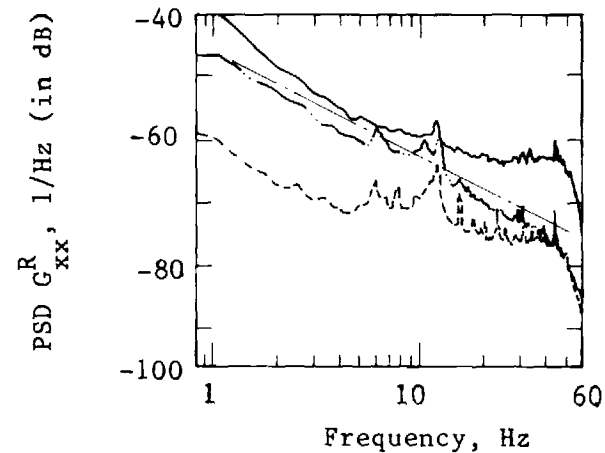
The window N4 is on the windward side, and therefore the turbulence characteristics of the distant flow are expected to be preserved. If the flow is visualized as 2-dimensional, location N4 is expected to be downstream of the stagnation point and in the reattached region (flow separation at the corner). The pressure fluctuations, therefore, must be due to the wind turbulence and possibly due to roughening caused by the mullions. It appears reasonable to consider the pressure characteristics at N4 as reference, and examine with respect to these, the pressure fluctuations at aerodynamically more complex locations such as W10 and NE1. It is recognized, however, that in reality the flow is 3-dimensional and pressure characteristics at the higher floors may be considerably altered due to upflow.

Magnitudes of the pressure coefficients are seen from Fig. 2.5-4a to be considerably higher at  $z = 9$  than at the upper floors. The positive pressure coefficients at the 9th floor and the mostly negative values at the 21st and 28th suggest the existence of an upflow. At the same time, discrete frequencies are seen in the low frequency band spectra as follows:

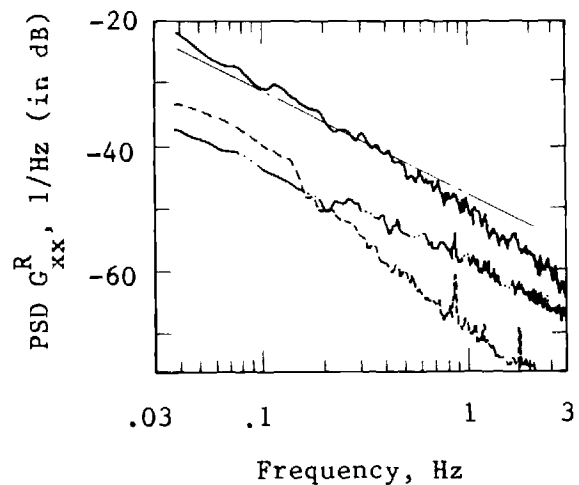
1. 28th	$f_1 = 0.86 \text{ Hz}$	$f_2 = 1.7 \text{ Hz}$
2. 21st	$f_1 = 0.86 \text{ Hz}$	
3. 9th	None	



(a) Normalized Histogram



(b) High Frequency Band PSD



(c) Low Frequency Band PSD

Notes:

————— 9<sup>th</sup> Floor  
 — · — · — 21<sup>st</sup> Floor  
 - - - - - 28<sup>th</sup> Floor

PSD in dB Relative to 1 1/Hz  
 e.g. -20 dB =  $10^{-2}$

Figure 2.5-4. Spectra and Normalized Histograms,  
 Window N4,  $\bar{\theta} = 160^\circ$



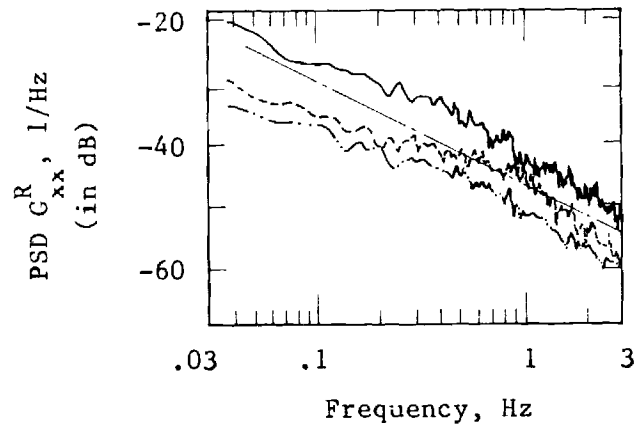
These frequencies could be due to vortex shedding from the reverse setbacks at the 19th and the 24th floors, Fig. 2.3-1. In the event that there is some upflow, the pressure fluctuations at  $z = 21$  are affected by a single change in cross-section, whereas at  $z = 28$  the flow has traversed two cross-sectional changes. The fact that  $f_1 = 0.86$  Hz is seen in the PSD for the 21st as well as the 28th and that  $f_2 = 2f_1$ , serves to reinforce the observation that there is an upward flow component (not necessarily vertical) and periodicities are induced by the setbacks. That  $f_1 = 0.86$  Hz is generated by the characteristic corner dimension,  $d_{19} = 14.42'$ , and  $f_2 = 1.7$  Hz is generated by the characteristic corner dimension,  $d_{24} = 7.32'$ , is evident from the Strouhal number,  $S = \frac{fd}{V} = 0.48$  in both cases. (The Reynolds number based on the corner dimension is of the order of  $10^6$ ).

The high frequency band PSD, Fig. 2.5-4b, reveal that the mean square value of the pressure fluctuations decreases with increasing height. The PSD levels in dB are with respect to  $G_{xx}^R = 1$  sec as defined in section 2.4. The histograms in Fig. 2.5-4a show, however, that the variance at the 28th and the 21st floors is approximately the same. This is because of the increased energy (higher PSD level) at the 28th floor for low frequencies. In Fig. 2.5-4b, a discrete frequency component is seen at approximately 12 Hz in all three spectra. The value corresponds to the fundamental frequency of the glass plates and appears in the spectrum due to feedback to the wind support. Association of this frequency with the flow characteristics, therefore, seems tenuous at best, but its appearance does point out the unusual sensitivity of the evaluation method. The narrow-band noise around

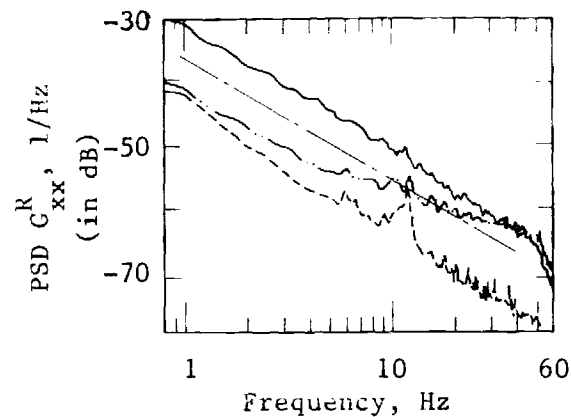
6 Hz for  $z = 21$  and  $28$  could be due to the external architectural features such as the protruding mullions affecting the flow. At  $z = 9$ , this component has lost prominence due to higher energy levels. The flattening of the 9th and 28th floor PSD's beyond about 20 Hz (Fig. 2.5-4b) is unusual as regards turbulence spectra, but was also observed in a number of other instances. The most plausible reason is that in these particular runs the combination of transducer range and tape recorder gain served to restrict the dynamic range and thus suppress the low amplitude, high frequency pressure fluctuations into the noise floor. The effect of the anti-aliasing filter is clearly seen as the sharp roll-off beyond 50 Hz.

Figures 2.5-5a, b, and c show the results for the aerodynamically more complex case of window W10 with  $\bar{\theta} = 160^\circ$ . The location is downstream of a varying number of corners depending on the height being considered.

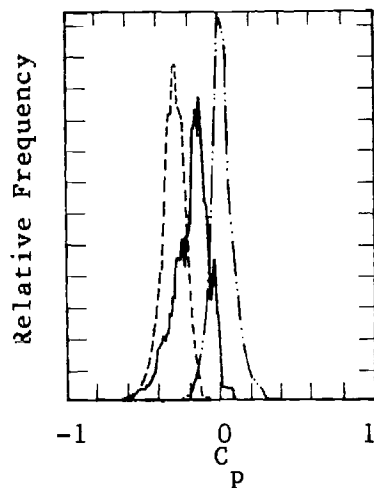
The magnitudes of the pressure coefficients are smaller than for N4 at  $z = 9$ . The pressure coefficients at  $z = 9$  and  $z = 28$  are predominantly negative, whereas at  $z = 21$  they are distributed such that  $\bar{C}_p \approx 0$  (actually  $\bar{C}_p = 0.008$ ). One explanation for this behavior is the roughening effect provided by the corners. The large corner dimensions and the small number of corner at  $z = 9$  may not dramatically affect the flow at location W10. Hence, the pressure coefficients at W10 are expected to be the same as they would be for flow around a right corner, thus accounting for negative  $C_p$ 's. However, negative  $C_p$ 's are also predominant at  $z = 28$ . In this case the much smaller corner dimensions and the large number of corners tend to smooth out



(c) Low Frequency Band PSD



(b) High Frequency Band PSD



(a) Normalized Histogram

Notes:

————— 9<sup>th</sup> Floor  
 - - - - - 21<sup>st</sup> Floor  
 - . . - . 28<sup>th</sup> Floor

PSD in dB Relative to 1, 1/Hz

e.g. -20 dB =  $10^{-2}$ .Figure 2.5-5 Spectra and Normalized Histograms, Window W10,  $\bar{\theta} = 160^\circ$

along the diagonal from face N to face W. Therefore, W10 again happens to be around a corner with respect to the flow, although the corner is not as sharp as at  $z = 9$ , and negative  $C_p$ 's can be expected. At the 21st floor, on the other hand, the corners may be expected to roughen up and thus randomize the flow. This results in negative and positive  $C_p$ 's occurring equally often. Considering the roughening effect in terms of the ratios  $r = \frac{d_1}{d_2}$  (See Fig. 2.4-11 for notation.) it is seen that  $r_9 = 0.34$  and  $r_{28} = 0.085$  provide the extremes, whereas  $r_{21} = 0.17$  tends to alleviate the pressure loads.

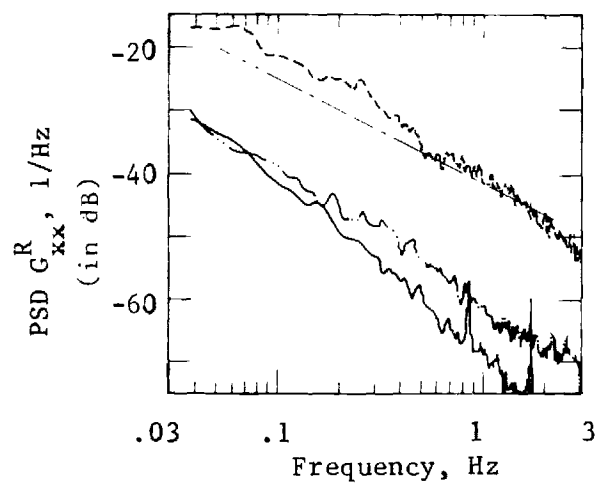
The high frequency band spectra are somewhat different from those at N4. The energy levels of the spectra at  $z = 9$  and  $z = 28$  are considerably higher for W10 (of the order of 10dB or by a factor of 3 on the rms values of the pressure) than on N4. However, the spectra for W10 at  $z = 9$  and 28 roll off much faster with frequency. The high frequency flattening in the PSD for  $z = 9$  and 28 is again attributed to the energy levels being in the noise floor for those recording sessions. The trends in PSD levels, however, are similar in that the levels decrease with height. At W10 for  $z = 21$ , the PSD levels, Fig. 2.5-5b, are almost the same as in Fig. 2.5-4b for  $z = 21$ , although the slopes differ somewhat. These observations suggest that higher levels of PSD, i.e., higher energies in a frequency band, occur in regions of negative pressure insofar as  $z = 9$  and  $z = 28$  are concerned. On the other hand, the small changes in energy levels at  $z = 21$  could be due to flow randomization. The observation that the energy levels decrease with height in the two cases considered, suggest that the lower building regions are affected by the terrain features so as to increase the

band-width of the wind loads. The peaks in the spectra around 12 Hz may again be fictitious as pointed out before. The low frequency band spectra, Fig. 2.5-5c, except for minor differences in levels are similar.

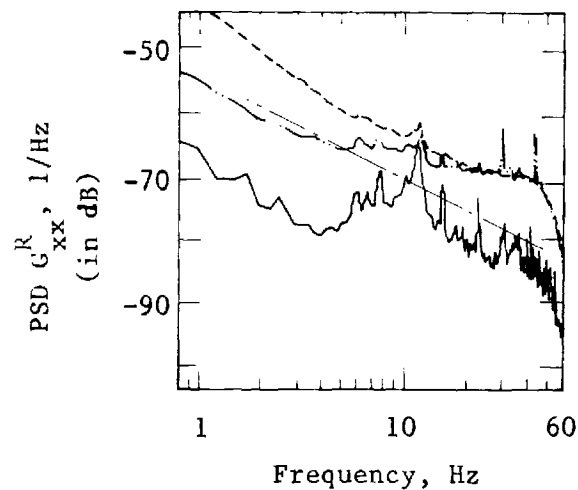
Window NE1 is located at a reentrant corner with respect to the flow direction being considered. The results for this case are shown in Figs. 2.5-6a, b, and c. Referring to (a), it is seen that the pressure coefficients for  $z = 9$  and 21 are predominantly negative, whereas at  $z = 28$  they are more evenly distributed. Table 2.5-2 shows the different  $C_p$  values and the factor  $g$  from Eq. (2.2-3) for the various window locations considered. The aerodynamic effects are more complex at this location than the previous cases, as is evidenced by the shape of the histograms. Even after taking into account variance in the pdf estimates, the histograms at  $z = 21$  and 28 appear to have periodic components as manifest by the presence of dual peaks. A cross-check with the respective spectra does not reveal any discrete frequency component of comparable amplitude.

The high frequency band spectra are now reversed in energy levels, in that the overall energy levels, and hence the mean square values, increase with height. This suggests that in a region of separation, such as can be expected at NE1, the rms values of the pressure fluctuations increase with height. For  $z = 21$ , the PSD contains a discrete frequency peak at 30 Hz and it is suspected that rotating machinery in the building (such as 1800 rpm motors used in the air handling unit) may have caused it to appear in the signal.

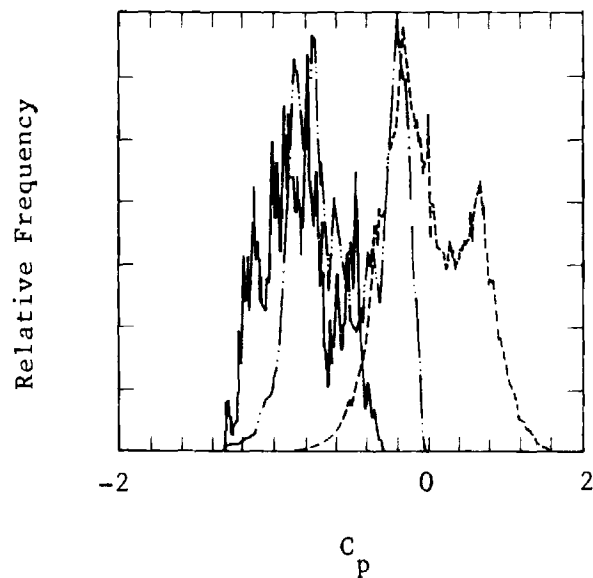
The low frequency band spectra of Fig. 2.5-6c show level variations similar to the high frequency band; however, the PSD for the 9th floor



(c) Low Frequency Band PSD



(b) High Frequency Band PSD



(a) Normalized Histograms

Notes:

————— 9<sup>th</sup> Floor  
 - . . - 21<sup>st</sup> Floor  
 - - - - - 28<sup>th</sup> Floor

PSD in dB Relative to 1, 1/Hz

e.g. -20 dB =  $10^{-2}$ .Figure 2.5-6. Spectra and Normalized Histograms, Window NE1,  $\bar{\theta} = 160^\circ$

Table 2.5-2 Pressure Coefficients and Gust Factors  
for  $\bar{\theta} = 160^\circ$

Floor	Window	$\bar{C}_p$	$C_{p_{rms}}$	$C'_p$	$C_{p_{max}}$	$g = \frac{C_{p_{max}}}{C'_p / \bar{C}_p}$
28	W10	-.307	.313	.061	-.7	3.52
	N4	-.203	.210	.054	-.38	1.43
	NE4*	+.014	.277	+.277	+.86	.043
21	W11	+.008	.075	.075	+.384	.041
	N3	-.197	.206	.06	-0.53	1.73
	NE1	-.55	.622	.29	-1.36	2.6
9	W10	-.398	.524	.339	-2.64	3.08
	N4	+.46	.472	.106	.84	3.66
	NE1	-.591	.764	.234	-1.77	4.47
	NW12	.539	.753	.526	+2.80**	2.87

\*NE4 presents the same profile to the flow as NE1 except for an additional upstream corner.

\*\*Positive  $C_p$  greater than +1 since the local velocities may be considerably different from the reference velocity.

reveals the discrete frequencies, 0.86 Hz and 1.7 Hz. The flow effects in this region are extremely complex, and although the characteristic frequencies could have been produced by a downwash, such a description is at best crude conjecture at this time.

Wind Incidence,  $\bar{\theta} = 280^\circ$ : Figures 2.5-7a, b, and c show the PSD's and histograms for window NE1 at  $z = 9$  and 21. At  $z = 9$  the windward location of NE1 together with the larger corner dimension results in a  $C_p$  histogram which is almost entirely on the positive side.\*

The high frequency band spectra, Fig. 2.5-7b, show that the mean square value of the pressure fluctuations is larger at  $z = 21$  than at  $z = 9$ . On the other hand, in the low frequency band, (c), the PSD levels at frequencies below 0.5 Hz are greater at  $z = 9$ . Considering the negative  $C_p$ 's at  $z = 21$  and differences in PSD levels, it appears that an upflow component may exist and that the building facade features tend to produce increased higher frequency loading.

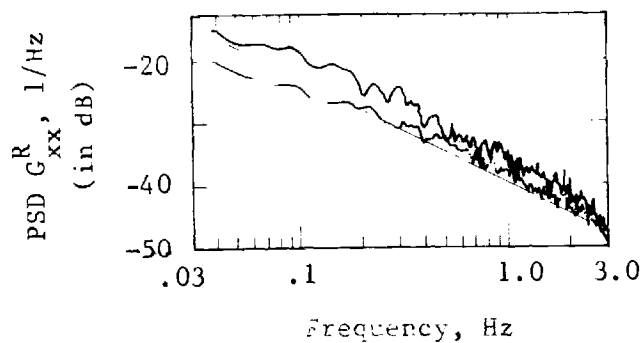
Window N4 is downstream of NE1 with respect to the flow and the wind direction is almost parallel to the wall on which the pressure tap is located. From Fig. 2.5-8a the randomization effect of the corners at the 21st floor is apparent in that at a  $z = 9$ , the  $C_p$ 's are mostly positive and at  $z = 28$  large negative  $C_p$  excursions occur, whereas at  $z = 21$  the distribution is almost symmetric with respect to a zero mean.

The high frequency band spectra, Fig. 2.5-8b, have almost the same slope and closely follow the  $f^{-5/3}$  Kilmogorov inertial subrange spectrum. The PSD level at N4 is reduced from that seen at NE1 for  $z = 21$  whereas for  $z = 9$  it is almost the same. The periodicity at 12 Hz is again asso-

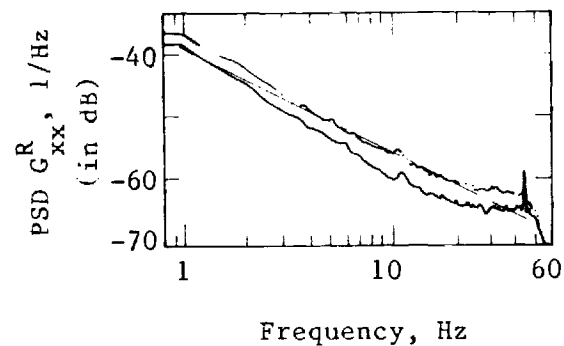
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\*The erroneous values  $C_p > +1.0$  result because of differences in the local speeds and the remote reference velocity of the wind used for normalization.

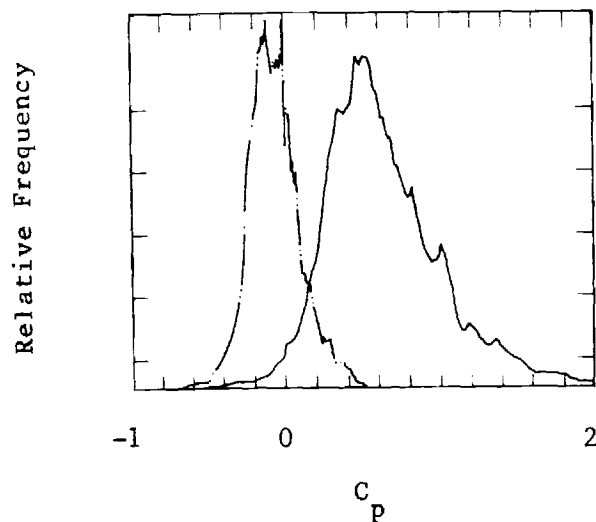




(c) Low Frequency Band PSD



(b) High Frequency Band PSD



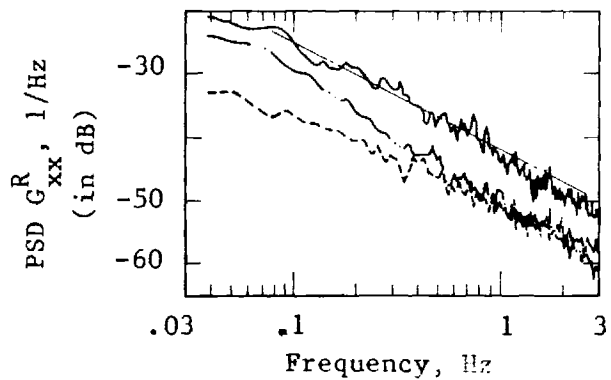
(a) Normalized Histogram

Notes:

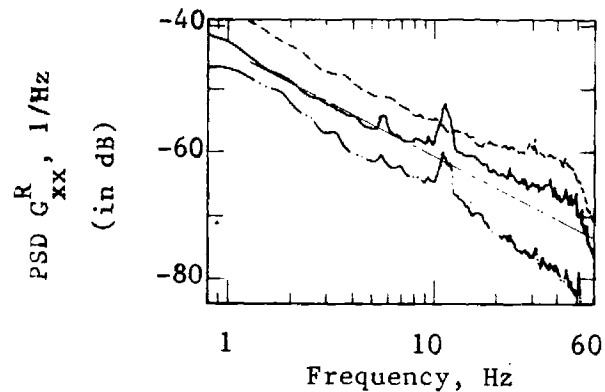
\_\_\_\_\_ 9<sup>th</sup> Floor  
 \_\_\_\_\_ .. \_\_\_\_\_ 21<sup>st</sup> Floor  
 - - - - - 28<sup>th</sup> Floor

PSD in dB Relative to 1, 1/Hz

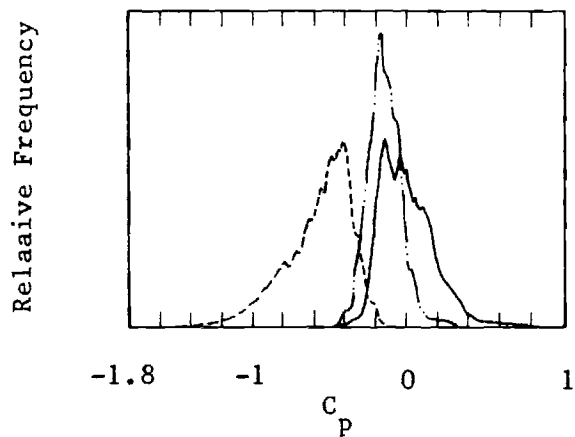
e.g. -20 dB =  $10^{-2}$ .Figure 2.5-7 Spectra and Normalized Histograms, Window NE1,  $\bar{\theta} = 280^\circ$



(c) Low Frequency Band PSD



(b) High Frequency Band PSD



(a) Normalized Histogram

Notes:

————— 9<sup>th</sup> Floor  
 — · — · — 21<sup>st</sup> Floor  
 - - - - - 28<sup>th</sup> Floor

PSD in dB Relative to 1, 1/Hz

e.g. -20 dB =  $10^{-2}$ .Figure 2.5-8 Spectra and Normalized Histograms, Window N4,  $\bar{\theta} = 280^\circ$

ciated with the window response, although the amplitude of this component is higher than those observed in other cases. Discrete flow effects such as peaks are not seen in the low frequency band spectra of Fig. 2.5-8c, although the relative levels are the same as in (b).

In concluding this section, some significant observations about the general character of the pressure spectra are made and the identified features of the cladding loads are summarized. With regard to dynamic cladding loads, the overall PSD levels or the mean square value of the loading are of concern.

Higher spectral levels are seen to occur on faces where the pressure coefficients are predominantly negative. In other words the highest spectral levels can be expected when the bulk of the samples in a pressure histogram lie on the negative  $C_p$  side and the magnitude of the negative mean  $C_p$  is largest. The variance of the normalized histogram (pdf) and the area under a power spectrum are both equal to the mean square value of the random process (Parseval's theorem). This implies that the higher variance should be associated with the histogram having the larger magnitude of the negative mean pressure coefficient  $\bar{C}_p$ . That this is indeed the case is manifested by the longer negative tails in the histograms as  $\bar{C}_p$  tends to be more negative.

The PSD levels on the windward side increase as the building ground level is approached, which is consistent with the assumption that gusts due to the effects of the surrounding terrain are largely responsible for the increase. The additional observation that in regions of separation, i.e., where the wind effects are not direct, the spectral levels increase with height, serves to reinforce the speculation of the effects of surrounding terrain.

Discrete flow effects such as vortex shedding, narrow band pressure fluctuations (bandwidth approximately 1 Hz), and spikes in the high frequency band spectra, are seen only for some particular combination of  $\phi$ ,  $\theta$ , and  $z$ . Exterior architectural features, such as mullions, are suspected to be the prime contributors to spectrum levels by way of narrow band fluctuations in the flow. Existence of an upflow and consequent vortex shedding at the reverse setbacks is responsible for the low frequency spectral lines corresponding to a Strouhal number of 0.48.

The multiple corners in the building planforms appear to act as roughness elements on the building profile. The relatively small number of large dimension corners, described by the nondimensional coefficient value  $r = 0.34$ , essentially act as bluff extensions of the building floor plan and are largely unaffected by the other corners. On the other hand, for a relatively large number of smaller dimension corners, described by  $r = 0.85$ , the roughening is considerably less. A side with such corners acts essentially as a flat wall, the corners at the ends of which are nearly unaffected by the intermediate corners. A building profile roughness described by  $r = 0.17$ , however, serves to randomize the flow effects, resulting in diminished mean and rms values of the pressure fluctuations (small mean and variance in the pdf, and low PSD levels). A preliminary indication, therefore, is that a roughness ratio in the vicinity of  $r = 0.2$  could have an alleviating effect on the cladding loads.

A feature of the low frequency band spectra that deserves mention is that the spectra do not level off or show a peak at the extremely low

frequencies. This behavior is expected because of the finite energy in the pressure signal. The absence of a peak or limiting behavior could be due to the restricted bandwidth of the spectra. At these low frequencies the external wind pressures are no longer distinct from the internal building pressure variations due to normal building operations. The period of the internal pressure fluctuations could be considerably longer, perhaps extending over an entire day. Periods longer than about 30 secs. are not encompassed by the low frequency band spectra. Several attempts were made to explore this point by analyzing the data with ultra-low  $F_{\max}$  and  $\Delta f$  but were unsuccessful because of the inability to adequately reduce the variance by processing enough sample records over the 1 - 2 hours during which stationarity could be assumed. Also, attempts to measure an absolute internal pressure and compute the resulting PSD with sufficient accuracy and fidelity proved beyond the present transducer limits.

An inspection of the histograms for pressure taps located on a flat face with the flow nearly parallel to the face, shows that the distribution of the instantaneous  $C_p$ 's moves to the negative side as the height increases. This is not a surprising result, but only an instance substantiating two of the conclusions drawn earlier. Specifically, it was seen that the spectral levels increase with height in regions of separation, and that higher spectral levels are associated with predominantly negative pressure coefficient histograms, i.e., where the bulk of the samples in a pressure histogram lie on the negative  $C_p$  side. The above observation, therefore, provides increased confidence in the consistency of the data.

The overall levels of the spectra in the Reynolds number range of  $1 - 5 \times 10^6$  are in general quite small and would require considerable dynamic amplification by structural vibration to approach static load levels. For this type of geometry at least, this must be taken as a negative conclusion to the original conjecture that significant wind pressure excitation power could be found at frequencies of concern for local structural dynamics. Finally, the slopes of the spectra differ only slightly from the Kolmogorov spectrum, although for energy levels they tend to be flatter.

#### 2.5(b) Window Response Causality

The salient features of the localized wind pressure loads, particularly the dynamic aspects, were identified from experimental data in the previous section. It is strikingly apparent from those results that there is an extremely small amount of power in the pressure loading over those frequencies of concern for localized structural dynamic response. However, the present field study, while of a somewhat general nature owing to the variable geometry of the building, cannot be considered as completely representative case insofar as wind loads are concerned. Consequently, it is still essential to identify the window response levels associated with the wind loads, and equally important to ascertain over what frequencies these take place. For example, it is reasonable to expect the static or zero frequency response to depend on factors other than wind (e.g., stack effect, thermal loading, etc.). At the ultra-low frequencies (periods down to about 10 sec.) it is likely that the response is due largely to wind effects, while at frequencies above this point the causality is not clear and is the subject of the present study.

The identification of causality can be accomplished by means of the coherence function, the derivation of which has been previously discussed. In the present applications the coherence function as estimated from the measured wind load and response data serves to establish the causality of the wind loads for the response. In a practical measurement situation the measurement noise also affects the coherence function and can lead to erroneous, low values. These considerations and the evaluation of the coherence function are discussed in the following section in light of the experimental results.

The present implementation of the coherence function is based on the single input/single output problem of Fig. 2.5-9. The true source, or input, of the system is  $u(t)$ , assumed to be stationary, and the output is  $v(t)$ . Extraneous noise at the input and output measurement points is represented by  $n(t)$  and  $m(t)$ , respectively, and is assumed to be incoherent with  $u(t)$  and  $v(t)$ . The quantities  $x(t)$  and  $y(t)$  are the measurement signals and correspond to the wind pressure and window displacement, respectively. Note that  $x(t) = u(t) + n(t)$  and  $y(t) = v(t) + m(t)$ .

Assuming that  $m$  and  $n$  are uncorrelated\* with  $x$  and  $y$ , it follows:

$$G_{xx}(f) = G_{uu}(f) + G_{nn}(f)$$

$$G_{yy}(f) = G_{vv}(f) + G_{mm}(f)$$

$$G_{xy}(f) = G_{uv}(f)$$

Then using Eq. 2.4-11 for the present case:

$$G_{vv}(f) = |H(f)|^2 G_{uu}(f)$$

$$G_{uv}(f) = H(f) G_{uu}(f)$$

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\*For the effect of correlated noise on the coherence function, See Ref. 51.

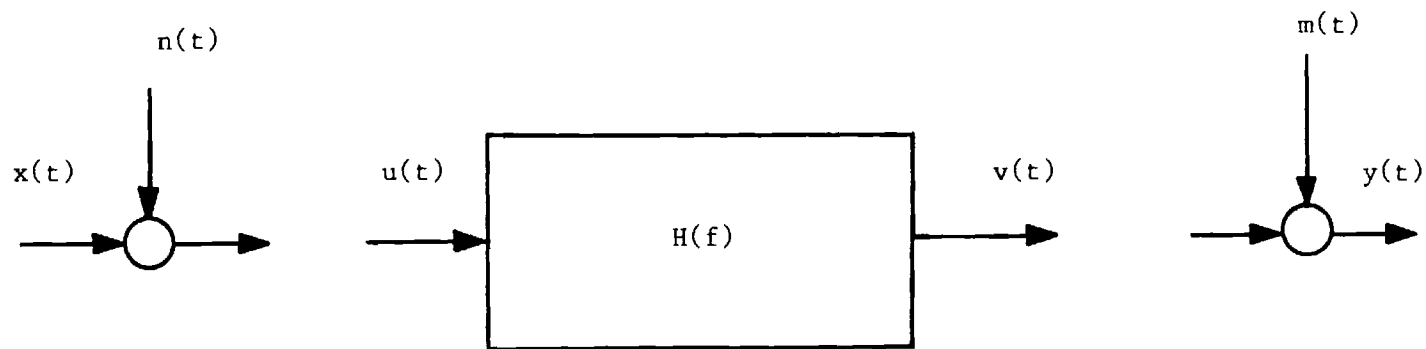


Figure 2.5-9 Ordinary Coherence Function Measurement Model



The true system coherence function  $\gamma_{uv}^2(f)$  is defined as:

$$\gamma_{uv}^2(f) = \frac{|G_{uv}(f)|^2}{G_{uu}(f)G_{vv}(f)}$$

and for a CPLS system is identically equal to one for all frequencies.

The true coherence is related to the measured coherence function,

$\gamma_{xy}^2(f)$ , as follows:

$$\begin{aligned}\gamma_{uv}^2(f) &= \frac{|G_{xy}(f)|^2}{G_{xx}(f)G_{yy}(f)} \\ &= \frac{|G_{uv}(f)|^2}{[G_{uu}(f) + G_{nn}(f)][G_{vv}(f) + G_{mm}(f)]} \\ &= \frac{\gamma_{uv}^2(f)}{[1 + \frac{G_{nn}(f)}{G_{uu}(f)}][1 + \frac{G_{mm}(f)}{G_{vv}(f)}]}\end{aligned}$$

Thus, the measured value of the coherence function is diminished from the true value by terms involving the measurement signal-to-noise ratios. Clearly  $0 \leq \gamma_{xy}^2(f) \leq 1$ . In the absence of any system nonlinearities the ratio  $\gamma_{xy}/1-\gamma_{xy}$  can be interpreted as the signal-to-noise ratio at the output. Alternatively, for a noise-free system  $\gamma_{xy}/1-\gamma_{xy}$  can be interpreted as the linear-to-nonlinear ratio. In practical measurement systems the various spectra are averaged over a sufficiently large number of samples to remove the uncorrelated noise terms.

In the present experimental study the coherence function was applied to the systems shown in Fig. 2.4-5(a) to determine the degree of linear dependence of the response on the measured differential pressures,  $\Delta p$ . Several preliminary remarks are in order before the computational procedure and the results are detailed.

1. From the physics of the phenomenon it is apparent that sustained vibration of the window system cannot be expected. Hence selected sample records need to be used to compute the coherence function.

2. The sample records should be such that the response signal-to-noise ratio is considerable, i.e., the response amplitude should be above the noise floor of the measurement and data acquisition system.

3. Such a selection can be accomplished by triggering the data input to the Fourier Analyzer using the pressure (or response) signal.

4. Sorting the data as in 3, it was found, restricts the number of sample records that can be extracted from an hour-long data record with stationary wind characteristics. Alternatively, therefore, visual inspection and sorting as described in Section 2.4(c) can be used.

5. For inputs which are not triggered in (This applies to the processing equipment used and other similar analyzers.) overlap processing of the data can be used to reduce the estimation errors in the coherence function(62). This estimation technique consists of applying the FFT operation to overlapped data segments. The result is a reduced variance at the expense of the number of FFT's to be performed, i.e., computational time. Overlap processing is readily implemented on the 5451B Fourier analyzer. The overlap fraction can be varied by the addition of a delay loop to the program.

The coherence function was computed from the measured data for the following cases:

1. Coherence between the differential pressure and the interanl glass plate displacement, i.e.,  $\gamma_{\Delta p, w_1}^2$

$$2. \gamma_{\Delta p, w_e}^2$$

$$3. \gamma_{w_1, w_e}^2$$

The first two of the above determine the fraction of the response power due to  $\Delta p$ . The function  $\gamma_{w_1, w_e}^2$  is used to verify the existence of coupling between the two plates.

All of the previously-discussed computational procedures, namely, triggered (without Hanning), free run and overlap processing were utilized to obtain the best possible coherence function estimate for the various signal pairs. Differences in the coherence functions estimated by the four techniques (This includes triggering the input on response.) were insignificant and were attributed to inherent uncertainty in the estimation procedures. Overlap processing did smooth the coherence function to some extent although the difference was not substantial to warrant an increase in computation time.

The coherence functions shown in Figs. 2.5-10 thru 2.5-12 were computed from record segments selected according to the sorting procedures described earlier. The input was free run, the sample records were disjoint and weighted by a single Hanning function. The high frequency band (0.8 - 50 Hz) functions were computed from the respective spectra averaged over 400 disjoint segments. The low frequency band (0.03 - 3 Hz) functions utilized 30 disjoint records. The coherence functions of Figs. 32, 33, are representative of the response dependence on the differential pressure. The actual wind and structural variables shown are for reference purposes only.

Figs. 2.5-10a, b, and c show the coherence functions  $\gamma_{\Delta p, w_e}^2$ ,

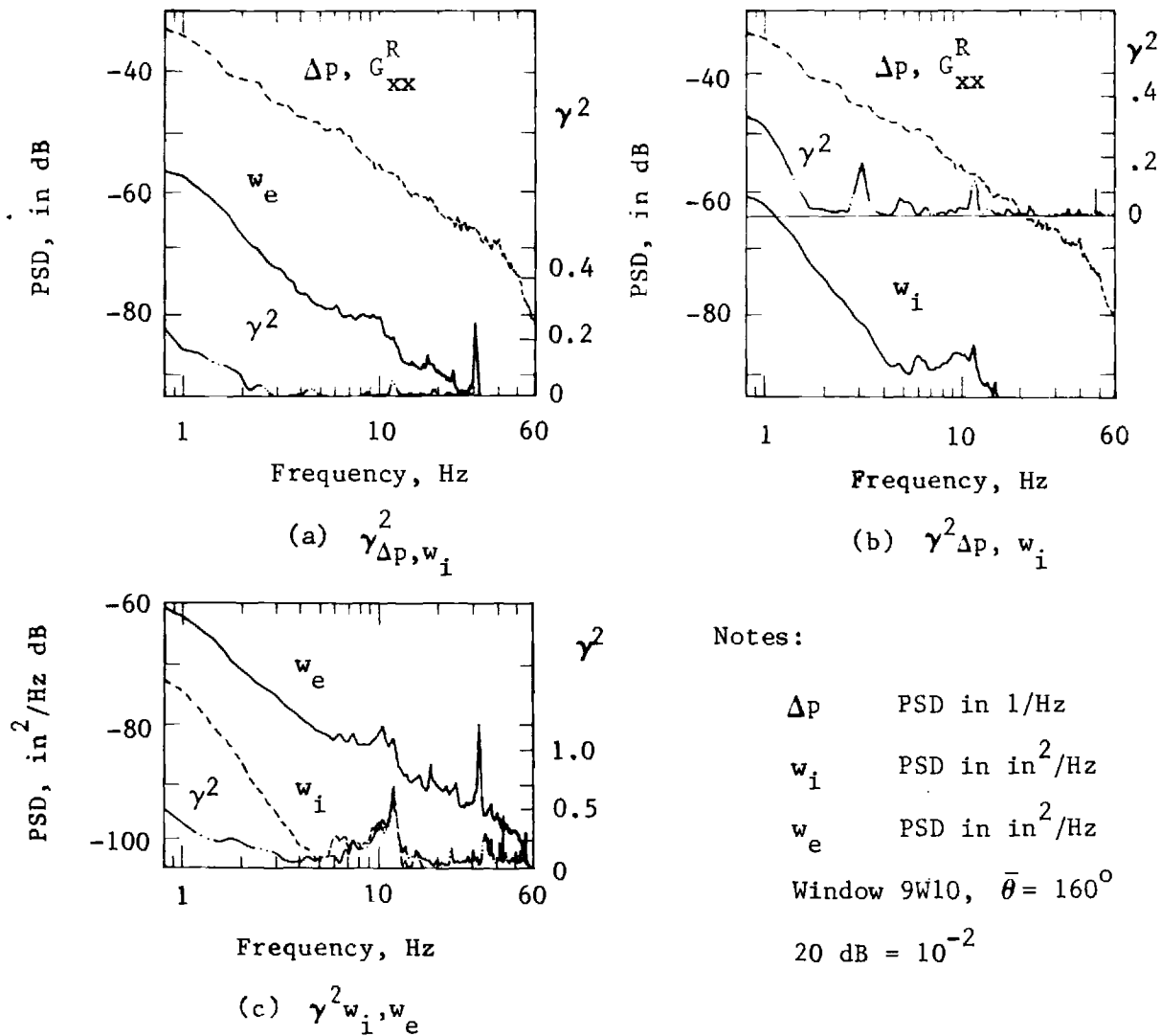


Figure 2.5-10 Measured Coherence Function, High Frequency Band

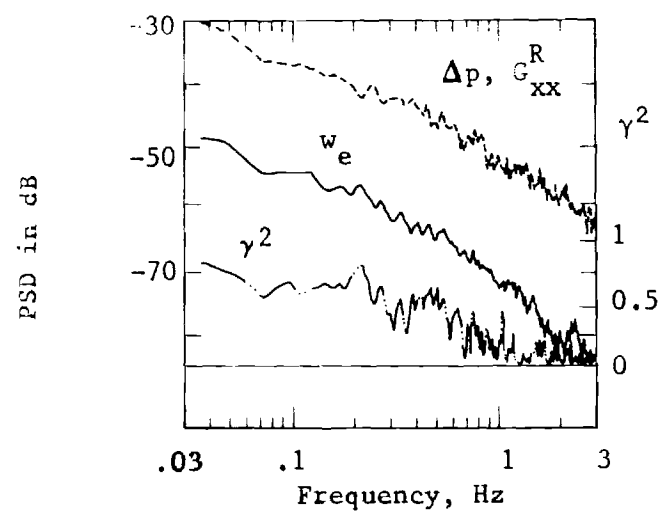
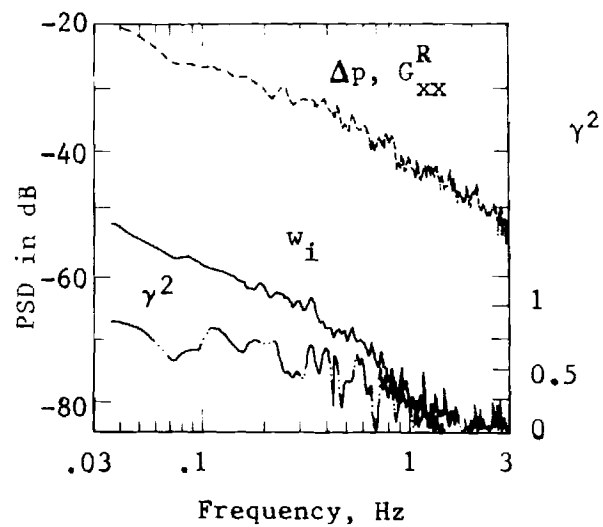
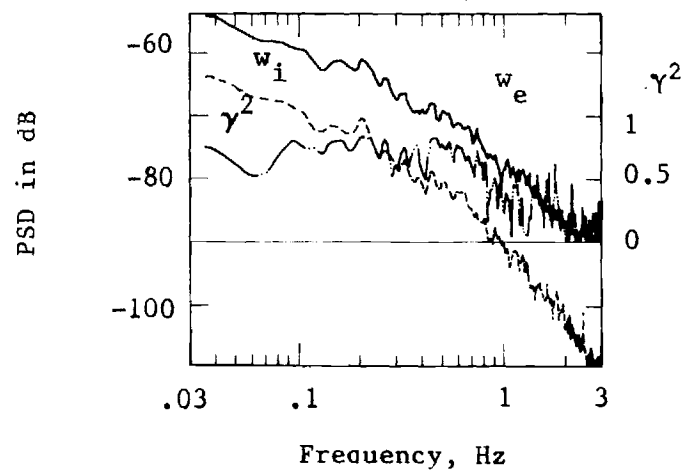
$\gamma_{\Delta p, w_i}^2$ , and  $\gamma_{w_e, w_i}^2$ , respectively, together with the input and output spectra, for the high frequency band. Significant values of the coherence functions in Figs. 2.5-10a, b, are seen only below 1.7 Hz. The peak values in the two cases of  $\gamma_{\Delta p, w_e}^2 = 0.26$  and  $\gamma_{\Delta p, w_i}^2 = 0.35$  at 0.8 Hz differ due to the associated statistical uncertainty. At frequencies above 1.7 Hz there appears to be no causality between the differential pressure and the response. The values of the coherence function between 0.8 - 1.7 Hz are too small compared to 1.0 and majority of the response power is due to sources other than wind loads. Examination of Fig. 2.5-10c reveals similar behavior in  $\gamma^2(f)$  up to 3 Hz. From the spectra it can be seen that the root-mean-square value of  $w_e$  is higher (at all frequencies) than of  $w_i$ . The difference in rms values in the vicinity of 1 Hz is about 6 dB corresponding to a ratio  $\frac{w_e}{w_i} = 4$ . This difference can be attributed to the compliance of the entrapped air. A high value of the coherence function  $\gamma^2 = 0.7$  indicating extremely good causality is seen at 11.5 Hz which is close to the natural frequency for inphase motion of the two plates in the (1,1) mode calculated from the model (Chapter 3). However, Figs. 2.5-10a and b do not show any pressure-to-response causality at this frequency. This implies that the vibratory motion of the window is due to sources other than the wind loading. The internal glass plate response rolls off much faster than the external plate response at frequencies above 1 Hz. A strong peak in  $w_e$  is seen at 31 Hz and compares with the calculated (3,1) inphase motion frequency. A corresponding peak in  $w_i$  is not so obvious. The coherence function at this frequency is of the order of 0.26. An explanation for the discrepancy at 31 Hz is that the

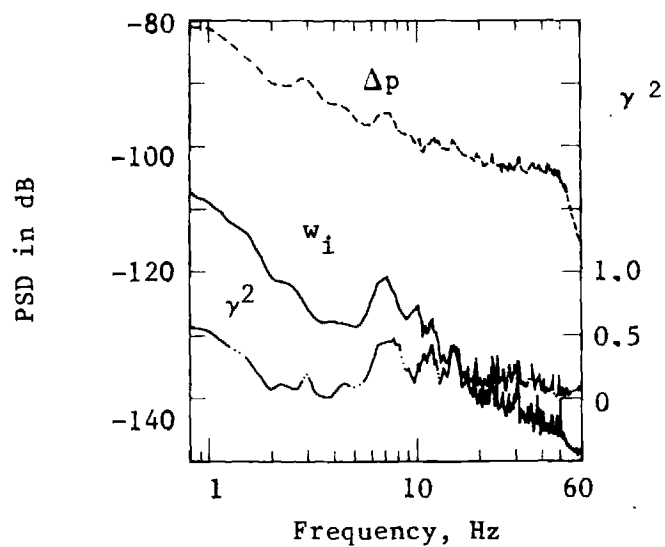
peak is fictitious and not due to glass plate motion. The frequency could have been introduced in the signal due to floor vibrations arising from rotating machinery used in building services (generally 1800 rpm motors). The LDM being more sensitive to floor vibrations, the frequency shows up prominently in  $w_e$ . The erroneous coherence function value is typical of the false coherence case described in Ref. 31.

Figures 2.5-11a, b, and c show the coherence functions in the low frequency band (0.03 - 3 Hz) corresponding to the data used in Fig. 2.5-10. This low frequency loading can be assumed to be quasi-static. Excellent causality between pressure and response is seen up to 0.6 Hz from Figs. 2.5-11a and b. This leads to the conclusion that the "static" response of the window systems is largely due to pressure loads. It is noted here that since only the differential pressure was measured and the internal pressure contributes to the quasi-static differential, a distinction as to the source of the pressure loads cannot be made from the data collected, in this low frequency range.

The roll-off of  $w_e$  and  $w_i$  spectra, as seen from Fig. 2.5-11c is almost identical up to 1 Hz and the difference in spectrum levels at 1 Hz is approximately 12 dB corresponding to a ratio  $\frac{w_e}{w_i} \approx 4$ .

Figure 2.5-12 shows the high frequency band spectra and coherence functions for a case which differed markedly from the representative response spectra and coherence function. The difference is seen in the frequency band extending from 6 - 16 Hz, where the coherence function is very close to 0.5. Considering the statistical variation in the estimates this is within the 95% confidence limits of true  $\gamma^2 = 0.5$ .

(a)  $\gamma_{\Delta p, w_e}^2$ (b)  $\gamma_{\Delta p, w_i}^2$ (c)  $\gamma_{w_i, w_e}^2$ Figure 2.5-11. Measured Coherence Function,  
Low Frequency Band



(a) Input  $\Delta p$ , Response  $w_i$ ,  $\gamma^2 \Delta p$ ,  $w_i$

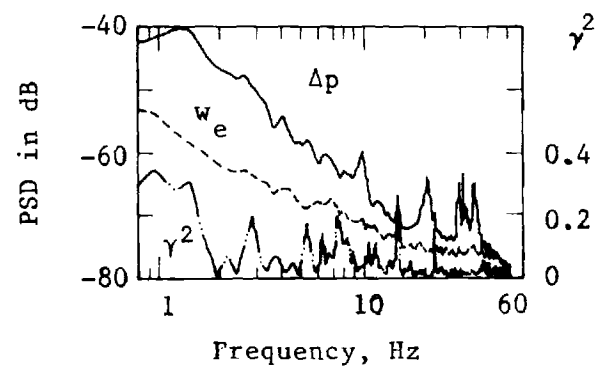
Notes:

$\Delta p$  PSD in 1/Hz

$w_i$  PSD in  $\text{in}^2/\text{Hz}$

$w_e$  PSD in  $\text{in}^2/\text{Hz}$

Window 9N4,  $\bar{\theta} = 160^\circ$



(b) Input  $\Delta p$ , Response  $w_e$ ,  $\gamma^2 \Delta p$ ,  $w_e$

Figure 2.5-12 Measured Coherence Function, Window 9N4,  $\bar{\theta} = 160^\circ$



The pressure spectrum shows a bulge between 6 - 8 Hz indicating that the wind pressure is responsible for the response at these frequencies. Since the window at which these data were measured was on the 9th floor, the relatively high frequency response is possibly due to gust effects.

The coherence function, as has been demonstrated, serves to identify the response levels associated with the wind or pressure loads. Several features of the window response dependence on wind loads are now apparent. High coherence between the input and output up to 1 Hz shows that the "quasi-static" response of the glass plates is primarily due to pressure loading. Beyond 1 Hz the coherence value decreases and becomes insignificant for frequencies less than 1.7 Hz. In the region from 1 to 1.7 Hz the values of the coherence function are in the vicinity of 0.25 implying that pressure loads may no longer be the sole excitation source. Interpretation of the coherence function at resonant peaks requires that care be exercised so as to isolate false coherence which may be due to an extraneous source affecting both the input and the output. At the fundamental frequency of 11.5 Hz (for window W10 on the 9th floor) which appears in the response spectrum the coherence between the pressure and the response is extremely low. Therefore, the fundamental frequency is excited by sources other than the wind loading. In some instances the fundamental was also seen in the pressure spectra yielding a coherence of almost unity (it is noted that for purely sinusoidal signals the coherence function can only be 0 or 1 ). This was diagnosed as a typical false coherence since the fundamental was strongly excited and appeared, by way of motion of the pressure tap located on the support and/or feedback

of panel vibrations to the flow, in the pressure spectra. False coherence also appears to be the reason for the high coherence function values seen at 31 Hz. This frequency corresponds to the rpm of some rotating machinery in the building.

The isolated case of Fig. 2.5-12 demonstrates that high values of the coherence function in the high frequency band may occur for a particular combination of the building profile and flow characteristics. In this case the wind loads are capable of exciting window dynamic response, in a relatively narrow frequency band extending from 6 - 16 Hz. In view of these observations the hypothesis of "critical response locations" appears more realistic.

#### 2.5(c) Window Dynamic Properties

The previous two subsections have attempted to characterize the localized wind pressure loads on windows and to assess the causality between these loads and the response. This subsection deals with the problem of measuring the transfer function of the window system from which the modal properties of the double plate system and the support transmissibility can, in turn, be estimated. The "inverse" problem thus posed is that of obtaining a mathematical model for the windows system from measured data and is referred to as System Identification (81). There are a variety of time domain and frequency domain techniques that have been employed for system identification, and the application of transfer function-based schemes is dealt with in Refs. 82 - 85.

These types of measurements are usually handled most conveniently in a laboratory and could be performed on either model or full-scale window assemblies. Such a program, including design and construction

of a dynamic pressure loading facility for full-scale testing and response measurements for window assemblies similar to those at the field site, is reported in Chapter 4. However, there still exists a need for measurements of the dynamic response characteristics of window units as they exist in the field under actual service conditions. This can best be argued on the basis of three points:

(a) The dynamic characteristics of the window unit depend on the nature of its installation and the fact that it is affected by adjoining elements of the curtain wall. It would be unrealistic to simulate the entire curtain wall. In addition, any significant resonant frequency variation from one window to another can be treated as indicative of the variation in installation characteristics. A correlation between the frequency variation and the nature of the installation would enable detection of a faulty installation during assembly of the curtain wall or at a later time in its service life when structural degradation might have occurred.

(b) Measured transmissibility functions for a window provides an estimate of the support force attenuation arising from the presence of the rubber gasket. The transmissibility function can, therefore, be used to assess the quality of the gasket and gasket installation.

(c) In-situ dynamic measurements such as resonant frequencies and damping characteristics are needed for comparison with the results of laboratory simulation tests in order to verify the fidelity of those tests or improve the methodology.

In Sections 2.1 and 2.2 it was envisioned that the measured ambient wind pressure loading, at least in certain critical locations

about the building, would contain sufficient power to excite an appreciable dynamic response in windows and cladding components. The full-scale wind pressure and window response measurements reported in Section 2.5(a) and (b) have, however, shown that the dynamic wind loads (pressure loading with power spanning the first few natural frequencies of the plates, say from 5 - 30 Hz) are extremely small in magnitude and do not bear significant causality for the window dynamic response. The low values of the coherence function at higher frequencies have another important implication. As mentioned in Section 2.4(b), the coherence function is a measure of the validity of the transfer function estimates. It is evident, therefore, that at higher frequencies the transfer function of the window system cannot be estimated from ambient data.

An examination of the response levels from the  $w_i$  and  $w_e$  spectra shown in Figs. 2.5-10 to 2.5-12 reveals that the response of the double plate system is essentially quasi-static.\* This characteristic of the measured window response provides an increased confidence in the measured wind load spectra which were also seen to be quasi-static.\* It can be concluded, therefore, that at all of the locations where the measurements were conducted, the wind loads do not possess sufficient energy to excite window vibrations. The mullion acceleration response spectra and its coherence function were also computed and are shown in Fig. 2.5-13. As can be anticipated, estimation of the transmissibility function from these data is not possible either.

Considering these factors, the dynamic properties of a window in an actual curtain wall must be determined from some type of forced excitation

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\*Response period large compared to the period of the fundamental.

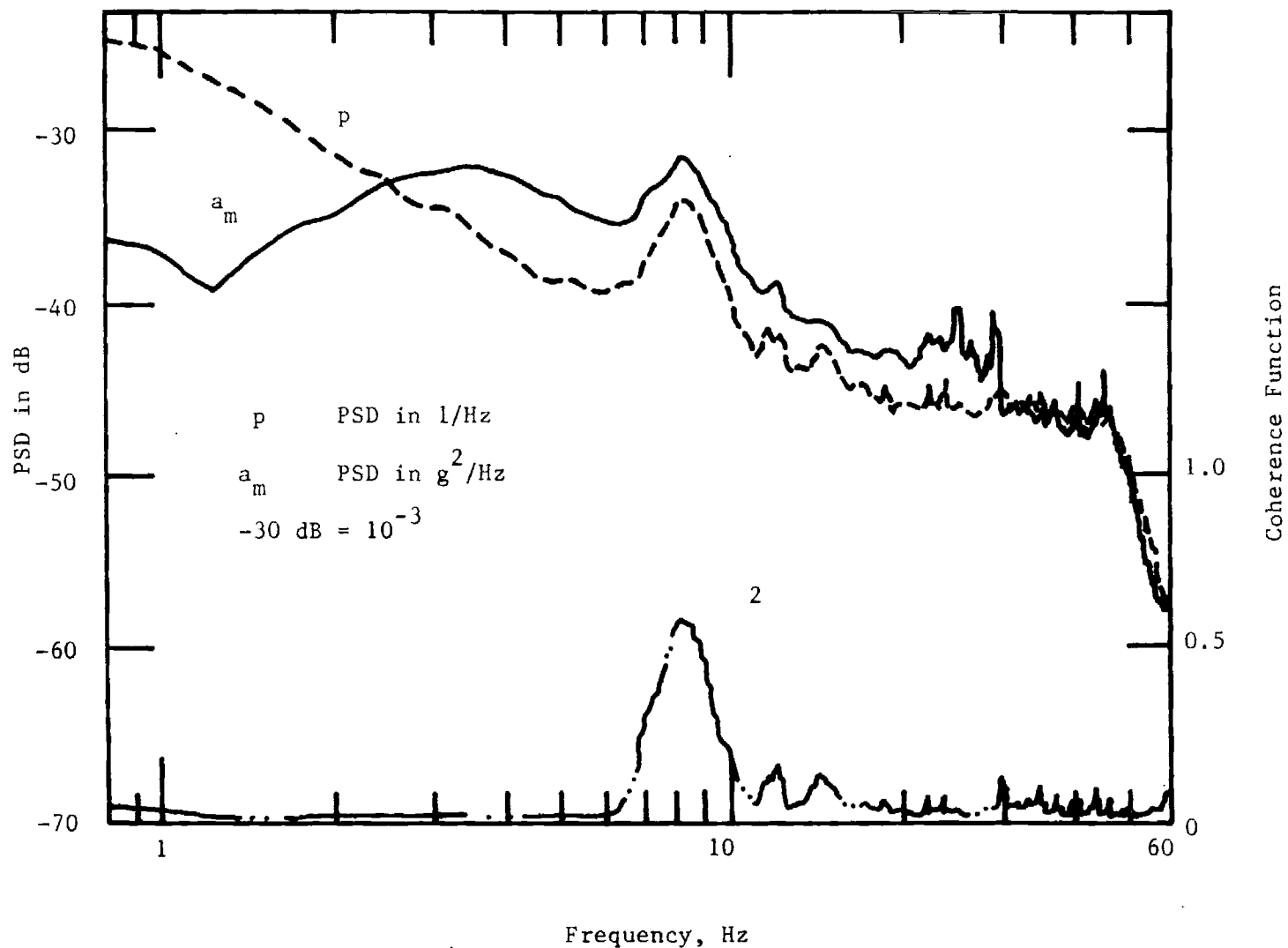


Figure 2.5-13 Mullion Acceleration Response Spectrum and Coherence Function

rather than from ambient loading. A variety of techniques are available for this purpose, but all can be classified as either (i) periodic, (ii) random, or (iii) transient. There is controversy as to which method is the most efficient for a given application (86), since some approaches trade off test simplicity for computational complexity, and vice versa. For the present application where field set-up time and test simplicity are of paramount importance, a transient forcing method was adopted. In conjunction with complex digital processors, such as the Fourier Analyzer used in these tests, this technique is capable of providing faster measurements compared to traditional sinusoidal excitation since the actual testing phase is considerably simplified.

A simple method of performing transient tests is to use a hand-held hammer with a load cell mounted on it to strike the structure and then measure the response to this transient excitation. The impulsive input has a flat frequency spectrum whose bandwidth depends on the "hardness" of the impacting head, a harder head providing a larger bandwidth. The transfer function computed from such an input-response pair can be used to determine the resonant frequencies and the damping coefficients. Multi-point measurement of the transfer function, i.e., computation of the transfer function between various excitation and response points can be used to determine the mode shapes corresponding to these frequencies. The modal mass and modal stiffness can also be estimated from the transfer function data, although this almost invariably requires some type of data fitting (84).

Measurement of the frequency response functions for the present application to window systems was directed at obtaining the resonant

frequencies,  $f_n$ , the viscous damping coefficient,  $\zeta$ , and the mode shapes of the inner plate. The mode shapes of the outer glass plate were not obtained since multiple point response measurements by the LDM proved time-consuming and inconvenient. However, results of the analytical model of the window system (Chapter 3) are used to obtain some idea of the double plate mode shapes corresponding to measured frequencies. In addition, the transfer function between  $w_i$  and  $w_e$  provides the vibration phase angle between  $w_i$  and  $w_e$  at the resonant frequencies and, therefore, the nature of the motion.

The procedures described below for the estimation of the vibration parameters are based on the following assumptions:

- (i) The modes of the system are widely spaced in frequency.
- (ii) Although modal coupling exists, a lightly-damped multi-degree-of-freedom system can be closely approximated by a single-degree-of-freedom system of appropriate natural frequency and damping near each of the system's natural frequencies.

The measured dynamic properties are:

- (i) The resonant frequencies of a single window unit and the damping coefficients for each of the first 6 modes;
- (ii) Variation in the resonant frequency (single mode) for a sample of 12 windows;
- (iii) The mode shapes for the inner plate of the double pane system.

No attempt has been made to develop computational methods for determining these quantities directly from the raw transfer function data since this is beyond the scope of this research, and in addition, several commercial

software systems have already been developed for this purpose.\* Rather, most of the interpretation has been done manually with limited curve-fitting. This provides acceptable accuracy of the lower modes which are relatively widely spaced, but owing to the amount of work involved, only a limited amount of data has been processed.

Modal Parameters: The inner plate of the window was marked at 3 spatial locations as shown in Fig. 2.5-14. The outer plate response was monitored at a single point 2' corresponding to point 2 of the inner plate. An accelerometer was used to measure the response at 2 and its location was fixed. The LDM measured the displacement response at point 2'. An instrumented hammer was used to excite the inner glass plate at each of the 3 points. Transfer functions for these 3 cases were computed from the measured response. The measured transfer function  $H_{2i}(f)$ ,  $i = 1, 2, 3$ , are the elements of the second row of a  $3 \times 3$  transfer function matrix. It can be shown (84) that only a single row or column of the transfer function matrix needs to be measured to compute the modal amplitudes at the 3 points. The LDM output provides the natural frequencies of the outer plate.

The compliance of the hammer head was first adjusted to provide a near uniform energy spectrum out to 50 Hz. Then, the hammer was used to repeatedly excite the structure at a designated point, and the resulting ensemble-averaged auto and cross power spectra were used to compute the transfer and coherence functions. The excitation was stopped when an acceptable coherence function was obtained, and the entire process was then repeated for each designated point. Transfer functions for the inner plate of window N4 on the 28th floor are shown

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\*None of these were available for the present work.



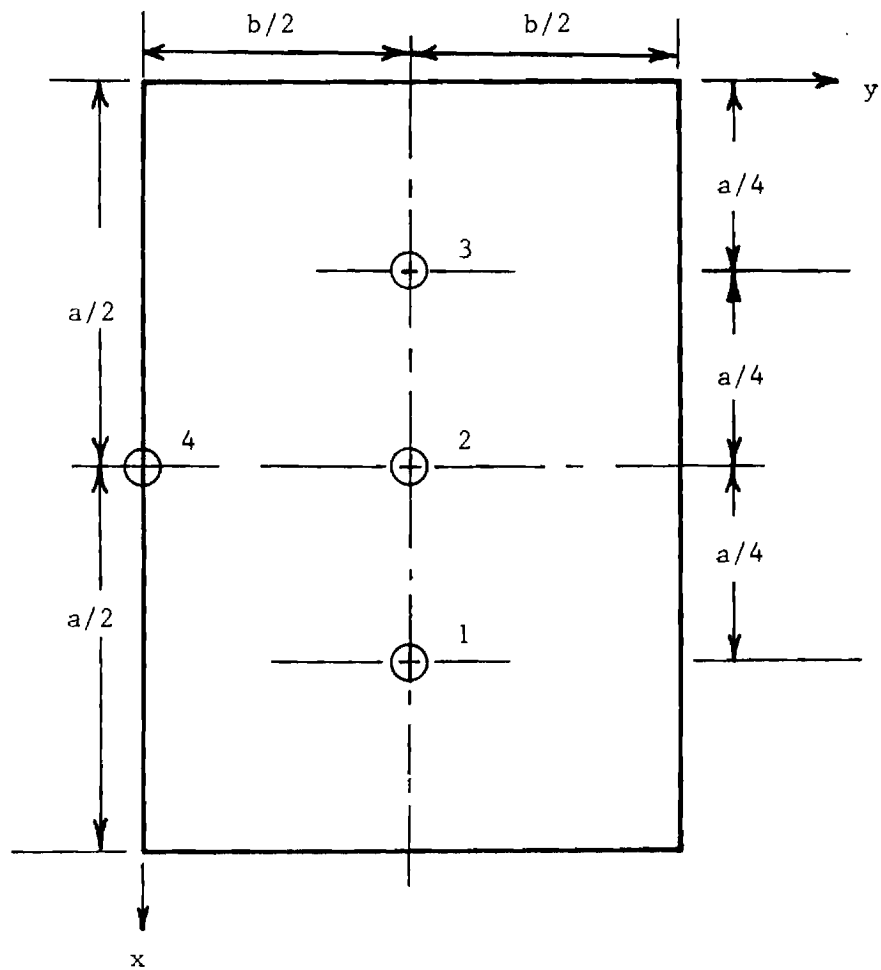


Figure 2.5-14. Location of Accelerometers on the Inner Plate

In Fig. 2.5-15 and Fig. 2.5-16 shows the  $H_{2,1}(f)$  transfer function of the outer plate. The modal parameters computed from these transfer functions are as follows:

Resonant Frequencies: The acceleration-to-force transfer function is called inertance whereas the displacement-to-force transfer function is called compliance. In inertance or compliance plots the natural frequency corresponds to the following transfer function characteristics:

- (a) Peak of the magnitude of the transfer function
- (b) Phase angle of  $90^\circ$
- (c) Zero of  $\text{Re } H(f)$
- (d) Peak of  $\text{Im } H(f)$

In practice, (c) and (d) were used to locate a resonance and the frequency was then determined by examining the imaginary peak.

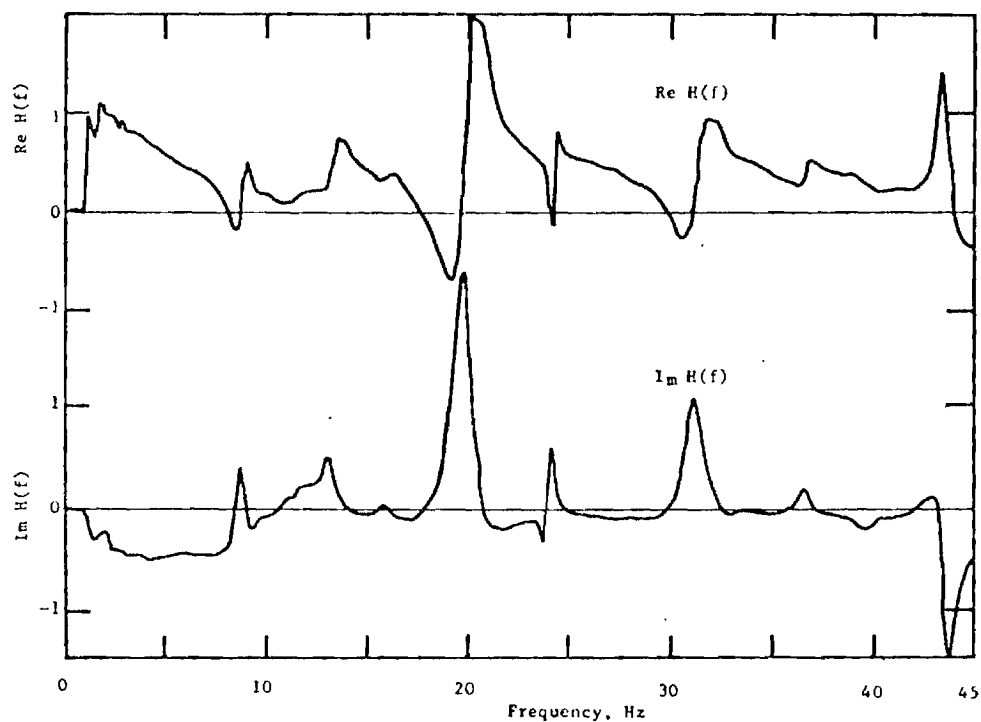
Damping Coefficient: Several techniques for estimating the damping coefficient from the measured transfer function are available. Nyquist plots of the transfer function, where the real and imaginary parts are plotted in the complex plane, can be used to estimate the modal damping (87); however, in practice, estimation of damping from these plots requires careful curve fitting.

Damping is often measured using the amplification factor,  $Q$ , defined as:

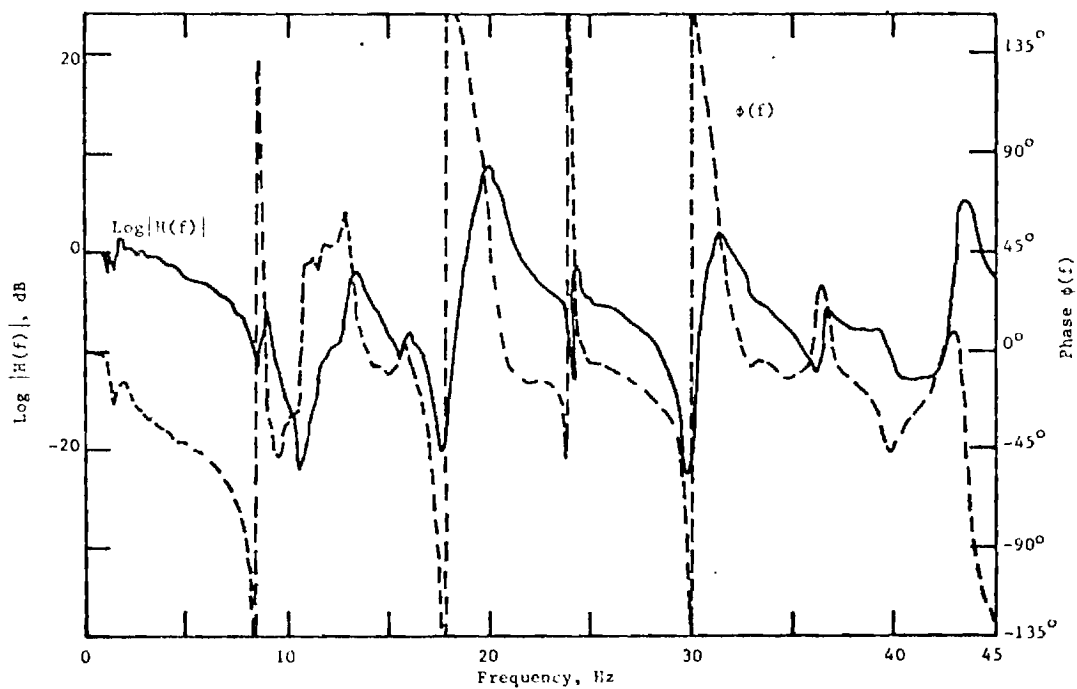
$$Q = \frac{1}{2\zeta}$$

$Q$ , in turn, may be measured from the magnitude of the transfer function as follows:

$$Q = f_n / \Delta f_n$$

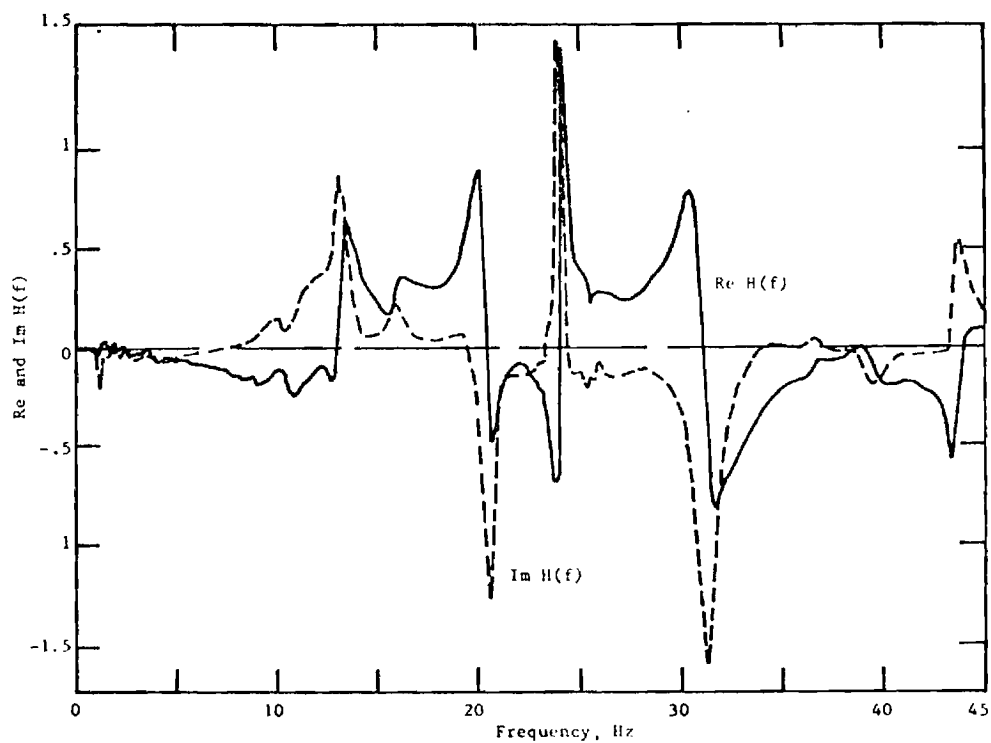


Real and Imaginary

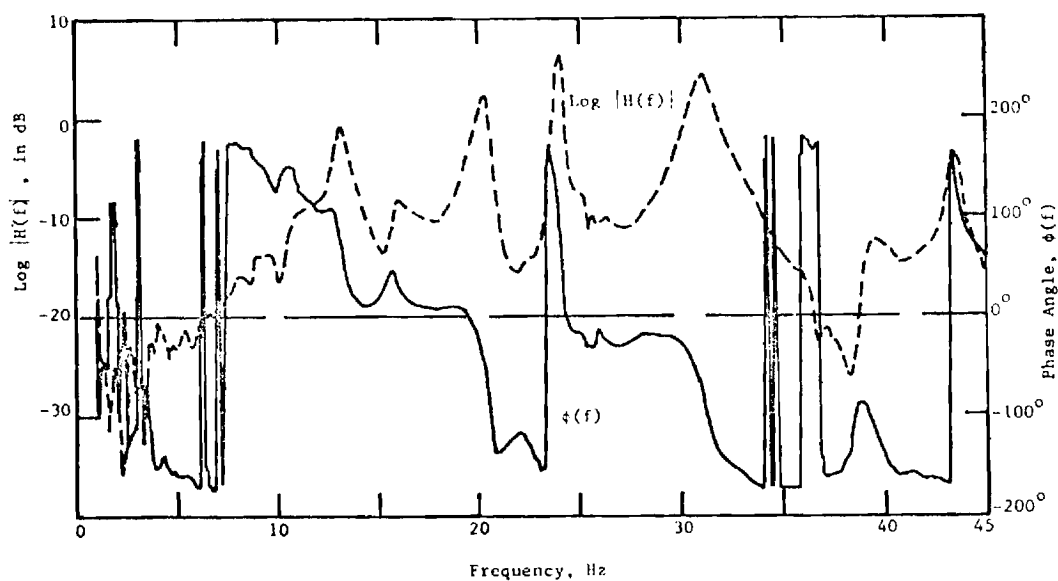


Magnitude and Phase

Figure 2.5-15(a) Frequency Response Function:  
Impact at 1, Response at 1

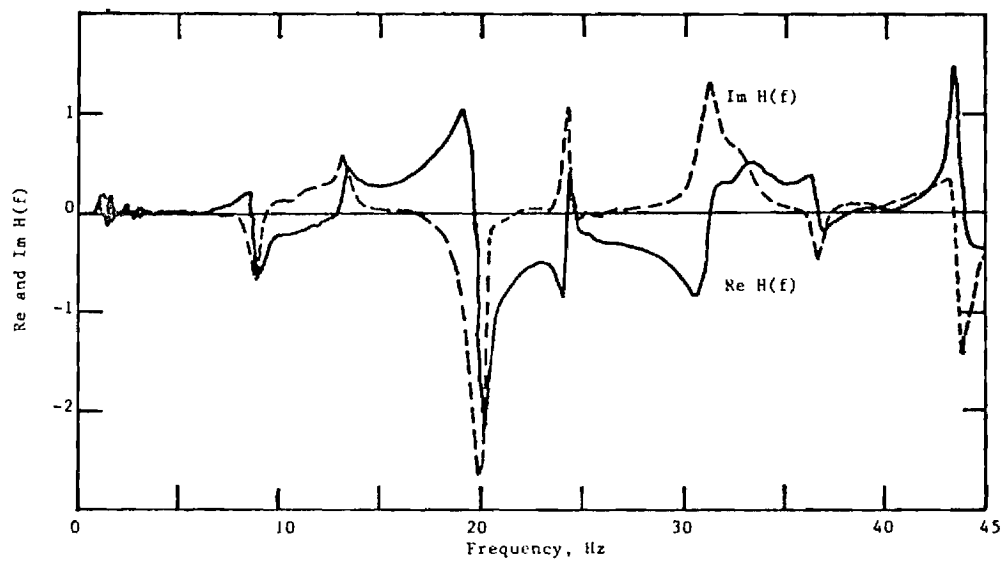


Real and Imaginary

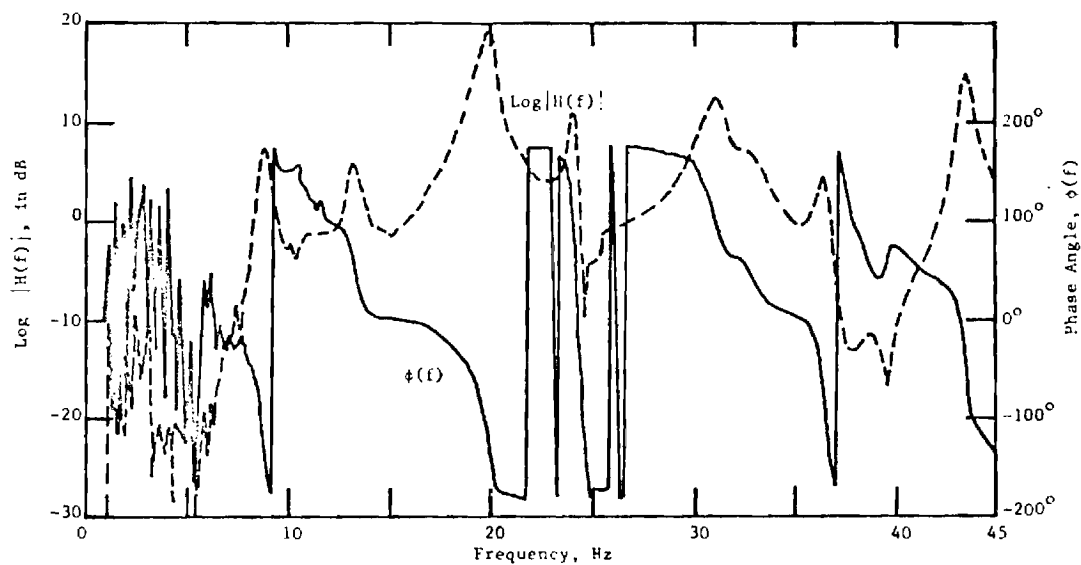


Magnitude and Phase

Figure 2.5-15(b) Frequency Response Function:  
Impact at 2, Response at 1

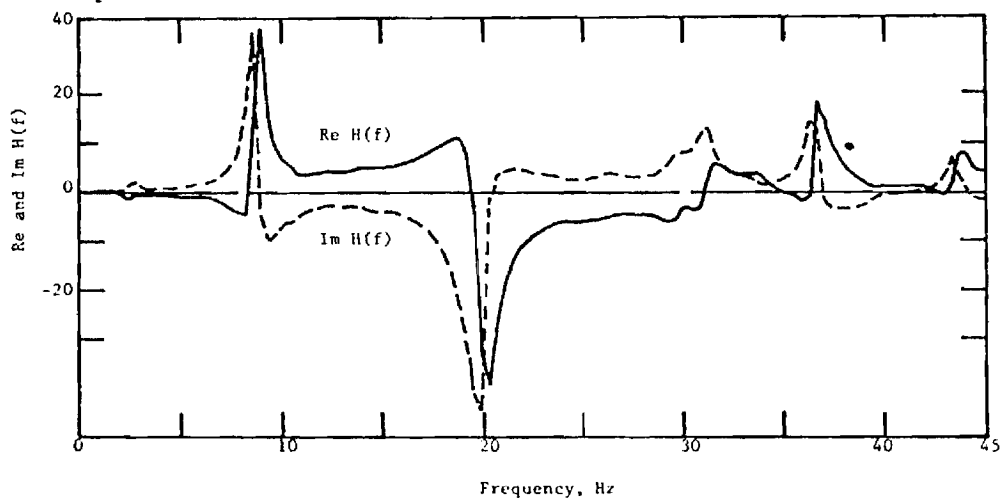


Real and Imaginary

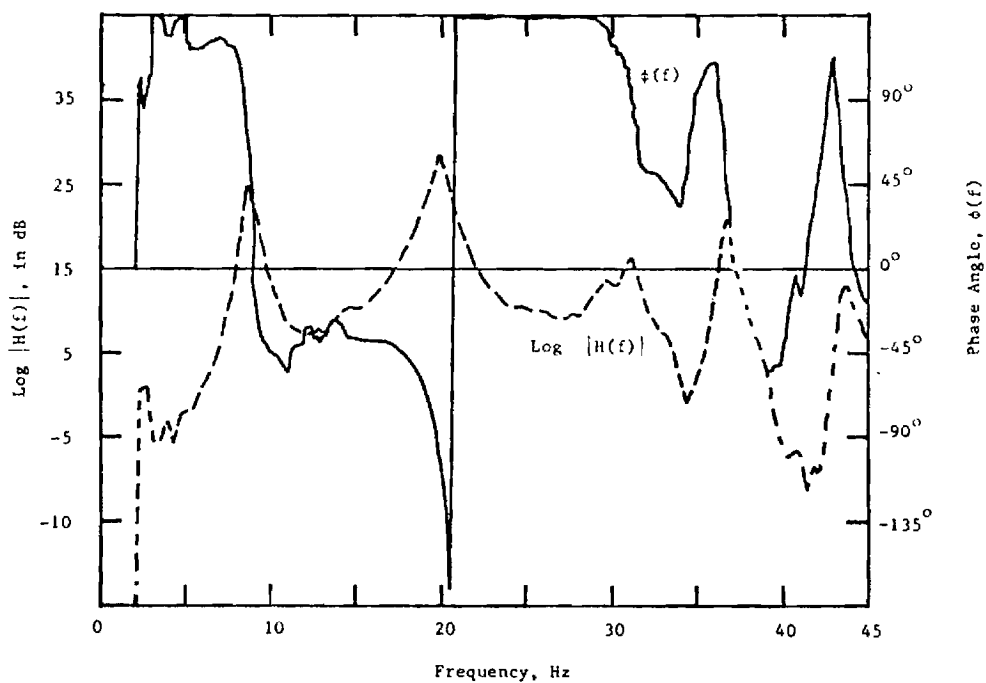


Magnitude and Phase

Figure 2.5-15(c) Frequency Response Function:  
Impact at 3, Response at 1



Real and Imaginary



Magnitude and Phase

Figure 2.5-16 Frequency Response Function:  
Impact at 2, Response at 2'

where  $\Delta f_n$  is the half-power bandwidth of the resonance peak corresponding to  $f_n$ .

Alternatively, a method based on features of the real and imaginary response spectra can be used to estimate  $Q$ :

$$Q = \frac{(f_a/f_b)^2 + 1}{(f_a/f_b)^2 - 1}$$

where  $f_a$  is the frequency above resonance where the real part of compliance or inertance reaches a peak and  $f_b$  is the frequency below resonance, where the real part of the compliance or inertance reaches a peak of opposite sign. The latter method has proven most useful for the present transfer functions which are computed in real and imaginary form directly.

Mode Shape: The mode shape is estimated by measuring quadrature (imaginary) response values from all of the measured inertance or compliance functions at each resonant frequency. The  $N$  mode shapes are normalized with respect to the maximum value of the quadrature component at each of the  $N$  natural frequencies.

These techniques were applied to the transfer functions of Figs. 2.5-15 and 2.5-16 to yield the mode shapes shown in Fig. 2.5-17 at the frequencies tabulated in Table 2.5-3. Since the test points were constrained to lie on the vertical centerline, only the vertical mode shapes are defined, and it is not possible to distinguish between symmetric horizontal mode shapes (odd wave numbers). From Fig. 2.5-17 it is apparent that the first few modes have been detected but the precise shape of higher modes is not clearly defined. Modes 1, 2, and 4 are the (1,1), (2,1), and (3,1) shapes, while modes 3 and 5 involve

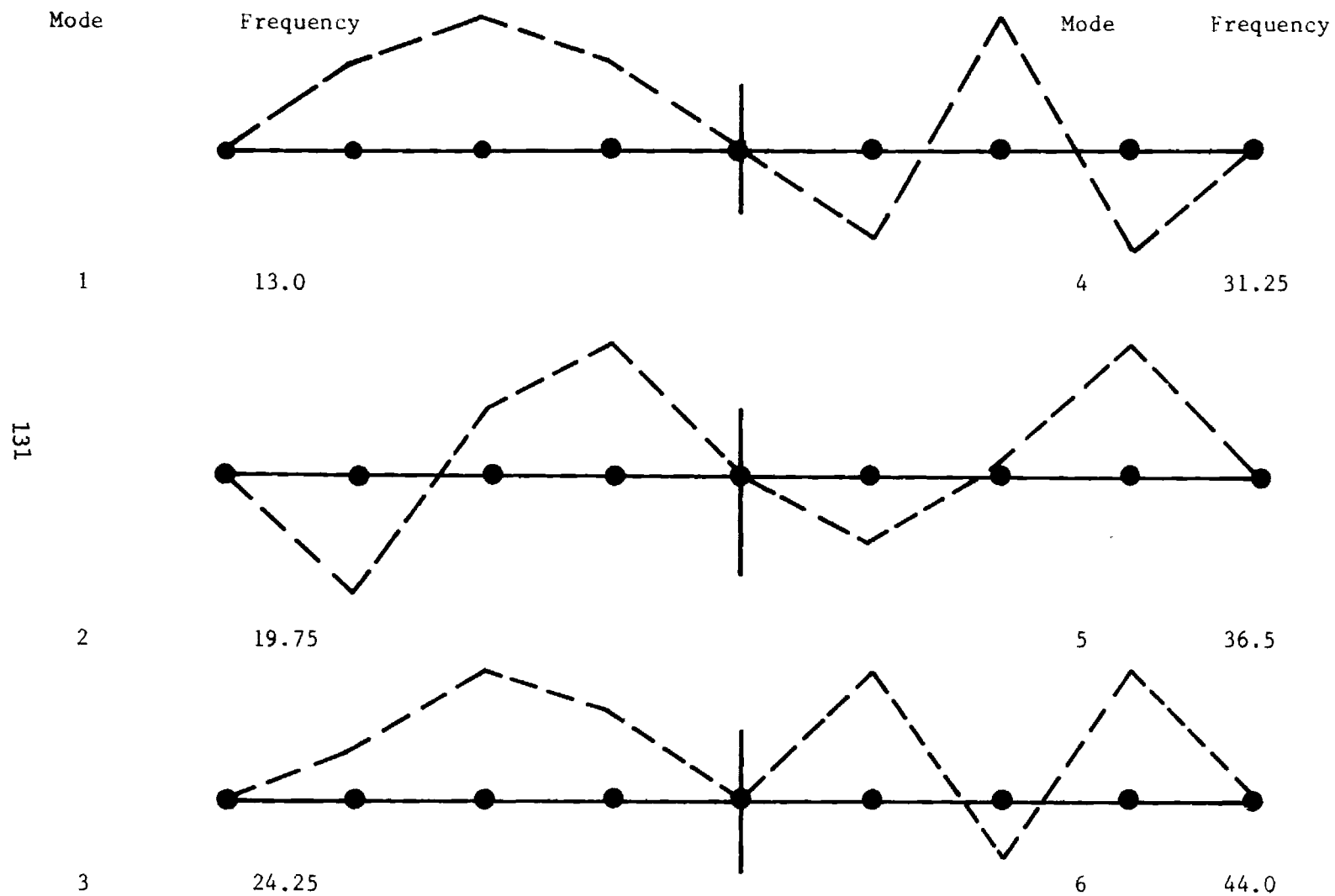


Figure 2.5-17. Mode Shapes



higher but undefined horizontal modes. The data from Table 2.5-3 shows that the natural frequencies of the window system are substantially higher than the bandwidth over which most of the wind power appears. There appears to be some evidence, however, from certain of the wind pressure psd and coherence plots, that in a few critical locations or situations, the (1,1) mode may be mildly excited.

The transfer function data show several other interesting features. There appears to be a strong mode at about 8.5 Hz that would seem at first to be the (1,1) fundamental. Closer study, however, indicates that this mode has a node at the midpoint and appears to be a type of rocking mode associated with the frame and adjoining windows. In addition, there appear to be a number of other unidentified modes with very small amplitudes as evidenced by the minor peaks between the major resonances. These are most likely associated with other elements in the adjoining cladding and were disregarded here. The damping coefficients shown in Table 2.5-3 are somewhat larger than the 0.5 to 1% usually associated with material damping alone. The 2 - 5% values reflect the additional viscoelastic behavior provided by the rubber glazing gaskets at the boundary. The largest values are associated with the lower modes.

The transfer function in Fig. 2.5-16 was obtained for single point (midpoint) on the outer pane by using the LDM displacement signal. As shown in Table 2.5-4, the modal frequencies generally agree with those for the inner pane; however, the corresponding mode 1 is not present. Instead, it is replaced by a strong resonance at the 8.5 Hz value previously associated with a generalized rocking mode.

Table 2.5-3 Longitudinal Modes, Window 28N4, Inner Plate

Mode #	Resonant Freq. $H_z$	Damping Coeff. $\zeta$	Mode Shape i.e. $n =$
1	13.0	.045	1
2	19.75	.028	2
3	24.0	.01	1
4	31.25	.02	3
5	36.5	---	2
6	44.0	.01	3

Table 2.5-4 Resonant Frequencies of Outer Glass Plate

Mode #	Frequency, $f_n$
1	----
2	20.75
3	31.0
4	36.75
5	43.75

Similar results were obtained for other windows examined but with variations in the associated frequencies. This is most clearly shown in Table 2.5-5 which is a tabulation of the (2,1) mode frequencies for 12 adjoining windows on the north side of the 28th floor. The range from 18.25 to 21.25 Hz represents a 15% variation in behavior which from the distribution in the table does not appear to be correlated with location.

In general the field tests to determine the dynamic response were troublesome to perform. The outer pane was accessible to measurement only with the LDM instrument which proved accurate in operation but somewhat difficult to move over the window area. In addition, the windows under study were clearly a part of the overall cladding and the measured response included certain modes associated with more generalized motion. The primary objective of this phase of the work was to ascertain, in general, what the in-situ modes looked like and their relation to the measured wind power. More detailed modal studies are described in the lab tests in Chapter 4.

Table 2.5-5. Variation in Resonant Frequency (2, 1) Mode

Window	Frequency, Hz
N1	18.5
N2	19.25
N3	19.25
N4	19.75
N5	18.25
N6	19.0
N7	18.5
N8	19.25
N9	18.5
N10	18.25
N11	21.25
N12	21.25

### 3.0 ANALYTICAL INVESTIGATIONS

#### 3.0(a) Need for Discrete Element Model

The laboratory window test facility described in this report can be used to test a number of glass cladding designs to determine their adequacy in carrying actual static and dynamic loads in the field. Naturally, the range of panel sizes and configurations as well as the loading types that can be applied in the laboratory are limited for reasons of economy. For example, the present test rig cannot be used to apply in-plane racking distortions to the frame-panel structure (although such modification to the test rig is possible), and the range of full-scale panels that can be tested is limited by the physical size and performance characteristics of the rig. Therefore, a computer model which was capable of accurately reproducing actual lab tests as well as permitting a broader range of panel sizes and support conditions and loading types to be evaluated was needed. To meet these requirements, a discrete element model of the frame, gasket, and glass panel was constructed using the finite element method. All basic components of the actual window installation are present in the analytical model, and complex static and dynamic loading conditions not possible in the laboratory can be handled with relative ease. In addition, the window model presented herein can be expanded to represent an entire glass curtain wall in a modern high-

rise building. The window model is sufficiently general that it can be directly incorporated into existing three-dimensional analytical models for static and dynamic analysis of multistory buildings. This extension of the present work would permit the determination of the interaction forces between frame and cladding for a wide range of loading conditions. The work described in this chapter constitutes the first step in an analytical study of cladding-structure interaction in highrise buildings.

### 3.0(b) Objectives

The basic objectives guiding this portion of the overall research program are as follows:

1. To develop a discrete element analytical model which is an accurate representation of the actual frame-panel structure;
2. To develop a FORTRAN computer program which implements the analytical model and permits the performance of a wide range of window designs of varying size, properties, and support conditions, and with static or dynamic applied loadings to be evaluated;
3. And, to conduct sensitivity studies to determine key parameters affecting the static and dynamic response of the total window-frame system.

The basic analytical model developed to satisfy objective 1. above is described in the next section. Model properties and support conditions were adjusted until both static displacements and vibration

frequencies for the analytical model were in good agreement with laboratory findings. The model was then considered to be a reasonably accurate representation of the actual structure which could be used in later sensitivity studies.

A FORTRAN computer program, WINDOW, based upon the finite element displacement method was developed to meet objective 2. above. A macro-flow chart of the program, a detailed user's guide, a listing of the program and several example problems illustrating its use for both static and dynamic loadings are presented in this chapter and in Appendix C.

To meet the final objective, basic properties of the frame-panel model were varied and the structure response computed for a number of representative loadings. Results are presented in the form of nondimensional curves which show the influence of system components on window response.

### 3.1 Analytical Model

#### 3.1(a) Description of Discrete Element Model

The discrete element model developed for this study employs a combination of two-dimensional plane stress and plate bending finite elements to represent the glass panel, and one-dimensional space frame elements to model the surrounding framework. Translational and rotational linear springs are used to represent the gasket on the panel boundary, and frame and panel are spring-connected at panel node points as shown in Figure 3.1-1. Nonrigid connections at frame corners are also included to account for frame connection rotational flexibility.

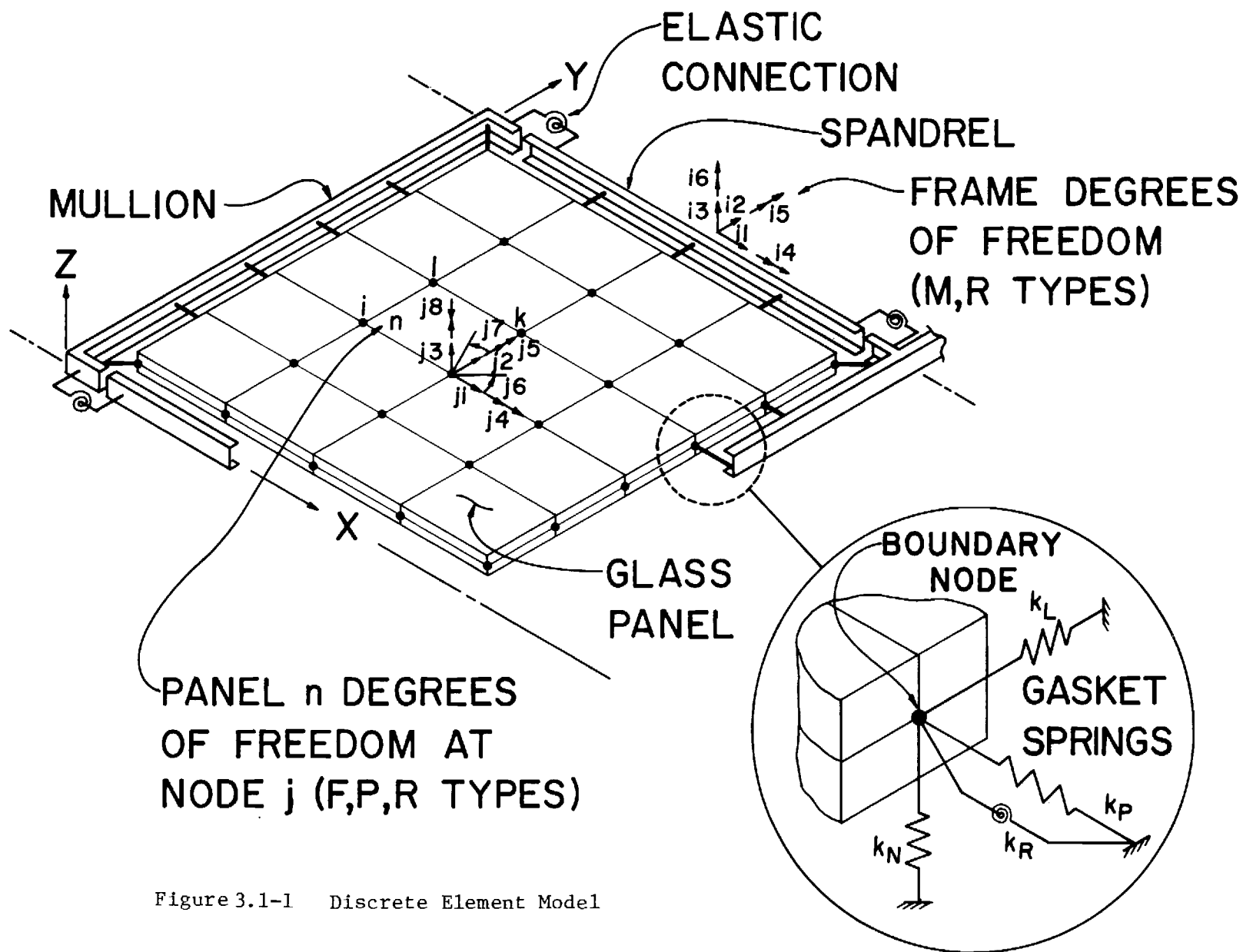


Figure 3.1-1 Discrete Element Model



Both in-plane and out-of-plane effects can be handled with this model, but the two are assumed uncoupled. Applied loads or racking displacements of the frame in the plane of the panel introduce distortions into the panel which are determined from the plane stress finite elements, while loads and frame movements normal to the panel cause bending distortions which activate the plate bending elements in the model. Deterministic planar and normal loadings may be applied simultaneously but coupled bending and stretching are not permitted since the model is assumed to be both geometrically and materially linear.

The stiffness method was used to assemble consistent stiffness and mass descriptions of the structure. Condensation of both stiffness and mass matrices was employed to eliminate extraneous degrees of freedom from the model thereby reducing the number of equations to be solved. A special purpose computer program was written to assemble the model, reduce the number of equations, and print and plot the system response for both static and time-dependent loadings.

Finite Elements: The plane stress and plate bending finite elements are superposed to form a single rectangular finite element, with eight degrees of freedom per node, referred to as the basic element. A refined plane stress element (88) and a conforming plate bending element (89), each having four degrees of freedom per node, were chosen because of their excellent past performance in modeling the static and dynamic behavior of plate-like structures (90). The geometry, nodal numbering scheme, and local coordinate system for the two elements are shown in Figure 3.1-2.

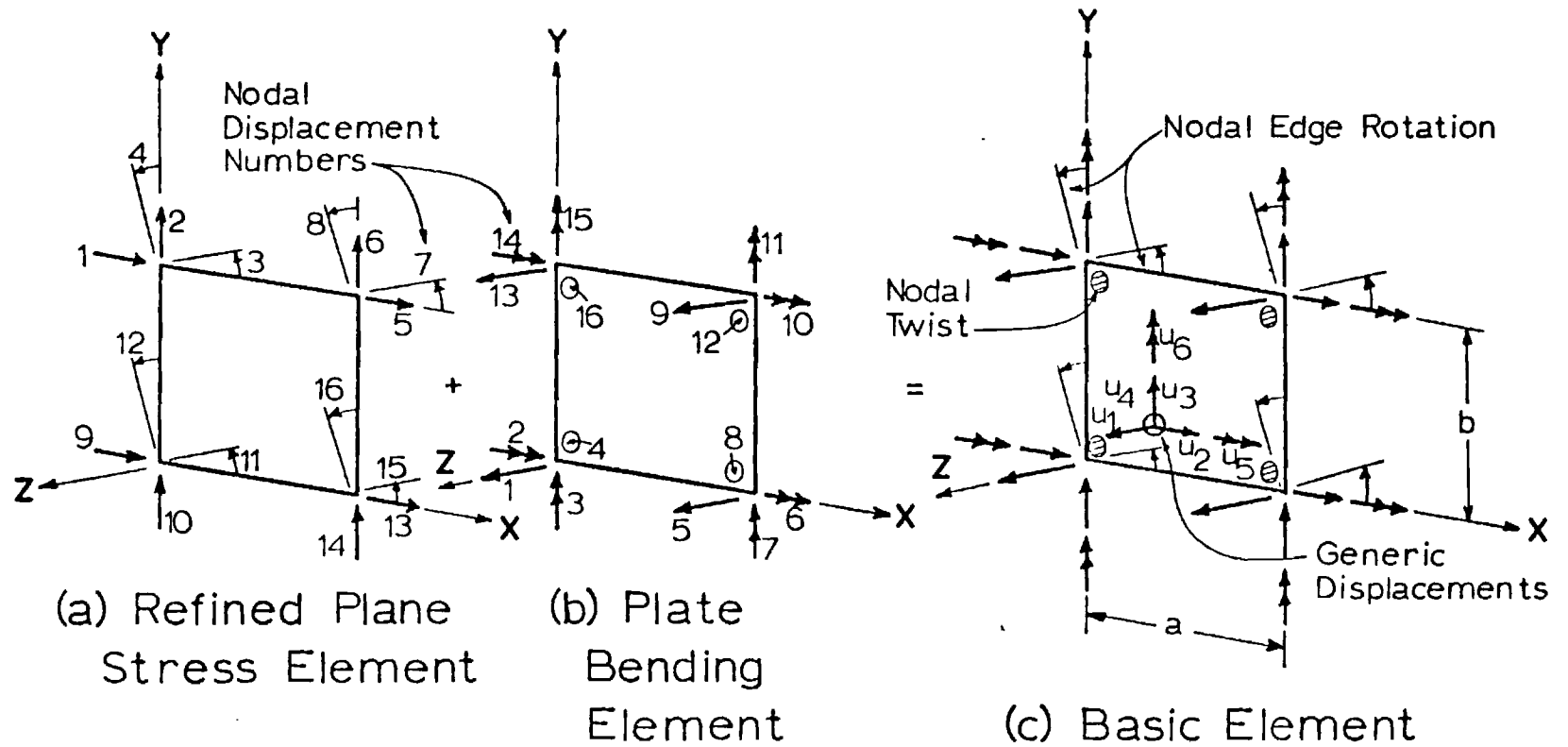


Figure 3.1-2 Formation of Basic Element

The generic displacements  $\tilde{u}$  at any point in the basic element are expressed in terms of the nodal displacements  $\tilde{q}$  for the element by the matrix equation:

$$\tilde{u} = T_1 \tilde{q} \quad (3.1-1)$$

where  $\tilde{u}' = \{u_1 \ u_2 \ u_3 \ u_4 \ u_5 \ u_6\}$  1 x 6

$$\tilde{q}' = \{q_1 \ q_2 \ \dots \ q_{32}\}$$
 1 x 32

and  $T_1$  is the matrix of shape functions for the basic element, assembled from the shape function arrays of the plane stress and plate bending finite elements (90). Generic displacements  $u_1$ ,  $u_4$ ,  $u_5$ , and  $u_6$  and nodal displacements  $q_1, \dots, q_{16}$  pertain to the plate bending finite element; while generic displacements  $u_2$  and  $u_3$  and nodal displacements  $q_{17}, \dots, q_{32}$  are associated with the plane stress element. Generic displacements  $u_4$ ,  $u_5$ , and  $u_6$  are dependent on  $u_1$  as follows:

$$u_4 = -\frac{\partial^2 u_1}{\partial x \partial y}, \quad u_5 = \frac{\partial u_1}{\partial y}, \quad u_6 = -\frac{\partial u_1}{\partial x} \quad (3.1-2)$$

where  $u_4$  represents a twist curvature and  $u_5$  and  $u_6$  represent rotations out of the plane of the plate bending element. For the plane stress element, rotations of the element edges parallel to the x and y axes are expressed as partial derivatives  $\frac{\partial u_3}{\partial x}$  and  $-\frac{\partial u_2}{\partial y}$ , respectively, evaluated at the element corners.

The plate bending element is both complete and conforming. A complete cubic polynomial plus six additional terms are employed in the derivation of the 16 x 16 element stiffness matrix  $\tilde{S}_{PB}$  for this element based on the principle of minimum potential energy. The 16 x 16 element stiffness matrix  $\tilde{S}_{PS}$  for the refined plane stress element

is obtained in the same manner. A linear-cubic displacement function is used, however, and the refined plane stress element is compatible with line elements in bending in frame and shear-wall models (88).

As a final step in the assembly of the basic element, a 32 x 32 element stiffness matrix  $\tilde{S}_E$  is constructed from arrays  $\tilde{S}_{PB}$  and  $\tilde{S}_{PS}$  as follows:

$$\tilde{S}_E = \begin{bmatrix} \tilde{S}_{PB} & 0 \\ 0 & \tilde{S}_{PS} \end{bmatrix} \quad (3.1-3)$$

$\begin{matrix} 16 \times 16 & & \\ & 0 & \\ & & 16 \times 16 \end{matrix}$

32x32

Standard matrix operations were used to rearrange  $\tilde{S}_E$  to obtain the counterclockwise ordering of node numbers  $k$ ,  $j$ ,  $k$ , and  $l$  and degree of freedom ordering at any node  $j$  as shown in Figure 3.1-1. In particular, a combined element stiffness matrix  $\tilde{S}_C$  for the 32 degree of freedom element in Figure 3.1-2 can be obtained from  $\tilde{S}_E$  as follows:

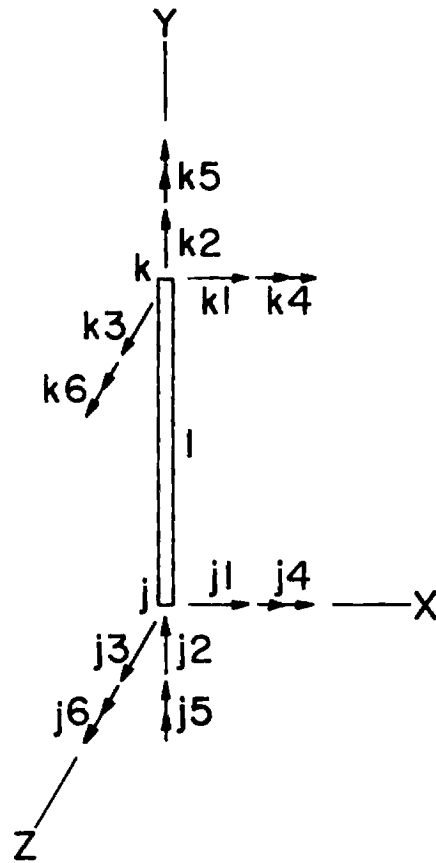
$$\tilde{S}_C = \tilde{T}' \tilde{S}_E \tilde{T} \quad (3.1-4)$$

where  $\tilde{T}$  is a matrix which correlates degree of freedom numbers in Figures 3.1-1 and 3.1-2 and  $\tilde{T}'$  is the transpose of  $\tilde{T}$ . Arrays  $\tilde{S}_C$  for each element in the panel are assembled to form the panel stiffness matrix.

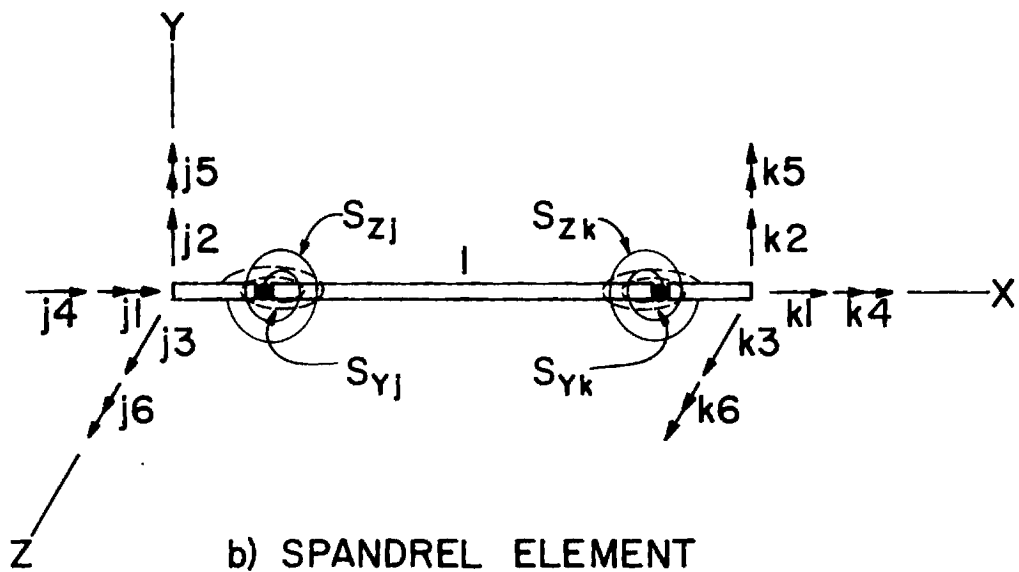
Frame Elements: The frame surrounding the panel is divided into mullion and spandrel elements each having six degrees of freedom per node. Mullion elements are space frame elements with local member axes parallel to the structure y axis while spandrel elements are space frame elements parallel to the structure x axis (Figure 3.1-3). The same number of frame and panel elements and nodes must be selected along each structure boundary so that frame and panel nodes correspond. However, the actual number of elements used will depend upon the finite element mesh employed.

The rotational connection stiffness of spandrel elements may be selected arbitrarily at the two ends of the member element. Connection stiffnesses at node j are denoted  $Sy_j$  and  $Sz_j$  for stiffness about the y and z axes, respectively; similarly stiffnesses at node k are  $Sy_k$  and  $Sz_k$ . Ordinarily, only spandrel element connections at mullion-spandrel junctures would be treated as nonrigid to account for framework flexibility at the joints; connection stiffnesses at all interior spandrel elements, with or without nonrigid connections, can be developed and assembled with mullion elements to form the window framework.

Spring Elements: The neoprene gasket which holds the glass panel in its frame was modeled as a series of discrete linear elastic springs joining adjacent frame and panel nodes along the boundary. Four basic springs are used as shown in Figure 3.1-1. Axial springs model the resistance of the gasket to motions in the plane of the panel, either perpendicular ( $k_p$ ) or parallel to ( $k_L$ ) a panel edge. Another axial spring represents the gasket's resistance to boundary displacements normal to the panel ( $k_N$ ), while a rotational spring ( $k_R$ ) accounts for



a) MULLION ELEMENT



b) SPANDREL ELEMENT

Figure 3.1-3 Frame Elements

gasket resistance to relative rotation between frame and panel on the boundary. All four spring constants are expressed as stiffness per unit length of panel boundary and must be multiplied by the tributary panel length to obtain the actual stiffness terms added to attached panel and frame nodes. Numerical values for unit spring constants  $k_P$ ,  $k_L$ ,  $k_N$ , and  $k_R$  were determined using simple laboratory static tests to be described later.

Mass and Damping Models: If dynamic behavior of the system is of interest, mass and damping properties must be specified for each element of the structure. The formation of a mass model for each element parallels the development of element stiffness matrices discussed above. A mass matrix for the basic finite element is constructed from element mass matrices for the plane stress and plate bending elements as follows:

$$\tilde{M}_E = \begin{bmatrix} \tilde{M}_{PB} & 0 \\ 0 & \tilde{M}_{PS} \end{bmatrix} \quad (3.1-5)$$

$\begin{matrix} 16 \times 16 & & \\ & 16 \times 16 & \end{matrix}$

32x32

If a consistent mass formulation is used, arrays  $\tilde{M}_{PB}$  and  $\tilde{M}_{PS}$  will be filled matrices. An alternative to consistent mass (CM) is assembled lumped mass (ALM) in which the mass of the element is lumped at nodal translational degrees of freedom. With ALM, arrays  $\tilde{M}_{PB}$  and  $\tilde{M}_{PS}$  are diagonal matrices; only entries (1,1), (5,5), (9,9), and (13,13) of  $\tilde{M}_{PB}$  and (1,1), (2,2), (5,5), (6,6), (9,9), (10,10), (13,13), and (14,14) of  $\tilde{M}_{PS}$  are non-zero for ALM, each term having the value  $1/4 \rho a b t$  where

a, b, t are element dimensions and  $\rho$  is the mass density of the material. (See Figure 3.1-4.) The ALM formulation has been shown to be superior to CM in determination of vibration frequencies for certain structures (90). In either case, element mass matrix  $\tilde{M}_E$  must be rearranged to produce  $\tilde{M}_C$  as in equation (3.1-4). Arrays  $\tilde{M}_C$  are then assembled for each panel element to form the panel mass matrix.

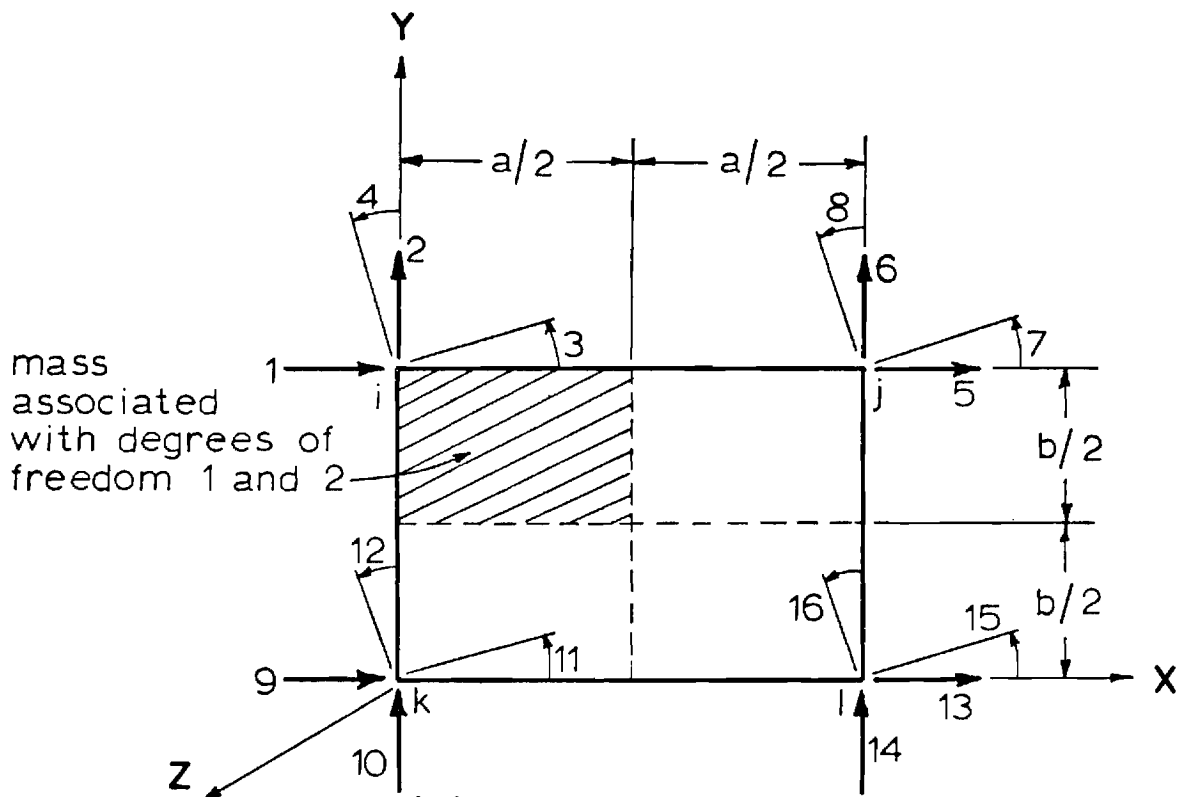
Consistent or assembled lumped mass matrices can be developed for frame elements in the same manner as for panel finite elements. However, linear springs representing the gaskets are assumed massless.

Element stiffness and mass matrices for panel finite elements, frame member elements, and spring elements are finally assembled, using procedures described below, to form overall structure stiffness and mass matrices for the entire window-frame structure. A structure damping matrix is not developed in explicit form; rather simple modal viscous damping was used with the fraction of critical viscous damping arbitrarily specified in each normal mode of vibration.

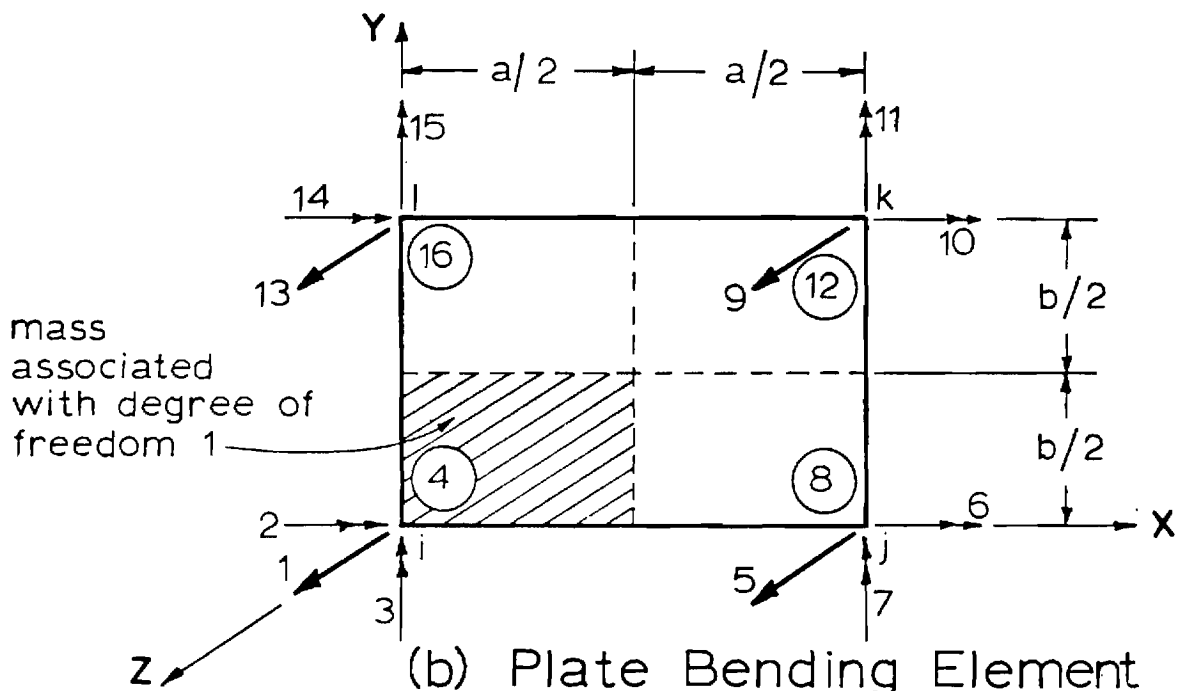
### 3.1(b) Structural Assemblage and Condensation

Panel, frame, and spring elements are assembled to form stiffness and mass properties of the total structure model. Extraneous degrees of freedom on the panel and frame are eliminated using a two-stage condensation procedure. The primary motivation for this approach was to construct an efficient model for dynamic analysis with as few dynamic degrees of freedom as practicable. In the first stage nonessential panel (or P type) degrees of freedom are condensed out, while in the second stage unnecessary frame (or M type) degrees of freedom are eliminated. Stiffness and mass arrays are partitioned into subarrays to conserve computer storage and to





(a) Plane Stress Element



(b) Plate Bending Element

Figure 3.1-4 Assembled Lumped Mass (ALM)

expedite the assembly and condensation procedures.

Displacement Types: Structure degrees of freedom are categorized as F, P, M, and R displacement types as a basis for partitioning stiffness and mass matrices and for orderly elimination of unessential degrees of freedom. F displacement types are those panel degrees of freedom which are to remain at the end of the condensation operations as master degrees of freedom. Any number and distribution of panel displacements may be selected as F types but points of load application or mass concentration should be chosen and corresponding degrees of freedom designated as F types. For example, nodal degrees of freedom of type  $j_3$  (Figure 3.1-1) at all unrestrained panel nodes could be selected as F types if the response normal to the plane of the panel were of interest. The remaining unrestrained panel degrees of freedom are designated as P type displacements to be eliminated in the first stage of condensation. All unrestrained frame displacement coordinates called M types are removed in the second condensation operation. Finally, R types refer to all restrained panel and frame displacement coordinates, which together with all F types, remain at the end of the condensations. The F and R types, then, constitute a reduced set of degrees of freedom which can be used to describe the response of particular points of interest on the panel. Applied loadings corresponding to the F degrees of freedom and support displacements at R displacement coordinates are acceptable loading conditions in the reduced structure model. Furthermore, static and dynamic response at eliminated (P,M) displacement coordinates may still be determined through a backsubstitution operation once the response at the master degrees of freedom is known.

The undamped equations of motion for the total window-frame model can be written as follows:

$$\tilde{M}_J \ddot{\tilde{D}}_J + \tilde{S}_J \tilde{D}_J = \tilde{A}_J \quad (3.1-6)$$

or in partitioned form:

$$\begin{bmatrix} \tilde{M}_{FF} & \tilde{M}_{FP} & \tilde{M}_{FM} & \tilde{M}_{FR} \\ \tilde{M}_{PF} & \tilde{M}_{PP} & \tilde{M}_{PM} & \tilde{M}_{PR} \\ \tilde{M}_{MF} & \tilde{M}_{MP} & \tilde{M}_{MM} & \tilde{M}_{MR} \\ \tilde{M}_{RF} & \tilde{M}_{RP} & \tilde{M}_{RM} & \tilde{M}_{RR} \end{bmatrix} \begin{Bmatrix} \tilde{D}_F \\ \tilde{D}_P \\ \tilde{D}_M \\ \tilde{D}_R \end{Bmatrix} + \begin{bmatrix} \tilde{S}_{FF} & \tilde{S}_{FP} & \tilde{S}_{FM} & \tilde{S}_{FR} \\ \tilde{S}_{PF} & \tilde{S}_{PP} & \tilde{S}_{PM} & \tilde{S}_{PR} \\ \tilde{S}_{MF} & \tilde{S}_{MP} & \tilde{S}_{MM} & \tilde{S}_{MR} \\ \tilde{S}_{RF} & \tilde{S}_{RP} & \tilde{S}_{RM} & \tilde{S}_{RR} \end{bmatrix} \begin{Bmatrix} \tilde{D}_F \\ \tilde{D}_P \\ \tilde{D}_M \\ \tilde{D}_R \end{Bmatrix} = \begin{Bmatrix} \tilde{A}_F \\ \tilde{A}_P \\ \tilde{A}_M \\ \tilde{A}_R \end{Bmatrix} \quad (3.1-7)$$

where F, P, M, and R subscripts refer to the displacement types described above. Mass and stiffness matrices  $\tilde{M}_J$  and  $\tilde{S}_J$  are symmetric so only those subarrays on and above the diagonals are actually assembled.

Assemblage: Selected submatrices of structure stiffness matrix  $\tilde{S}_J$  are assembled and condensed first, then a similar operation is performed to build and condense mass matrix  $\tilde{M}_J$ . To begin the procedure panel finite elements are processed one at a time, element stiffness matrices  $\tilde{S}_C$  (Equation 3.1-4) formed, and contributions of  $\tilde{S}_C$  to submatrices in  $\tilde{S}_J$  recorded. If successive panel elements are identical,  $\tilde{S}_C$  need not be regenerated until an element of different size is encountered. Next, mullion element and spandrel element (with or without nonrigid connections) stiffness matrices are formed for each framing element and are used to augment submatrices of  $\tilde{S}_J$ . Finally, spring elements add stiffness terms at panel and frame degrees of freedom along the boundary of the model to complete the assemblage of matrix  $\tilde{S}_J$ . Following condensation of  $\tilde{S}_J$ , mass matrix  $\tilde{M}_J$  is formed.

Panel element degrees of freedom may be F, P, or R types. Consequently, matrices  $\tilde{S}_C$  and  $\tilde{M}_C$  for each element add stiffness and mass terms

to FF, FP, FR, PP, PR, and RR subarrays in Equation 3.1-7. Frame elements form MM and MR subarrays and augment the RR submatrix. Spring elements tie the panel and frame together to form a stable model; spring stiffness terms are used to form FM and PM stiffness submatrices and to add stiffness terms to FF, PP, MM, RR, FR, PR, and MR subarrays. Springs are assumed massless and do not contribute to  $\tilde{M}_J$ . If the consistent mass (CM) formulation is used, all ten of the mass submatrices which are formed will have nonzero entries, while assembled lumped mass (ALM) will result in diagonal FF, PP, MM, and RR subarrays and all remaining entries in  $\tilde{M}_J$  are zeros.

Condensation: Elimination of displacement coordinates which have very little or no mass associated with them or whose response is not of primary interest produces a more efficient model for dynamic analysis. The two-stage reduction procedure employed here begins with static condensation of  $\tilde{S}_J$  to eliminate first P and then M type displacements. The static equilibrium equations,

$$\tilde{S}_J \tilde{D}_J = \tilde{A}_J \quad (3.1-8)$$

with the partitioned forms of arrays  $\tilde{S}_J$ ,  $\tilde{D}_J$ , and  $\tilde{A}_J$  shown in Equation (3.1-7), are reduced to

$$\tilde{S}_J^* \tilde{D}_{J_1} = \tilde{A}_{J_1} \quad (3.1-9)$$

where

$$\tilde{S}_J^* = \begin{bmatrix} \tilde{S}_{FF}^* & \tilde{S}_{FM}^* & \tilde{S}_{FR}^* \\ \tilde{S}_{MF}^* & \tilde{S}_{MM}^* & \tilde{S}_{MR}^* \\ \tilde{S}_{RF}^* & \tilde{S}_{RM}^* & \tilde{S}_{RR}^* \end{bmatrix} \quad (3.1-10)$$

and

$$\underset{\sim}{D}_{J_1} = \begin{Bmatrix} \underset{\sim}{D}_F \\ \underset{\sim}{D}_M \\ \underset{\sim}{D}_R \end{Bmatrix}, \quad \underset{\sim}{A}_{J_1} = \begin{Bmatrix} \underset{\sim}{A}_F \\ 0 \\ \underset{\sim}{A}_R \end{Bmatrix} \quad (3.1-11)$$

and asterisk superscripts denote condensed arrays. Vectors  $\underset{\sim}{A}_P$  and  $\underset{\sim}{A}_M$ , representing applied forces at P and M displacement coordinates, are assumed zero for simplicity. Condensed matrix  $\underset{\sim}{S}_J^*$  is obtained by solving for  $\underset{\sim}{D}_P$  in the first of equations (3.1-8) as

$$\underset{\sim}{D}_P = \underset{\sim}{S}_{PP}^{-1} (-\underset{\sim}{S}_{FP}^T \underset{\sim}{D}_F - \underset{\sim}{S}_{PM} \underset{\sim}{D}_M - \underset{\sim}{S}_{PR} \underset{\sim}{D}_R) \quad (3.1-12)$$

and substituting this expression into the remaining equations. Only condensed matrices on and above the diagonal of  $\underset{\sim}{S}_J^*$  are actually computed due to symmetry. The form of  $\underset{\sim}{S}_{FF}^*$ , for example, is

$$\underset{\sim}{S}_{FF}^* = \underset{\sim}{S}_{FF} - \underset{\sim}{S}_{FP} \underset{\sim}{S}_{PP}^{-1} \underset{\sim}{S}_{FP}^T \quad (3.1-13)$$

In the next stage of condensation, static equilibrium equations (3.1-9) for F, M, and R degrees of freedom are reduced to final form, retaining only F and R types, yielding

$$\underset{\sim}{S}_J^{**} \underset{\sim}{D}_{J_2} = \underset{\sim}{A}_{J_2} \quad (3.1-14)$$

where

$$\underset{\sim}{S}_J^{**} = \begin{bmatrix} \underset{\sim}{S}_{FF}^{**} & \underset{\sim}{S}_{FR}^{**} \\ \underset{\sim}{S}_{RF}^{**} & \underset{\sim}{S}_{RR}^{**} \end{bmatrix} \quad (3.1-15)$$

and

$$\underset{\sim}{D}_{J_2} = \begin{Bmatrix} \underset{\sim}{D}_F \\ \underset{\sim}{D}_R \end{Bmatrix}, \quad \underset{\sim}{A}_{J_2} = \begin{Bmatrix} \underset{\sim}{A}_F \\ \underset{\sim}{A}_R \end{Bmatrix} \quad (3.1-16)$$

Here, an expression for displacement vector  $\underline{D}_M$  was obtained by solving the second of equations (3.1-9) with the result

$$\underline{D}_M = \underline{S}_{MM}^{*-1} (-\underline{S}_{FM}^* \underline{D}_F - \underline{S}_{MR}^* \underline{D}_R) \quad (3.1-17)$$

Substituting this expression into the remaining equations produced Equation (3.1-14). The forms of the  $\underline{S}_J^{**}$  submatrices are:

$$\underline{S}_{FF}^{**} = \underline{S}_{FF}^* - \underline{S}_{FM}^* \underline{S}_{MM}^{*-1} \underline{S}_{FM}^* \quad (3.1-18)$$

$$\underline{S}_{FR}^{**} = \underline{S}_{RF}^{**} = \underline{S}_{FR}^* - \underline{S}_{FM}^* \underline{S}_{MM}^{*-1} \underline{S}_{MR}^* \quad (3.1-19)$$

$$\underline{S}_{RR}^{**} = \underline{S}_{RR}^* - \underline{S}_{MR}^* \underline{S}_{MM}^{*-1} \underline{S}_{MR}^* \quad (3.1-20)$$

where array  $\underline{S}_{FF}^{**}$ , for example, represents holding forces at F degrees of freedom due to unit displacements at F degrees of freedom in a partially restrained structure with P and M displacement types allowed to occur freely.

With the condensation of  $\underline{S}_J^*$  complete, mass matrix  $\underline{M}_J$  can be reduced to  $\underline{M}_J^*$  and then to  $\underline{M}_J^{**}$  in parallel fashion, removing first P and then M displacement terms. The equations of motion after elimination of P types have the form

$$\underline{M}_J^* \ddot{\underline{D}}_{J_1} + \underline{S}_J^* \underline{D}_{J_1} = \underline{A}_{J_1} \quad (3.1-21)$$

where  $\underline{M}_J^*$  is similar in form to  $\underline{S}_J^*$  above. After elimination of M types, the final reduced set of undamped equations of motion are

$$\underline{M}_J^{**} \ddot{\underline{D}}_{J_2} + \underline{S}_J^{**} \underline{D}_{J_2} = \underline{A}_{J_2} \quad (3.1-22)$$

where

$$\underline{M}_J^{**} = \begin{bmatrix} \underline{M}_{FF}^{**} & \underline{M}_{FR}^{**} \\ \underline{M}_{RF}^{**} & \underline{M}_{RR}^{**} \end{bmatrix} \quad (3.1-23)$$

Condensation of  $\tilde{M}_J$  to  $\tilde{M}_J^*$  is accomplished by equating the virtual work of the inertia forces in equations (3.1-6) and (3.1-21), as discussed by Guyan (91), with the result

$$\tilde{M}_J^* = \tilde{T}_P' \tilde{M}_J \tilde{T}_P \quad (3.1-24)$$

where

$$\tilde{T}_P = \begin{bmatrix} \tilde{I} & \tilde{0} & \tilde{0} \\ -\tilde{S}_{PP}^{-1} \tilde{S}_{FP}' & -\tilde{S}_{PP}^{-1} \tilde{S}_{PM}' & -\tilde{S}_{PP}^{-1} \tilde{S}_{PR}' \\ \tilde{0} & \tilde{I} & \tilde{0} \\ \tilde{0} & \tilde{0} & \tilde{I} \end{bmatrix} \quad (3.1-25)$$

Stiffness array products in  $\tilde{T}_P$  are available from condensation of  $\tilde{S}_J$  to  $\tilde{S}_J^*$ , and  $\tilde{I}$  and  $\tilde{0}$  are identity and null matrices, respectively, of appropriate size. In the second condensation stage, inertia terms in equations (3.1-21) and (3.1-22) are made energy equivalent with the result

$$\tilde{M}_J^{**} = \tilde{T}_M' \tilde{M}_J^* \tilde{T}_M \quad (3.1-26)$$

where

$$\tilde{T}_M = \begin{bmatrix} \tilde{I} & \tilde{0} \\ -\tilde{S}_{MM}^{*-1} \tilde{S}_{FM}' & -\tilde{S}_{MM}^{*-1} \tilde{S}_{MR}' \\ \tilde{0} & \tilde{I} \end{bmatrix} \quad (3.1-27)$$

Stiffness array products in  $\tilde{T}_M$  were generated during the reduction of  $\tilde{S}_J^*$  to  $\tilde{S}_J^{**}$  (See equations 3.1-18, 3.1-19.). If the transformation in equation (3.1-26) is carried out, the explicit expression for  $\tilde{M}_{FF}^{**}$ , for example, is:

$$\tilde{M}_{FF}^{**} = \tilde{M}_{FF}^* + \tilde{M}_{FM}^* (-\tilde{S}_{MM}^{*-1} \tilde{S}_{FM}'') + (-\tilde{S}_{MM}^{*-1} \tilde{S}_{FM}'')' \tilde{M}_{FM}^* + (-\tilde{S}_{MM}^{*-1} \tilde{S}_{FM}'')' \tilde{M}_{MM}^* (-\tilde{S}_{MM}^{*-1} \tilde{S}_{FM}'') \quad (3.1-28)$$

At the end of the assembly and condensation procedures, only arrays  $\tilde{S}_{JJ}^{**}$  and  $\tilde{M}_{JJ}^{**}$  remain reflecting the elastic and inertial force coupling, respectively, among selected master (F type) and restraint (R type) displacement coordinates. Using these arrays, static and dynamic analyses of the reduced structure model may now be performed.

### 3.1(c) Static Analysis Procedures

The basic static equilibrium equations for the condensed structure model were presented above in equations (3.1-14), (3.1-15) and (3.1-16). If the double asterisk superscripts (denoting two stages of condensation) are dropped to simplify the notation, the equations can be written as

$$\begin{bmatrix} \tilde{S}_{FF} & \tilde{S}_{FR} \\ \tilde{S}'_{FR} & \tilde{S}_{RR} \end{bmatrix} \begin{Bmatrix} \tilde{D}_F \\ \tilde{D}_R \end{Bmatrix} = \begin{Bmatrix} \tilde{A}_F \\ \tilde{A}_R \end{Bmatrix} \quad (3.1-29)$$

where  $\tilde{D}_F$  = vector of static nodal displacements at master degrees of freedom (F types),  $\tilde{D}_R$  = vector of specified restraint displacements at panel or frame supports,  $\tilde{A}_F$  = vector of applied nodal loads at F types, and  $\tilde{A}_R$  = vector of support reactions. As many master degrees of freedom on the panel as desired can be retained to characterize the load distribution and displaced configuration of the panel. Distributed loadings must be replaced by equivalent concentrated nodal loads  $\tilde{A}_F$  at F type displacement coordinates, and points of application of concentrated loads should have F degrees of freedom associated with them. Support displacements  $\tilde{D}_R$  can be specified if panel response to racking displacements is of interest.

Solution for Nodal Displacements: In the stiffness method of analysis, displacements at structure nodes are the basic unknowns. Displacements at master degrees of freedom  $\tilde{D}_F$  due to applied nodal loads  $\tilde{A}_F$  and specified



support motions  $\tilde{D}_R$  are obtained by solving the first of matrix equations (3.1-29) as follows:

$$\tilde{D}_F = \tilde{S}_{FF}^{-1}(\tilde{A}_F - \tilde{S}_{FR}\tilde{D}_R) \quad (3.1-30)$$

Then, the second of equations (3.1-29) yields the support reaction forces

$$\tilde{A}_R = \tilde{S}'_{FR}\tilde{D}_F + \tilde{S}_{RR}\tilde{D}_R \quad (3.1-31)$$

If static frame displacements at M type degrees of freedom are of interest, equation (3.1-17) can be solved to obtain displacement vector  $\tilde{D}_M$  in the backward substitution phase of the analysis. Finally, unknown panel displacements  $\tilde{D}_P$  are found from equation (3.1-12). With all frame and panel displacements known, frame member end forces and panel element stresses can be computed as discussed below.

Solution for Member Forces and Stresses: If the member stiffness matrix for any frame member i is referred to as  $\tilde{S}_i$ , the forces  $\tilde{A}_i$  at member ends, corresponding to the member degrees of freedom in Figure (3.1-3), can be computed as

$$\tilde{A}_i = \tilde{S}_i\tilde{D}_i \quad (3.1-32)$$

where  $\tilde{D}_i$  = vector of displacements at the j and k ends of member i assembled from vectors  $\tilde{D}_M$  and  $\tilde{D}_R$ .

Likewise, stresses may be computed within any panel finite element once nodal displacements are known for that element. Recall that the generic displacements  $\tilde{u}$  at any point within the element depend upon the nodal displacements  $\tilde{q}$  as follows:

$$\tilde{u} = \tilde{T}_1\tilde{q} \quad (3.1-1) \text{ repeated}$$

where  $\tilde{T}_1$  is the shape function matrix for the element and  $\tilde{q}$  is assembled

from displacement vectors  $\underline{D}_F$  and  $\underline{D}_P$  (equations (3.1-30) and (3.1-12)). Element generic and nodal displacements are related to element strains,

$$\underline{\epsilon} = \underline{T}_2 \underline{u} = \underline{T}_2 \underline{T}_1 \underline{q} = \underline{B} \underline{q} \quad (3.1-33)$$

and element strains are related to element stresses,

$$\underline{\sigma} = \underline{T}_3 \underline{\epsilon} = \underline{T}_3 \underline{T}_2 \underline{T}_1 \underline{q} = \underline{T}_3 \underline{B} \underline{q} \quad (3.1-34)$$

at any point within the element. Vectors  $\underline{\epsilon}$  and  $\underline{\sigma}$  are the strain and stress vectors, respectively, expressed as

$$\underline{\sigma} = \begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix}, \quad \underline{\epsilon} = \begin{Bmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{Bmatrix} \quad (3.1-35)$$

where  $\sigma$  and  $\epsilon$  represent normal stress and strain components, and  $\tau$  and  $\gamma$  shear stress and strain components in the x-y plane. If plate bending and plane stress effects are superimposed, strain component  $\epsilon_x$ , for example, may be computed in the following manner:

$$\begin{aligned} \epsilon_x &= \left( \frac{\partial u_2}{\partial x} \right)_{\text{plate bending}} + \left( \frac{\partial u_2}{\partial x} \right)_{\text{plane stress}} \\ &= -z \frac{\partial^2 u_1}{\partial x^2} + \frac{\partial u_2}{\partial x} \end{aligned} \quad (3.1-36)$$

where  $u_1$  and  $u_2$  are the element generic displacements shown in Figure 3.1-2.

Matrix  $\underline{T}_2$  above is a matrix of derivatives and  $\underline{B}$  is the strain-displacement matrix. Matrix  $\underline{T}_3$  is the stress-strain matrix expressed as

$$\underline{T}_3 = \frac{E}{1-\nu^2} \begin{bmatrix} 1 & & 0 \\ & 1 & 0 \\ 0 & 0 & \frac{1-\nu}{2} \end{bmatrix} \quad (3.1-37)$$

for two-dimensional plane stress problems, where  $E$  is the modulus of elasticity and  $\nu$  is Poisson's ratio for the element.

The basic approach outlined above for obtaining stresses or stress resultants in frame and panel members in a static problem may be used to compute stress-time histories in a dynamics problem. In this case, displacement vectors  $\underline{D}_F$ ,  $\underline{D}_P$ ,  $\underline{D}_M$ , and  $\underline{D}_R$  represent the displaced configuration of the structure at any selected time of interest during dynamic response, and are used to assemble vectors  $\underline{D}_i$  and  $\underline{q}$  in equations (3.1-32) and (3.1-1) as described in the next section.

### 3.1(d) Dynamic Analysis Procedures

The undamped equations of motion for master (F) and restrained (R) displacement coordinates in the condensed structure model were presented earlier in equation (3.1-22) which is repeated here in partitioned form

$$\begin{bmatrix} \underline{M}_{FF} & \underline{M}_{FR} \\ \underline{M}'_{FR} & \underline{M}_{RR} \end{bmatrix} \begin{Bmatrix} \ddot{\underline{D}}_F \\ \ddot{\underline{D}}_R \end{Bmatrix} + \begin{bmatrix} \underline{S}_{FF} & \underline{S}_{FR} \\ \underline{S}'_{FR} & \underline{S}_{RR} \end{bmatrix} \begin{Bmatrix} \underline{D}_F \\ \underline{D}_R \end{Bmatrix} = \begin{Bmatrix} \underline{A}_F \\ \underline{A}_R \end{Bmatrix} \quad (3.1-38)$$

with asterisk superscripts omitted for simplicity. If the first of these equations is considered and applied forces  $\underline{A}_F$  and support motions  $\ddot{\underline{D}}_R$  and  $\underline{D}_R$  are assumed to be zero, the equation of motion governing undamped free vibration of the structure results:

$$\underline{M}_{FF} \ddot{\underline{D}}_F + \underline{S}_{FF} \underline{D}_F = 0 \quad (3.1-39)$$

This equation characterizes the free vibration behavior of an ideal elastic structure which, if properly disturbed, will execute harmonic motion in a natural mode of vibration. Each natural mode  $i$  has a natural frequency  $p_i$  and a characteristic shape  $\underline{X}_{M_i}$ , with the number of frequencies

and mode shapes equal to the number of master degrees of freedom.

The solution of equation (3.1-39) is of the form

$$\underset{\sim}{D}_{F_i} = \underset{\sim}{X}_{M_i} \sin (p_i t + \phi_i) \quad (3.1-40)$$

where  $\underset{\sim}{D}_{F_i}$  is the nodal displacement vector for master degrees of freedom in the  $i$ -th mode and  $\phi_i$  is the  $i$ -th mode phase angle. Substitution of equation (3.1-40) into equation (3.1-39) yields

$$\underset{\sim}{S}_{FF} \underset{\sim}{X}_{M_i} = p_i^2 \underset{\sim}{M}_{FF} \underset{\sim}{X}_{M_i} \quad (3.1-41)$$

the algebraic eigenvalue problem in nonstandard form in which  $p_i^2$  is the  $i$ -th mode eigenvalue and  $\underset{\sim}{X}_{M_i}$  the  $i$ -th mode eigenvector. Solution of equation (3.1-41) gives frequencies  $p_i$  and modal amplitudes  $\underset{\sim}{X}_{M_i}$ .

Solution of Equations of Motion: The normal mode method is a widely used technique for solution of the equations of motion for linear systems. In this approach, the system displacement ( $\underset{\sim}{D}_F$ ) and acceleration ( $\ddot{\underset{\sim}{D}}_F$ ) response can be determined by transforming to normal coordinates where arrays  $\underset{\sim}{M}_{FF}$  and  $\underset{\sim}{S}_{FF}$  are diagonal. The multi-degree of freedom system is reduced to a set of equivalent single degree of freedom oscillators whose damped responses are readily determined. The structure response in the original coordinate system is obtained by a back transformation from normal coordinates. Finally, reactions  $\underset{\sim}{A}_R$  at supports for all times of interest can be determined from the second set of matrix equations (3.1-38).

It is the orthogonality property of the eigenvectors that permits the transformation of the equations of motion to normal coordinates. Orthogonality with respect to  $\underset{\sim}{M}_{FF}$  means that, for any two modes  $i$  and  $j$  with  $i \neq j$ ,

$$\tilde{X}_{M_j}^T \tilde{M}_{FF} \tilde{X}_{M_i} = 0 \quad (3.1-42)$$

When  $i = j$ ,

$$\tilde{X}_{M_j}^T \tilde{M}_{FF} \tilde{X}_{N_j} = Q_j \quad (3.1-43)$$

where  $Q_j$  is an arbitrary constant which can be set to unity by normalizing the eigenvectors with respect to the mass matrix such that

$$\tilde{X}_{N_j}^T \tilde{M}_{FF} \tilde{X}_{N_j} = 1 \quad (3.1-44)$$

Vector  $\tilde{X}_{N_j}$  denotes the normalized eigenvector for the  $j$ -th vibration mode. It follows that

$$\tilde{X}_{N_i}^T \tilde{M}_{FF} \tilde{X}_{N_j} = \delta_{ij} \quad (3.1-45)$$

and

$$\tilde{X}_{N_i}^T \tilde{S}_{FF} \tilde{X}_{N_j} = p_{ij}^2 \quad (3.1-46)$$

where  $I$  is the identity matrix and  $p^2$  is the spectral matrix containing terms  $p_{ij}^2$  along its diagonal. Array  $\tilde{X}_N = \{\tilde{X}_{N_1} \tilde{X}_{N_2} \dots \tilde{X}_{N_j} \dots \tilde{X}_{N_n}\}$  is the normalized modal matrix composed of the normalized eigenvectors for all  $n$  master degrees of freedom. Now, ignoring support motion terms in the first of equations (3.1-38) for the moment, orthogonality relationships (3.1-45) and (3.1-46) can be used to uncouple the equations of motion as follows:

$$\tilde{X}_{N_i}^T \tilde{M}_{FF} \tilde{X}_{N_j} \ddot{D}_N + \tilde{X}_{N_i}^T \tilde{S}_{FF} \tilde{X}_{N_j} D_N^{-1} = \tilde{X}_{N_i}^T A_F \quad (3.1-47)$$

or

$$\ddot{D}_N + p^2 D_N = F_N \quad (3.1-48)$$

where

$$\ddot{\underline{D}}_{\underline{N}} = \underline{X}_{\underline{N}}^{-1} \ddot{\underline{D}}_{\underline{F}} = \underline{X}_{\underline{N}}' \underline{M}_{\underline{FF}} \ddot{\underline{D}}_{\underline{F}} \quad (3.1-49a)$$

$$\underline{D}_{\underline{N}} = \underline{X}_{\underline{N}}^{-1} \underline{D}_{\underline{F}} = \underline{X}_{\underline{N}}' \underline{M}_{\underline{FF}} \underline{D}_{\underline{F}} \quad (3.1-49b)$$

$$\underline{F}_{\underline{N}} = \underline{X}_{\underline{N}}' \underline{A}_{\underline{F}} \quad (3.1-49c)$$

Simple modal damping may now be inserted into equations (3.1-48) to produce the  $i$ -th mode damped equation of motion

$$\ddot{D}_{N_i} + 2\gamma_i p_i \dot{D}_{N_i} + p_i^2 D_{N_i} = F_{N_i} \quad (3.1-50)$$

where  $\gamma_i$  is the fraction of critical damping in mode  $i$ . Solution of equation (3.1-50) provides the response in normal coordinates for each single degree of freedom oscillator. Finally, the system response in the original coordinates is obtained from equations (3.1-49a) and (3.1-49b) as

$$\ddot{\underline{D}}_{\underline{F}} = \underline{X}_{\underline{N}} \ddot{\underline{D}}_{\underline{N}} \quad (3.1-51a)$$

$$\underline{D}_{\underline{F}} = \underline{X}_{\underline{N}} \underline{D}_{\underline{N}} \quad (3.1-51b)$$

and support reactions  $\underline{A}_{\underline{R}}$  are calculated using equation (3.1-38) in which  $\underline{D}_{\underline{R}}$  and  $\ddot{\underline{D}}_{\underline{R}}$  are either specified or zero. Further details may be found in reference (92).

Types of Loadings: While the basic solution procedure outlined above is quite straightforward, several details regarding various types of system excitations are worthy of specific mention. The most elementary case is concerned with transient response of the system due to initial conditions of displacement and velocity alone with forcing functions  $\underline{A}_{\underline{F}}$  zero. In this case, vectors of initial displacements  $\underline{D}_0$  and velocities  $\dot{\underline{D}}_0$  at master

degrees of freedom are specified and are transformed to normal coordinates as follows:

$$\underset{\sim}{D}_{N_0} = \underset{\sim}{X}_N^{-1} \underset{\sim}{D}_0 \quad (3.1-52a)$$

$$\dot{\underset{\sim}{D}}_{N_0} = \underset{\sim}{X}_N^{-1} \dot{\underset{\sim}{D}}_0 \quad (3.1-52b)$$

where  $\underset{\sim}{D}_{N_0}$  and  $\dot{\underset{\sim}{D}}_{N_0}$  are the initial displacement and velocity vectors in normal coordinates. Equation (3.1-50), with  $F_{N_i}$  taken as zero, can now be solved to obtain the free vibration response in normal coordinates due to initial conditions  $\underset{\sim}{D}_{N_{0i}}$  and  $\dot{\underset{\sim}{D}}_{N_{0i}}$ . Equations (3.1-51) provide acceleration and displacement response at master degrees of freedom in the original coordinate system.

Proportional loading is assumed for the case of applied forcing functions  $\underset{\sim}{A}_F$  at master degrees of freedom. Therefore,  $\underset{\sim}{A}_F$  can be expressed as the product of a matrix of load factors  $\underset{\sim}{F}_M$ , containing force amplitudes at master degrees of freedom, and normalized force-time histories  $\underset{\sim}{F}(t)$  as follows:

$$\underset{\sim}{A}_F = \underset{\sim}{F}_M \underset{\sim}{F}(t) \quad (3.1-53)$$

where example time-histories for x and y direction forces are shown in Figure (3.1-5). The time-histories are taken to be piecewise-linear functions of time which can be treated as a succession of trapezoidal pulses for computational purposes. Vector  $\underset{\sim}{A}_F$  is transformed to normal coordinates as shown in equation (3.1-49c) to produce the following expression for normal coordinate forces:

$$\underset{\sim}{F}_N = \underset{\sim}{X}_N^{-1} \underset{\sim}{F}_M \underset{\sim}{F}(t) \quad (3.1-54)$$

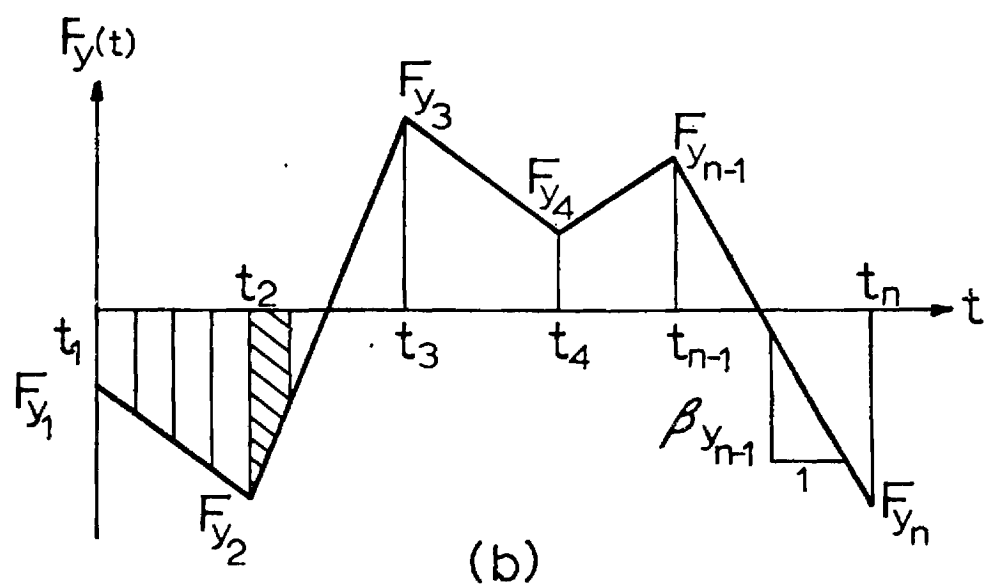
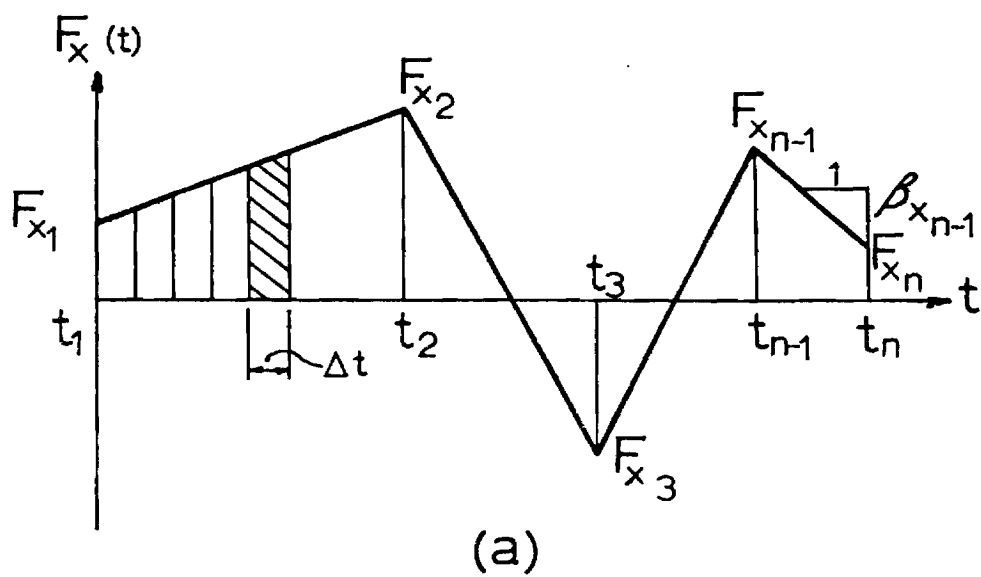


Figure 3.1-5 Piecewise-Linear Forcing Functions



These forces are used in equation (3.1-50) to determine normal mode response and equations (3.1-51) are used to obtain system response in the original coordinate system.

If the response of the structure to the acceleration of a rigid-body support system is sought, a new expression for the normal coordinate forces must be developed for use in the normal mode solution procedure described above. In this case, the vector  $\ddot{\underline{D}}_{\underline{R}}$  of independent restraint accelerations degenerates to  $\ddot{\underline{D}}_{\underline{G}}$  which describes the accelerations of the rigid body support system as piecewise-linear functions of time. Structure response is calculated relative to the supports and the equations of motion must be expressed in relative coordinates. Let  $\underline{r}$  equal displacements at master degrees of freedom due to unit support motions. Then, from equation (3.1-29),

$$\underline{r} = -\underline{S}_{\underline{FF}}^{-1} \underline{S}_{\underline{FR}} \quad (3.1-55)$$

and relative motions of master degrees of freedom can be expressed as

$$\underline{D}_{\underline{F}}^* = \underline{D}_{\underline{F}} - \underline{r} \underline{D}_{\underline{G}} \quad (3.1-56a)$$

$$\ddot{\underline{D}}_{\underline{F}}^* = \ddot{\underline{D}}_{\underline{F}} - \ddot{\underline{r}} \underline{D}_{\underline{G}} \quad (3.1-56b)$$

Writing the first of equations (3.1-38), with  $\underline{A}_{\underline{F}}$  zero, in relative coordinates gives

$$\underline{M}_{\underline{FF}} \ddot{\underline{D}}_{\underline{F}}^* + \underline{S}_{\underline{FF}} \underline{D}_{\underline{F}}^* = -(\underline{M}_{\underline{FF}} \underline{r} + \underline{M}_{\underline{FR}}) \ddot{\underline{D}}_{\underline{G}} \quad (3.1-57)$$

If this equation is transformed to normal coordinates, the appropriate normal coordinate forces are

$$\underline{F}_{\underline{N}} = -\underline{X}_{\underline{N}}' (\underline{M}_{\underline{FF}} \underline{r} + \underline{M}_{\underline{FR}}) \ddot{\underline{D}}_{\underline{G}} \quad (3.1-58)$$

and the normal mode response analysis may be performed using equation (3.1-50).

The final type of excitation to be considered is the case of arbitrary independent motion of multiple supports. In this case, support racking motions are specified as piecewise-linear functions using vectors  $\ddot{D}_R$  and  $\ddot{\tilde{D}}_R$  and total response at master degrees of freedom is obtained using the first of equations (3.1-38). Rearrangement of this equation gives

$$\tilde{M}_{FF}\ddot{\tilde{D}}_F + \tilde{S}_{FF}\tilde{D}_F = -\tilde{M}_{FR}\ddot{\tilde{D}}_R - \tilde{S}_{FR}\tilde{D}_R \quad (3.1-59)$$

Transformation to normal coordinates produces normal coordinate forces

$$\tilde{F}_N = -\tilde{X}'_N\tilde{M}_{FR}\ddot{\tilde{D}}_R - \tilde{X}'_N\tilde{S}_{FR}\tilde{D}_R \quad (3.1-60)$$

and normal mode response analysis proceeds with equation (3.1-50).

Backward substitution operations can be used in all of the cases discussed above to determine remaining frame and panel displacements and stresses at selected times of interest in the dynamic analysis. The solution for nodal displacements, member forces and panel stresses then follows the procedures outlined in the static analysis section above.

## 3.2 Computer Program

### 3.2(a) Description

A FORTRAN computer program, WINDOW, was developed during the course of this research program to evaluate the static and dynamic response of a wide range of glass panels of varying size and support conditions. This program, containing approximately 4200 source statements, implements the finite element model described earlier. All program development and de-

bugging was performed on Georgia Tech's CDC Cyber 74 digital computer. Disk storage is used extensively to save matrix condensation products for later use in the backsubstitution phase of static and dynamic problem solutions.

Program WINDOW is quite versatile and may perform any of a variety of tasks. The program can be used to:

1. Generate stiffness and mass arrays for a finite element model of a panel alone or of a window-framework assembly supported in an arbitrary manner;
2. Develop the framework model with or without elastic connections to account for joint flexibility;
3. Construct a reduced dynamic model of the structure, with arbitrarily-selected master degrees of freedom on the panel, while retaining the ability to compute all panel displacements and stresses through backsubstitution if desired;
4. Compute displacements, stresses, and reactions for applied static loadings or support movements;
5. Compute displacements, stresses, and reactions for applied dynamic nodal loads, time-dependent support displacements and accelerations, or initial conditions of displacement and velocity;
6. Print and/or plot response quantities of interest.

A macro-flow chart and a data input guide for WINDOW are contained in the following sections. Sample analyses using the program are presented in the discussion of results.

### 3.2(b) Macro-Flow Chart

Figure 3.2-1 contains a macro-flow chart showing the basic organiza-

Figure 3.2-1 Macro-Flow Chart of Program WINDOW

1. Read in number (NS) of structures to be analyzed.

Initialize structure number

2. Start loop on structural number

3. Read in structure data for structure number SN

Initialize code variable for stiffness and mass assemblage

4. Begin loop for assemblage of stiffness (ISM = 1) and mass (ISM = 2) matrices

Omit mass matrix assemblage for static problems (ISDA = 0)

Omit mass matrix assemblage if lumped mass matrix is input directly (MCODE = 3)

5. Form panel stiffness (ISM = 1) and mass (ISM = 2) submatrices

Test for existence of boundary frame

6. Form frame stiffness (ISM = 1) and mass (ISM = 2) submatrices

Test for existence of elastic spring members

7. Add spring stiffness terms to structure stiffness submatrices

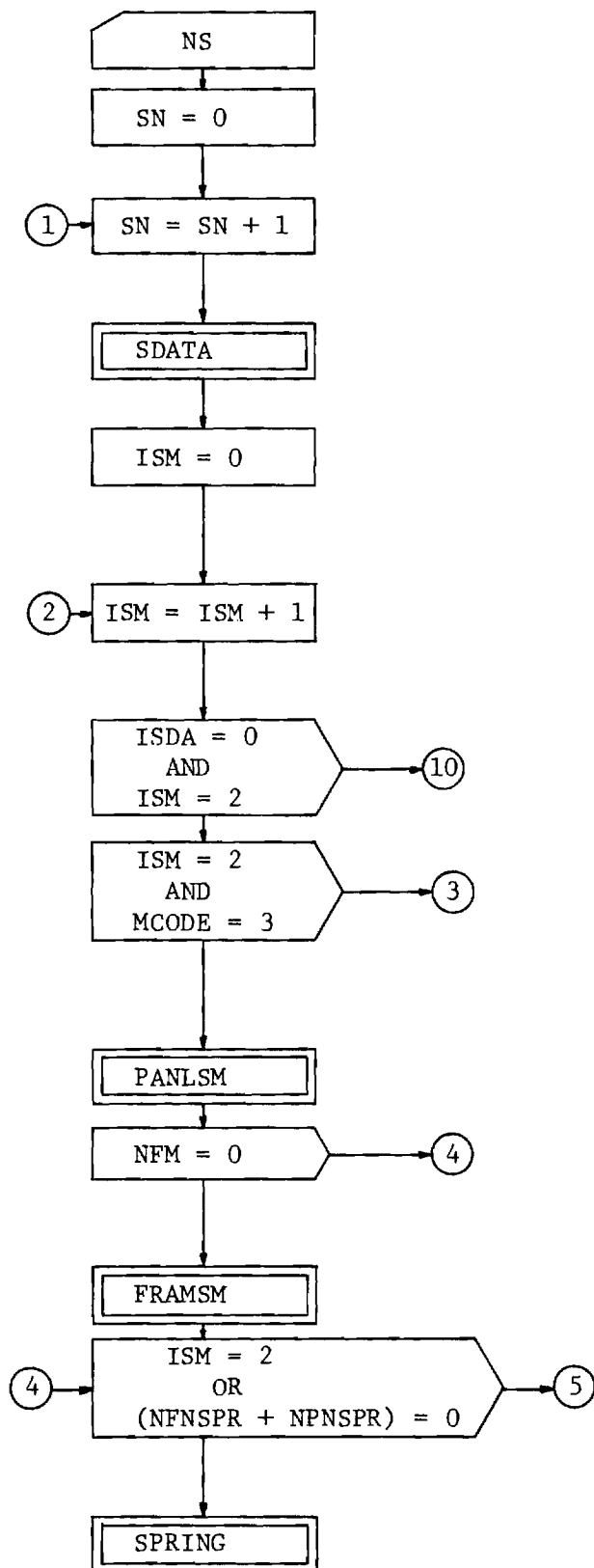


Figure 3.2-1--Cont'd.

8. Condense stiffness (ISM = 1)  
and mass (ISM = 2) matrices

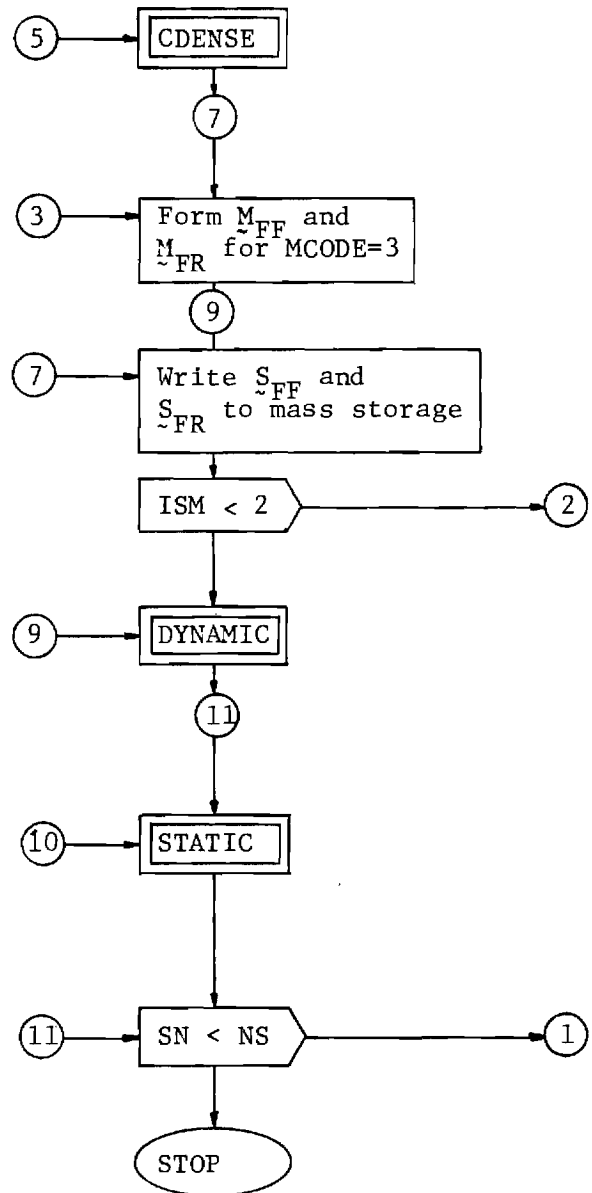
9. Form mass arrays from direct  
input data if MCODE = 3

10. Save stiffness subarrays on  
mass storage prior to mass matrix  
assemblage

11. Determine structure frequen-  
cies, mode shapes, and dynamic  
response

12. Determine all static nodal  
displacements and element stresses

13. Test for last structure



tional layout of the program. Double-lined boxes indicate subprograms and descriptive comments show program operations and flow. A listing of the program is contained in the Appendix.

### 3.2(c) Notation and Input Guide

Key program variables are defined in Table 3.2-1. A data input user's guide is provided in Table 3.2-2.

### 3.3 Discussion of Results

The discrete element analytical model and computer program were used to conduct parametric studies of the static and dynamic performance of insulating windows typical of those employed in the field study building. The initial objective of these studies was to construct a model of a single panel, including mullion and spandrel elements, between floor levels and to adjust model properties until a satisfactory comparison was achieved between measured field and laboratory data and computer program results for the analytical model. This task was not viewed as a system identification problem, however, and formal parameter estimation procedures were not used. Rather, each of the model properties, such as gasket spring constant values and framework area moments of inertia, was varied in turn and the response sensitivity of the model recorded. Model properties were then adjusted in steps until displacements and frequencies were reasonably close to measured laboratory values. System properties which had relatively more influence on panel displacement or frequency response were studied further for varying panel aspect ratios. The results of these more in-depth sensitivity studies are presented in the form of nondimensional curves showing panel response versus selected panel, frame, and gasket properties and panel aspect ratios.

Table 3.2-1 Notation for Program WINDOW

<u>Variable</u>	<u>Definition</u>
A( )	frame member cross-sectional area
AF( )	nodal load vector
AX( ), AY( ), AZ( )	x, y, z accelerations
DENSP, DENSF	panel and frame mass densities
DR( )	support displacements
D1	response calculation time interval
EP, EF	elastic moduli for panel and frame
EZJ( ), EZK( ), EYJ( ), EYK( )	elastic connection parameters
FX( ), FY( ), FZ( )	forcing function time histories
G( )	modal damping ratios
I	member number
IAF	forcing function code
ICNTY( , )	counter-clockwise element node numbers
IDTH, IDMAX	codes to print displacement time history and maximum response
IFM( )	frame member type code
IGA	support acceleration code
IND, INV	initial displacement and velocity codes
IREAC, ISTRES	codes for reaction and stress calculations
IREP( )	code for repeated finite elements
ISD	support displacement code
ISDA	indicator for static (0) or dynamic (1) problem
IX( ), IY( ), IZ( )	frame area moments of inertia
J	joint number
JJ( )	joint at j end of frame member
JK( )	joint at k end of frame member
JPN( )	panel node number
JRL( )	frame joint restraint list
K	joint number
KS( , )	spring stiffness constants
LF( )	load factors at nodes

Table 3.2-1--Cont'd.

<u>Variable</u>	<u>Definition</u>
MCODE	mass assemblage type code
MP( ), MF( )	added lumped masses at panel and frame nodes
NDLC	number of dynamic loading conditions
NFJ	number of frame joints
NFJSD	number of frame joints with support displacements
NFM	number of frame members
NFL( )	panel master degree of freedom list
NFNSPR	number of frame nodes with attached springs
NFT	number of F displacement types
NFR	number of frame restraints
NJPLOT	number frame displacement-time history plots
NLN	number of loaded nodes
NMEC	number of frame members with elastic connections
NMODES	number of modes in response analysis
NNIND, NNINV	number of nodes with initial displacements, velocities specified
NNLF	number of nodes with load factors
NNPLOT	number of element displacement-time history plots
NNWFT	number of nodes with F displacement types
NPE	number of panel elements
NPN	number of panel nodes
NPNSD	number of panel nodes with specified support displacements
NPNSPR	number of panel nodes with springs
NPR	number of panel restraints
NRFJ	number of restrained frame joints
NRL( )	nodal restraint list
NRPN	number of restrained panel nodes
NS	number of structures to be analyzed
NSATI	number of static stress analysis times of interest
NSLC	number of static loading conditions
NTX, NTY, NTZ	number of time points in forcing function vectors
PRP, PRF	Poisson's ratio for panel and frame
PLIST( )	degree of freedom plot list



Table 3.2-1--Cont'd.

<u>Variable</u>	<u>Definition</u>
SKPO, SKL, SKN, SKR	elastic spring constants for gaskets
TITLE	descriptive problem title
TP	panel thickness
TX( ), TY( ), TZ( )	times corresponding to data points in FX, FY, FZ
UNITS	active units in problem for calculations
VOP	initial velocity vector
XP( ), YP( )	panel node coordinates
XF( ), YF( )	frame joint coordinates
XOP( )	initial displacement vector

Table 3.2-2 Input Data for Program WINDOW

<u>Data</u>	<u>Restrictions</u>
1. NS	provide NS sets of data (all data free format)
2. TITLE	
3. UNITS	
4. NPE, NFM, ISDA	
5. NPN, NRPN, NPR, NNWFT, NFT, EP, PRP, DENSP, TP, MCODE	
6. N, XP(N), YP(N), MP(N)	NPN cards
7. K, (ICNTY(K,J)J=1,4), IREP(K)	NPE cards
8. J, NFL(8J-7),...,NFL(8J)	NNWFT cards
9. J, NRL(8J-7),...,NRL(8J)	If NRPN>0, NRPN cards
10. NFJ, NFR, NRFJ, NMEC, EF, PRF, DENSP	If NFM>0
11. J, XF(J), YF(J), MF(J)	If NFM>0, NFJ cards
12. I, JJ(I), JK(I), A(I), IX(I), IY(I), IZ(I), IFM(I)	If NFM>0, NFM cards
13. J, JRL(6J-5),..., JRL(6J)	If NFM>0, NRFJ cards
14. I, EZJ(I), EZK(I), EYJ(I), EYK(I)	If NFM and NMEC>0, NMEC cards
15. NFNSPR, NPNSPR	
16. SKPO, SKL, SKN, SKR	If NFNSPR>0
17. JPN(I), (KS(I,J),J=1,5)	If NPNSPR>0

Static Load Data: Items 18 - 22

18. NSLC	If ISDA = 0
19. NLN, NPNSD, NFJSD, IREAC, ISTRES	If ISDA = 0
20. K, AF(8K-7),...,AF(8K)	If ISDA = 0 and NLN>0, NLN cards
21. K, DR(6K-5),...,DR(6K)	If ISDA = 0 and NFJSD>0, NFJSD cards
22. K, DR(8K-7),...,DR(8K)	If ISDA = 0 and NPNSD>0, NPNSD cards

Dynamic Load Data: Items 23 - 48

23. NDLC	If ISDA = 1
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Table 3.2-2--Cont'd.

<u>Data</u>	<u>Restrictions</u>
24. IND, INV, IGA, IAF, ISD	If ISDA = 1 and NDLC > 0
25. NMODES, NNPLOT, NJPLOT, IDTH, IDMAX	If ISDA = 1 and NDLC > 0
26. N, PLIST(8N-7),...,PLIST(8N)	If ISDA = 1, NDLC > 0, and NNPLOT > 0, NNPLOT cards
27. N, PLIST(6N-5),...,PLIST(6N)	If ISDA = 1, NDLC > 0, NJPLOT > 0, and NFM > 0, NJPLOT cards
28. G( )	If ISDA = 1, and NDLC > 0
29. NNIND	If ISDA = 1, NDLC > 0, and IND = 1
30. N, XOP(8N-7),...,XOP(8N)	If ISDA = 1, NDLC > 0, and IND = 1, NNIND cards
31. NNINV	If ISDA = 1, NDLC > 0, and INV = 1
32. N, VOP(8N-7),...,VOP(8N)	If ISDA = 1, NDLC > 0, and INV = 1, NNINV cards
33. NNLF	If ISDA = 1, NDLC > 0, and IAF = 1
34. N, LF(8N-7),...,LF(8N)	If ISDA = 1, NDLC > 0, and IAF = 1, NNLF cards
35. K, LF(8K-8),...,LF(8K)	*If ISDA = 1, NDLC > 0, ISD = 1, and NPR > 0, NRPN cards
36. K, LF(6K-5),...,LF(6K)	*If ISDA = 1, NDLC > 0, ISD = 1, and NFR > 0, NRFJ cards
37. NTX, NTY, NTZ, D1	If ISDA = 1, NDLC > 0
38. FX( )	If ISDA = 1, NDLC > 0, NTX points
39. AX( )	If ISDA = 1, NDLC > 0 and ISD = 1, NTX points
40. TX( )	If ISDA = 1, NDLC > 0, NTX points
41. FY( )	If ISDA = 1, NDLC > 0, NTY points
42. AY( )	If ISDA = 1, NDLC > 0 and ISD = 1, NTY points
43. TY( )	If ISDA = 1, NDLC > 0, NTY points
44. FZ( )	If ISDA = 1, NDLC > 0, NTZ points
45. AZ( )	If ISDA = 1, NDLC > 0 and ISD = 1, NTZ points
46. TZ( )	If ISDA = 1, NDLC > 0, NTZ points
47. NSATI, IREAC, ISTRES	If ISDA = 1, NDLC > 0
48. PLIST( )	If ISDA = 1, NDLC > 0, NSATI points

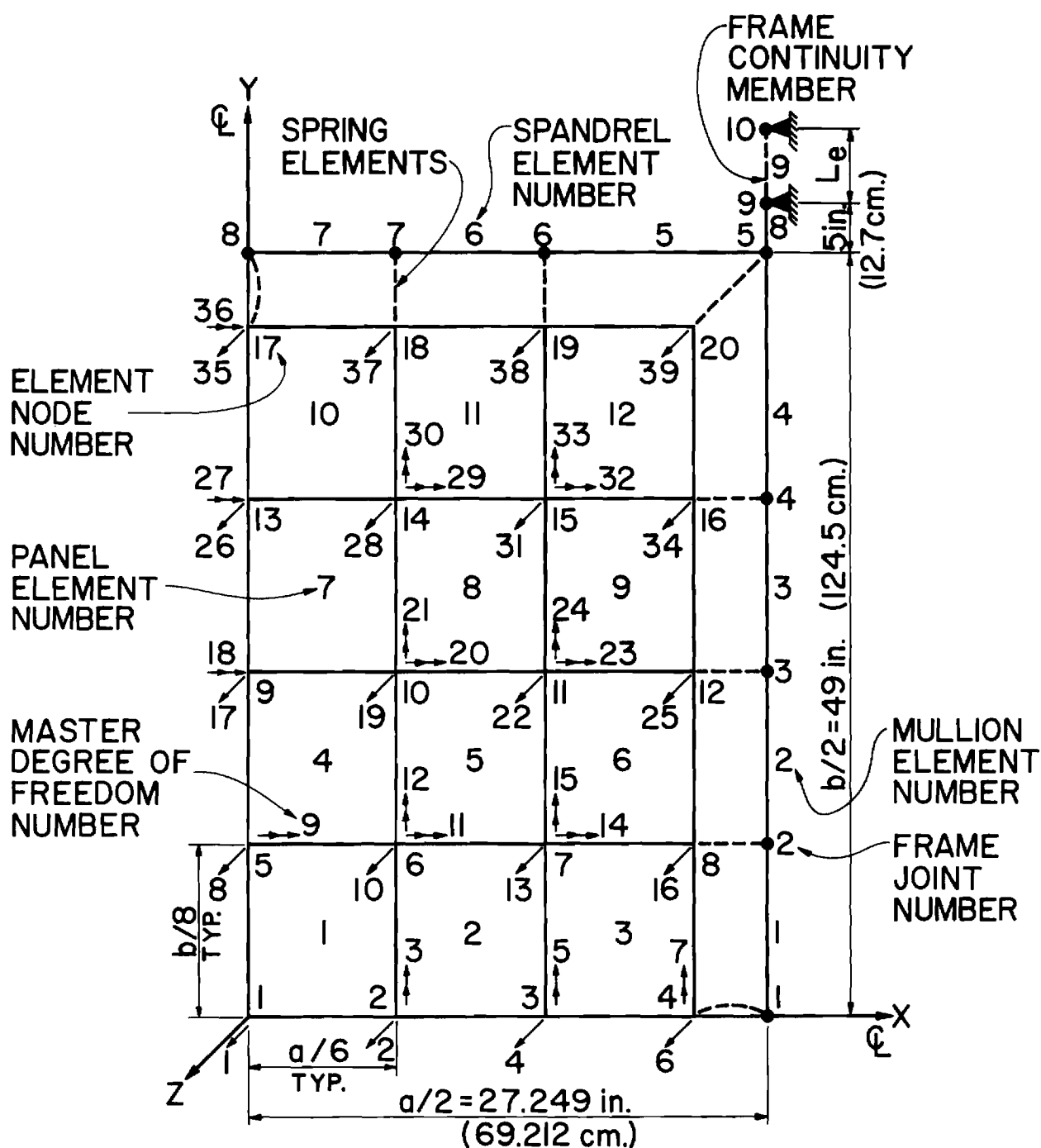
\* translational x, y, and z support motion only

Sample analyses in which model properties were adjusted to calibrate the analytical model are described first. Then, nondimensional ratios are defined and sensitivity analysis results presented for the static and dynamic response studies. Several displacement-time history plots showing the influence of key system parameters on dynamic response of the model are presented next. Finally, panel response to in-plane racking distortions caused by differential motions of adjacent floors is presented.

### 3.3(a) Sample Analyses

Model Description: Several finite element models of the panel-frame structure with different finite element meshes, properties, and boundary conditions were assembled, but that shown in Figure 3.3-1 produced the best correlation with experimental results and was used for most sensitivity studies. In this model, the x-z and y-z planes are planes of symmetry; therefore, only symmetric loadings, displacement patterns, and modes of vibration are considered. Appropriate symmetry boundary conditions are imposed along the x and y axes. Studies of the in-plane racking behavior of the system, however, employed a model in which symmetry was not considered. This model of the entire panel-frame structure between floor levels is described in Section 3.3(e).

The symmetry model features twelve finite elements, nine frame elements, and eight points of attachment of gasket spring elements at frame joints 1 to 8. Corresponding frame and panel nodes to which gasket springs are attached are assumed to have the same joint coordinates even though they are shown as distinct points in Figure 3.3-1. The frame is simply-supported at frame joint 9 but the length ( $L_e$ ) of member 9 can be



ASPECT RATIO  $R = b/a$

Figure 3.3-1 Finite Element Model of Quarter Panel

selected arbitrarily to provide values of rotational stiffness at joint 9 which lie between the extremes of zero (simple support) and infinity (fixed support). Length  $L_e$  is adjusted to simulate the effect of mullion continuity on panel response. If  $L_e$  is zero, member 9 is omitted from the frame and continuity is neglected. In addition, an elastic connection for rotation about the global y axis was inserted in frame member 5 at joint 5 to investigate the effect of a nonrigid spandrel-mullion connection on panel response.

Thirty-nine master degrees of freedom were selected from the 160 panel displacements available in the model. Enough master degrees of freedom must be chosen to ensure that distributed loadings, static displacement patterns, and vibration mode shapes are adequately represented.

The panel aspect ratio  $b/a$  was varied between 0.5 and 5.0 in the sensitivity studies with 1.8 being the actual ratio for insulating windows in the prototype structure. Discrete values of 0.5, 1.0, 1.8, 2.5, and 5.0 were selected. The area of the quarter panel model was held constant at  $1335.2 \text{ in}^2$  ( $8614.2 \text{ cm}^2$ ) so that total load for distributed static pressure loadings and total panel mass would remain constant for static and dynamic sensitivity studies.

Model Properties: Elastic constants, mass densities, cross-sectional areas and area moments of inertia of frame members, and gasket spring constants are summarized in Table 3.3-1. Area moments of inertia are taken with respect to x-y-z axes in the member cross-section which are parallel to the coordinate axes in Figure 3.3-1. Gasket spring constant values are expressed as spring constant per unit length of boundary and are arranged as shown in Figure 3.1-1. These values were determined from simple laboratory experiments described in Chapter 4.

Table 3.3-1 Model Properties

Glass Panel:

Elastic Modulus (E)	$1 \times 10^7$ psi ( $6.9 \times 10^7$ kN/m <sup>2</sup> )
Poisson's Ratio ( $\nu$ )	0.23
Mass Density ( $\rho$ )	5 slugs/ft <sup>3</sup> (2575 kg/m <sup>3</sup> )

Frame:

Elastic Modulus (E)	$1 \times 10^7$ psi ( $6.9 \times 10^7$ kN/m <sup>2</sup> )
Poisson's Ratio ( $\nu$ )	0.25
Mass Density ( $\rho$ )	5.37 slugs/ft <sup>3</sup> (2766 kg/m <sup>3</sup> )
Area Moments of Inertia	
Mullion Elements	
Ix	8.321 in <sup>4</sup> (346.3 cm <sup>4</sup> )
Iy	11.390 in <sup>4</sup> (474.1 cm <sup>4</sup> )
Iz	3.069 in <sup>4</sup> (127.7 cm <sup>4</sup> )
Spandrel Elements	
Ix	3.452 in <sup>4</sup> (143.7 cm <sup>4</sup> )
Iy	2.481 in <sup>4</sup> (103.3 cm <sup>4</sup> )
Iz	0.971 in <sup>4</sup> (40.4 cm <sup>4</sup> )
Cross-sectional Areas	
Mullion (A)	2.613 in <sup>2</sup> (1686 mm <sup>2</sup> )
Spandrel (A)	1.404 in <sup>2</sup> (906 mm <sup>2</sup> )

Gasket Equivalent Spring Constants:

$k_L$	318.0 lb/in/in (2194 kN/m/m)
$k_P$	81.1 lb/in/in (560 kN/m/m)
$k_N$	556.9 lb/in/in (3843 kN/m/m)
$k_R$	305.3 in-lb/radian/in (1358 N·m/radian/m)

The one inch (2.54 cm) thick insulating windows in the prototype structure consist of two 0.25 in. (0.64 cm) thick glass panels separated by a sealed 0.50 in. (1.27 cm) thick air space with metal spacer bars on the periphery of the panel. For simplicity, the double plate system was replaced by an equivalent single panel which could be represented by a single-layer finite element model. This model assumes that both plates act together and vibrate in phase. (A double plate continuum model and a parallel-plate finite element model were also developed to study out-of-phase motions of the double panel system and are discussed in Section 3.4.).

The thickness of the equivalent single panel was determined by assuming that the flexural rigidity of the single plate is equal to the flexural rigidity of the double plate system. Using the plate properties in Table 3.3-1, a thickness of 0.315 in. (0.8 cm) is computed for the equivalent single panel. However, measured glass thicknesses were somewhat less for the window tested in the laboratory pressure test assembly, having a measured total thickness of 0.9803 in. (2.49 cm) and a glass panel thickness of 0.237 in. (0.60 cm). If these thicknesses are used, a value of 0.298 in. (0.76 cm) results for the equivalent single panel thickness. This lower value was used in the panel sensitivity studies rather than the value of 0.315 in. (0.80 cm) based on nominal panel dimensions.

Results: Initial studies employed a finite element model of the panel only and were used to check program WINDOW output against classical plate theory results. Static displacement and vibration frequency comparisons are not presented here for brevity but finite element results were within one or two percent of continuum theory values.



Other calibration studies using the frame-panel model shown in Figure 3.3-1 investigated the effect of changes in panel thickness and density and in frame continuity member length ( $L_e$ ) on panel frequencies. Results of these studies are presented in Tables 3.3-2, 3.3-3, and 3.3-4. A consistent mass formulation was used for both frame and panel elements and the cross-sectional area of the continuity member was set to a very small value so that its mass would not affect panel vibrations.

Results presented in Table 3.3-2 show that increased panel thickness results in increased structure fundamental frequency for all cases except that of panel aspect ratio equal to 5. In this case increased panel thickness adds more mass than stiffness effect resulting in a reduction in panel frequency with thickness for the range of panel thicknesses and aspect ratios considered. Table 3.3-3 shows that panel frequencies decrease somewhat uniformly with increasing mass density for all panel aspect ratio cases as expected. Table 3.3-4 displays the variation in structure fundamental frequency with the length  $L_e$  of the continuity member. This member provides variable rotational stiffness at joint 9 equal to  $3EI/L_e$  where  $I$  is the area moment of inertia of the continuity member about the x axis in Figure 3.3-1. Results in Table 3.3-4 show that added rotational stiffness at joint 9 changes structure frequency by less than 4% for the range of values of length  $L_e$  considered. Similar results were obtained when the effect of a nonrigid mullion-spandrel connection at joint 5 (Figure 3.3-1) was investigated. In this case the effect of the elastic connection on static displacement and vibration frequency results was negligible.

Figures 3.3-2, 3.3-3, and 3.3-4 show the frequencies and mode shapes of the first three symmetric structure modes for model properties in Table

Table 3.3-2 Fundamental Frequency of Framed Panels  
of Varying Thickness and Aspect Ratio

Equivalent Single Panel Thickness, in inches (mm) (1)	Fundamental Frequency, in Hertz, for Varying Panel Aspect Ratios, R				
	R = 0.5 (2)	R = 1.0 (3)	R = 1.8 (4)	R = 2.5 (5)	R = 5.0 (6)
0.150 (3.81)	8.82	7.25	8.64	10.41	12.04
0.200 (5.08)	9.82	8.46	9.90	11.58	11.58
0.250 (6.35)	10.63	9.64	11.12	12.57	11.12
0.298 (7.57)	11.23 (13.23) <sup>a</sup>	10.77 (10.59) <sup>a</sup>	12.24 (12.46) <sup>a</sup>	13.31 (15.35) <sup>a</sup>	10.70 (27.53) <sup>a</sup>
0.350 (8.89)	11.64	11.94	13.31	13.81	10.28
0.400 (10.2)	11.83	12.97	14.14	14.04	9.90
0.450 (11.4)	11.90	13.86	14.75	14.08	9.56
<p>NOTE: See Figure 3.3-1 and Table 3.3-1 for panel dimensions and properties; panel density = 5 slugs/ft<sup>3</sup> (2575 kg/m<sup>3</sup>); no continuity member in frame.</p> <p><sup>a</sup>Classical plate theory frequency for simply-supported plate with same dimensions and properties.</p>					

Table 3.3-3 Fundamental Frequency of Framed Panels  
of Varying Density and Aspect Ratio

Mass Density of Panel, in slugs/ft <sup>3</sup> (kg/m <sup>3</sup> ) (1)	Fundamental Frequency, in Hertz, for Varying Panel Aspect Ratios, R				
	R = 0.5 (2)	R = 1.0 (3)	R = 1.8 (4)	R = 2.5 (5)	R = 5.0 (6)
4.0 (2060)	12.50	12.04	13.66	14.77	11.52
4.5 (2318)	11.81	11.35	12.89	13.98	11.09
5.0 (2575)	11.23	10.77	12.24	13.31	10.70
5.5 (2833)	10.72	10.27	11.68	12.72	10.35
6.0 (3090)	10.28	9.83	11.18	12.20	10.03
6.5 (3348)	9.89	9.45	10.75	11.74	9.74
7.0 (3605)	9.54	9.11	10.36	11.33	9.47
NOTE: See Figure 3.3-1 and Table 3.3-1 for panel dimensions and properties; panel thickness = 0.298 in. (7.57 mm); no continuity member in frame					

Table 3.3-4 Fundamental Frequency of Framed Panel  
with Continuity Member of Varying Length

Continuity Member Length, $L_e$ , in inches (cm) (1)	Fundamental Frequency, Hertz (2)
0.0 <sup>a</sup>	12.24
1000 (2540)	12.35
500 (1270)	12.44
324 (823)	12.52
270 (686)	12.55
252 (640)	12.57
234 (594)	12.58
<sup>a</sup> No continuity member in frame	
NOTE: See Figure 3.3-1 and Table 3.3-1 for dimensions and properties of panel and frame; panel thickness = 0.298 in (7.57 mm); panel aspect ratio $R = 1.8$ .	

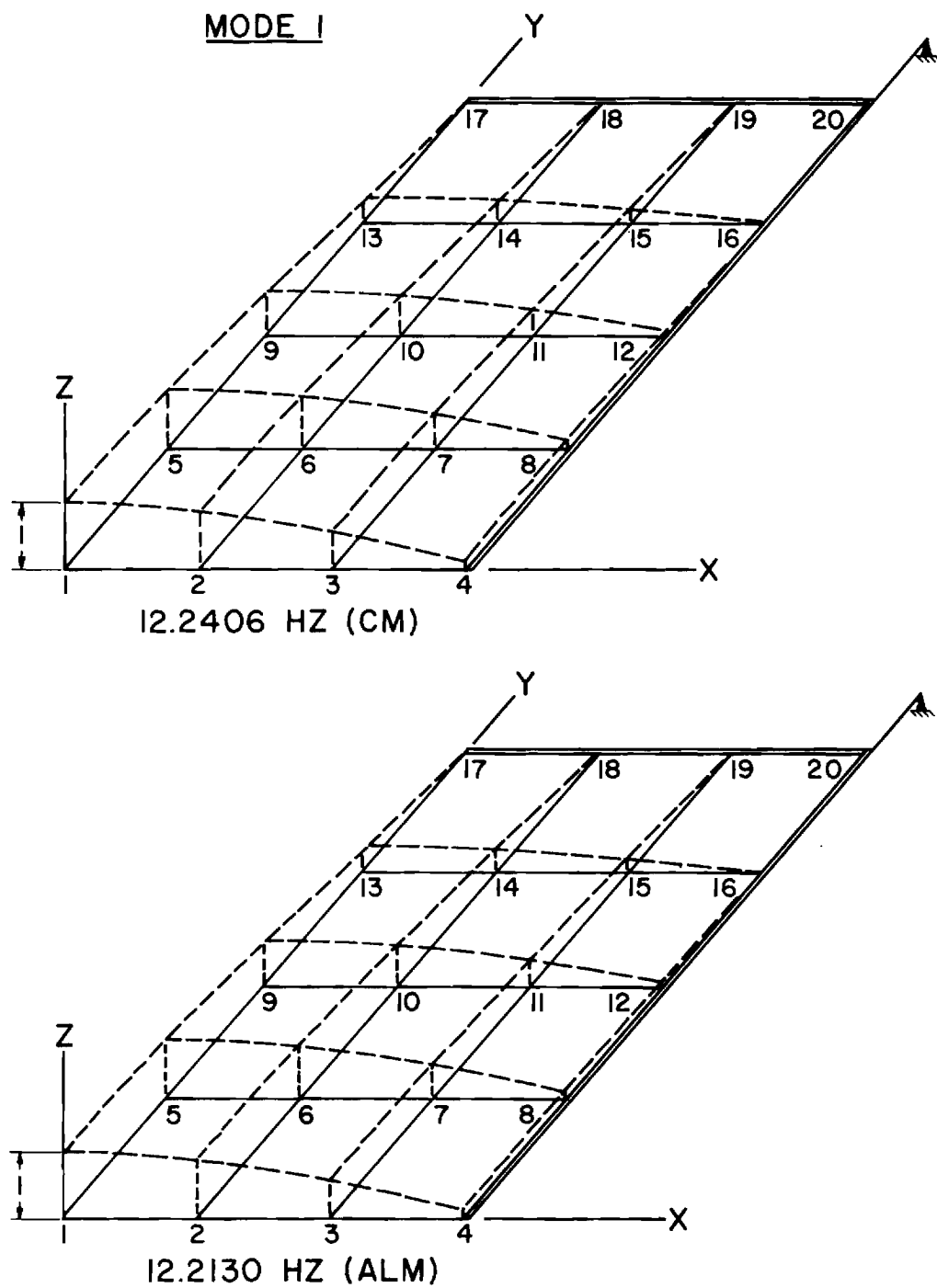


Figure 3.3-2 Plate Bending Frequencies and Mode Shapes, Mode 1

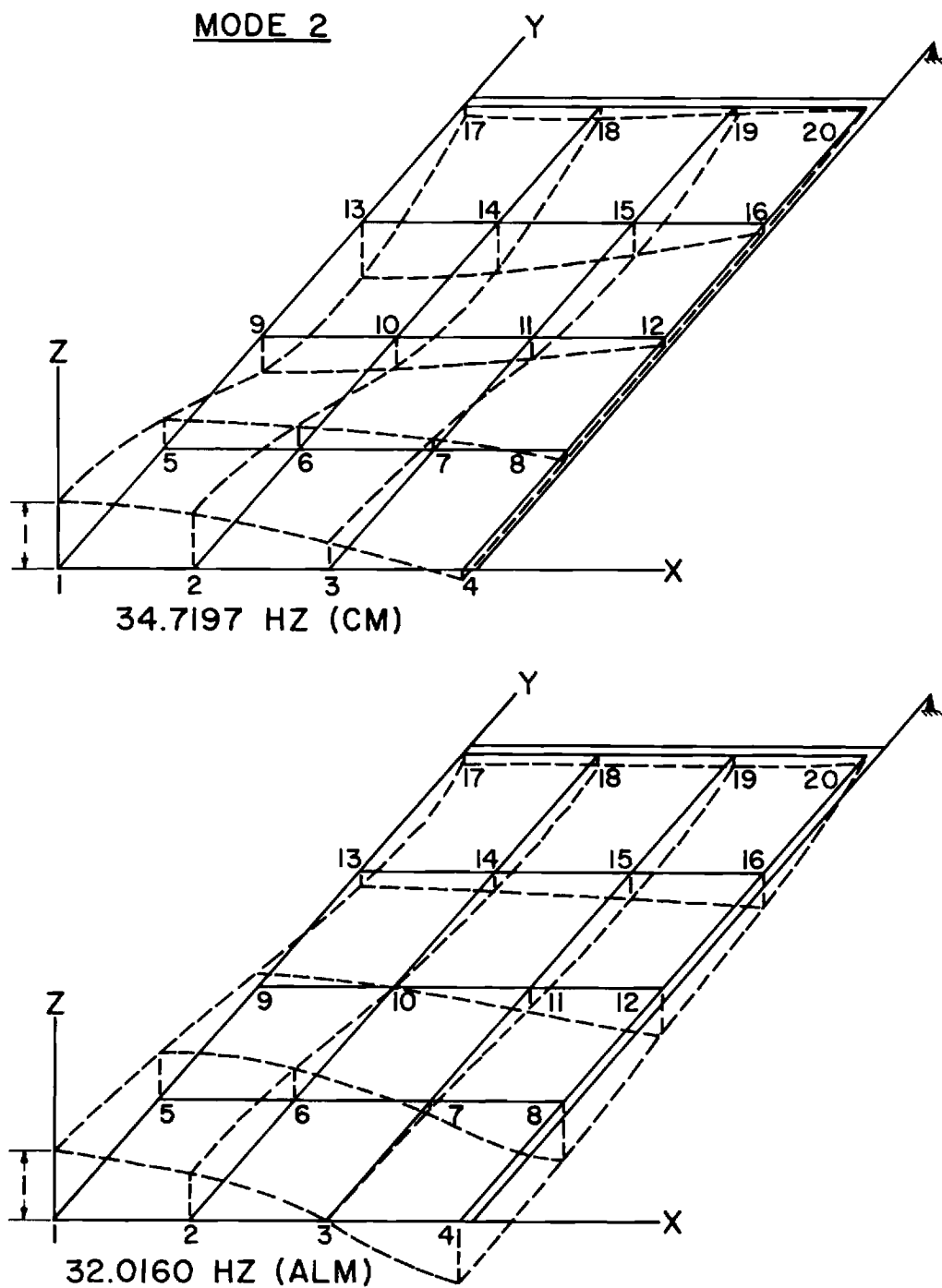


Figure 3.3-3 Plate Bending Frequencies and Mode Shapes, Mode 2

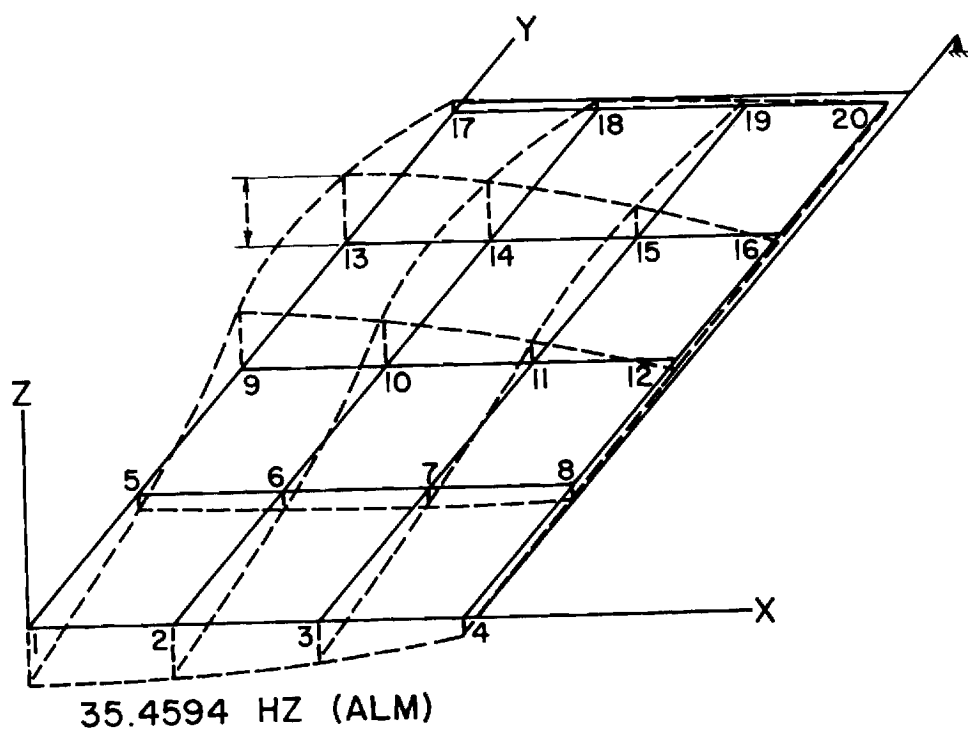
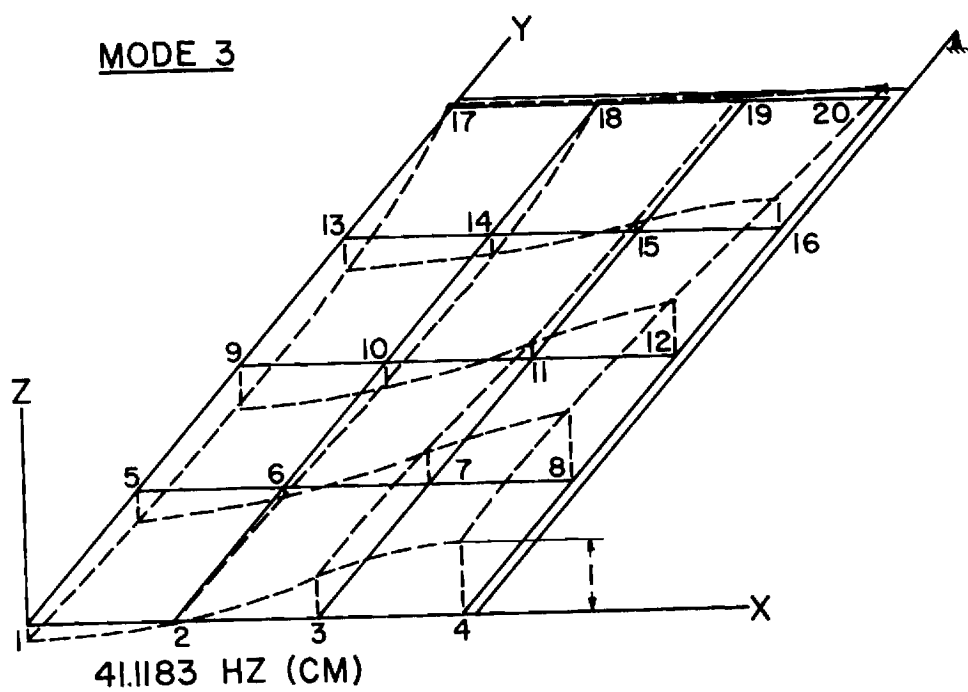


Figure 3.3-4 Plate Bending Frequencies and Mode Shapes, Mode 3

3.3-1 and panel thickness of 0.298 in. (7.57 mm) and  $L_e = 0$ . The consistent mass (CM) and assembled lumped mass (ALM) formulations are compared and show excellent agreement in frequency and mode shape for mode 1. The CM formulation has mass terms associated with all thirty-nine master degrees of freedom in Figure 3.3-1 while the ALM model assigns mass to the twenty nodal translations in the Z direction only. Mode 2 and 3 frequency values for CM and ALM differ by 8% and 14%, respectively. In addition, mode 2 for ALM contains significantly more boundary displacement than the CM shape, and mode 3 shapes are quite dissimilar. Because the CM model contains more master degrees of freedom with mass terms associated with both translational and rotational displacements, the CM mode shapes were assumed to be more accurate and the CM model was used in all sensitivity and dynamic response studies.

The effect of change in stiffness constant  $k_N$  for the boundary spring normal to the panel on mode shapes for mode 1 was also studied for panels of varying aspect ratio but of constant surface area. Model properties in Table 3.3-1 were used, with panel thickness equal to 0.298 in. (0.76 cm) and  $L_e$  equal to zero, but  $k_N$  took on values between 10 lb/in/in (69 kN/m/m) (soft spring) and  $1 \times 10^6$  lb/in/in ( $6.9 \times 10^6$  kN/m/m) (hard spring) with rotational spring constant  $k_R$  held fixed at 305.3 in-lb/radian/in (1358 N·m/radian/m). Mode 1 displacements along x and y symmetry axes are plotted in Figures 3.3-5 and 3.3-6 and corresponding frequencies are recorded in Table 3.3-5. The percent increase in frequency with change in  $k_N$  is also shown in Table 3.3-5. Figure 3.3-6 shows that the value of  $k_N$  has negligible effect on the first mode shape of the panel with aspect ratio of 5.0.



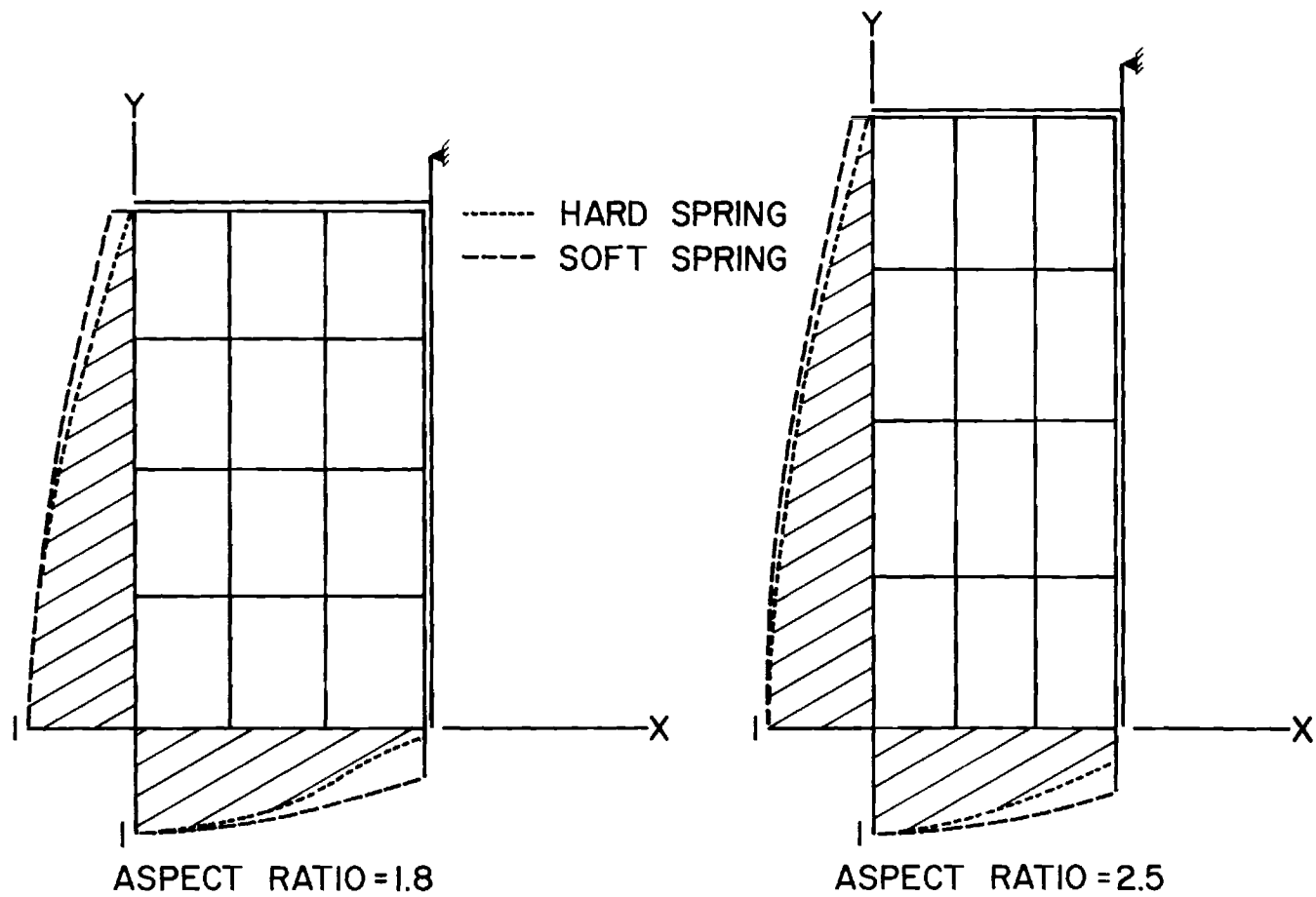


Figure 3.3-5 Effect of Normal Spring on Mode Shape 1,  $R = 1.8$  and  $2.5$

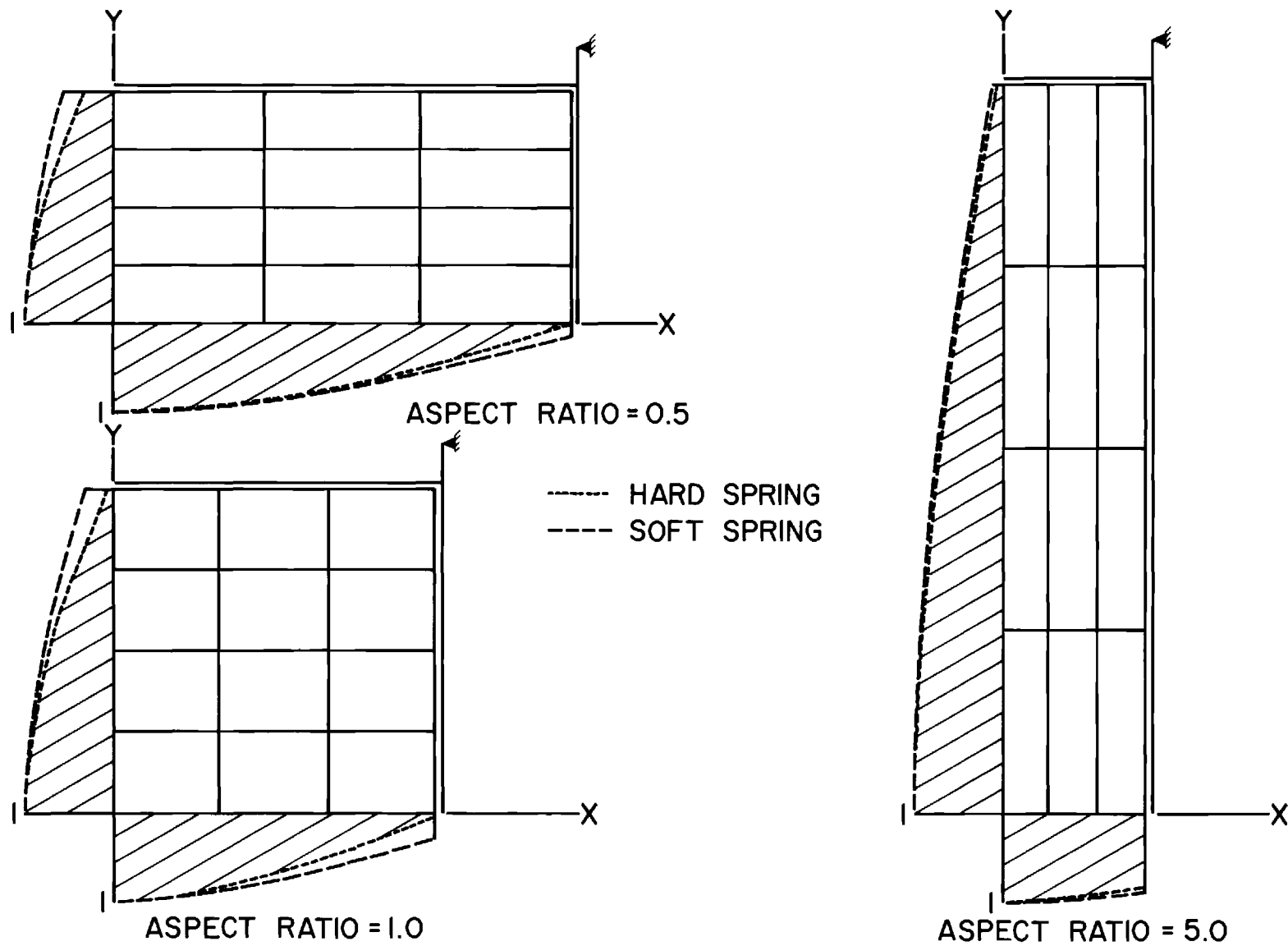


Figure 3.3-6 Effect of Normal Spring on Mode Shape 1,  $R = 0.5, 1, \text{ and } 5$

Table 3.3-5 Fundamental Frequencies of Framed Panels  
of Varying Normal Spring Constant and Aspect Ratio

Aspect Ratio, R (1)	Fundamental Frequency, in Hertz		Frequency Increase (%) (4)
	Soft Spring <sup>a</sup> (2)	Hard Spring <sup>b</sup> (3)	
0.5	8.74	11.31	29
1.0	8.51	10.87	28
1.8	9.19	12.24 <sup>c</sup>	33
2.5	9.73	13.43	38
5.0	9.20	10.73	17

<sup>a</sup>Normal spring constant  $k_N = 10 \text{ lb/in/in}$  (69 kN/m/m)

<sup>b</sup>Normal spring constant  $k_N = 1 \times 10^6 \text{ lb/in/in}$  ( $6.9 \times 10^6 \text{ kN/m/m}$ )

<sup>c</sup>Classical plate theory, simply-supported plate: 12.46 Hz

NOTE: See Table 3.3-1 for other model properties.

If the value of the boundary rotational spring constant  $k_R$  is varied between 50 in-lb/radian/in (222 N·m/radian/m)(soft spring) and  $1 \times 10^6$  in-lb/radian/in ( $4.45 \times 10^6$  N·m/radian/m), the range of frequencies shown in Table 3.3-6 is obtained. Comparison of percent frequency increase with increase in  $k_N$  or  $k_R$  in Tables 3.3-5 and 3.3-6 shows that a change in  $k_R$  affects the fundamental frequency of structures with aspect ratios of 1.0 and 1.8 most, while aspect ratio 0.5, 2.5, and 5.0 panels are more heavily influenced by a change in  $k_N$  for the range of values considered.

Further model sensitivity results are presented below in the form of nondimensional plots of static and dynamic structure response versus variation in selected system parameters.

### 3.3(b) Nondimensional Ratios

Static and dynamic sensitivity studies employ the nondimensional ratios in Table 3.3-7 in plots showing variation of system parameters versus system response. Table 3.3-1 should be consulted for specific values of terms in the ratios except that panel thickness  $t_p$  equals 0.298 in. (0.76 cm) and  $L_e = 0$  for all figures. The symmetry model in Figure 3.3-1 was used throughout.

Parameters  $K_1$  through  $K_4$  are nondimensional forms of normal and rotational spring constants  $k_N$  and  $k_R$ . If system response is plotted for varying  $k_N$  or  $k_R$  and varying aspect ratio  $R$ , terms  $K_3$  and  $K_4$  containing constant panel area  $A_p$  are used. If results are presented for one aspect ratio only, variables  $K_1$  and  $K_2$  are used.

Parameters  $I_1$  through  $I_8$  are dimensionless ratios for variation of mullion and spandrel area moments of inertia.  $I_1$  and  $I_2$  are flexural

Table 3.3-6 Fundamental Frequencies of Framed Panels  
of Varying Rotational Spring Constant and Aspect Ratio

Aspect Ratio, R (1)	Fundamental Frequency, in Hertz		Frequency Increase (%) (4)
	Soft Spring <sup>a</sup> (2)	Hard Spring <sup>b</sup> (3)	
0.5	10.83	13.47	24
1.0	10.19	15.28	50
1.8	11.61	17.07 <sup>c</sup>	47
2.5	12.80	16.43	28
5.0	10.66	10.94	3
<sup>a</sup> Rotational spring constant $k_R = 50 \text{ in-lb/radian/in}$ ( $222 \text{ N}\cdot\text{m/radian/m}$ ) <sup>b</sup> Rotational spring constant $k_R = 1 \times 10^6 \text{ in-lb/radian/in}$ ( $4.45 \times 10^6 \text{ N}\cdot\text{m/radian/m}$ ) <sup>c</sup> Classical plate theory, clamped plate: 24.42 Hz			

Table 3.3-7 Nondimensional Parameters

<u>Parameter</u>	<u>Definition</u>
R	Aspect Ratio, $b/a$
$K_1$	$k_n \bar{b}^3/D$ , $\bar{b} = b/2$
$K_2$	$k_R \bar{b}/D$
$K_3$	$k_N A_p t_p/D$ , $A_p = ab/4 = \bar{a} \bar{b}$ , $\bar{a} = a/2$
$K_4$	$k_R A_p / t_p D$
$I_1$	$EI_x / \bar{b} D$ , mullion
$I_2$	$EI_x t_p / A_p D$ , mullion
$I_3$	$EI_y / \bar{b} D$ , mullion
$I_4$	$EI_y t_p / A_p D$ , mullion
$I_5$	$EI_y / \bar{a} D$ , spandrel
$I_6$	$EI_y t_p / A_p D$ , spandrel
$I_7$	$EI_x / \bar{a} D$ , spandrel
$I_8$	$EI_x t_p / A_p D$ , spandrel
$W_1$	$w_{\max} / w_{\max,ss}^*$ , static displacement
f	$f_{11} / f_{11,ss}^*$ , fundamental frequency
$W_2$	$w_{\max} / w_{\max,ss}^*$ , dynamic displacement

\* Simply-supported plate with  $R = 1.8$

inertias for mullion elements and  $I_5$  and  $I_6$  for spandrel elements. Terms  $I_3$  and  $I_4$  are nondimensional torsion constants for mullion elements and  $I_7$  and  $I_8$  for spandrel elements. Quantities  $I_2$ ,  $I_4$ ,  $I_6$ , and  $I_8$  are used in plots containing system response for more than one aspect ratio; response plots using  $I_1$ ,  $I_3$ ,  $I_5$ , and  $I_7$  are for aspect ratio  $R = 1.8$  only.

Symbols  $W_1$ ,  $f$ , and  $W_2$  refer to the nondimensional static displacement, frequency, and dynamic displacement response of the system in Figure 3.3-1 for fixed or varying panel aspect ratio,  $R$ .  $W_1$  is the ratio of static displacement at node 1 (center of panel) for panels of fixed or varying  $R$  to the maximum static displacement of a simply-supported plate with  $R = 1.8$  and identical properties and loading. Parameter  $f$  is the fundamental structure frequency divided by the fundamental frequency of a simply-supported plate with  $R = 1.8$  from classical plate theory. Finally,  $W_2$  is the ratio of the maximum dynamic response of the structure to a specified pressure loading to the maximum dynamic response of a simply-supported plate ( $R = 1.8$ ) subjected to the same loading.

### 3.3(c) Static Analysis Results

Static displacement response to uniform pressure loading over the entire quarter panel versus variation in spring and frame properties for selected aspect ratios are presented in Figures 3.3-7 through 3.3-10. Data points on all plots were generated using program WINDOW and then connected by smooth curves. Numerical values of parameters defined in Table 3.3-7 which were held fixed are recorded on each figure.

Figure 3.3-7 displays panel center displacement  $W_1$  versus normal spring constant  $K_3$ .  $W_1$  is seen to be relatively insensitive to normal

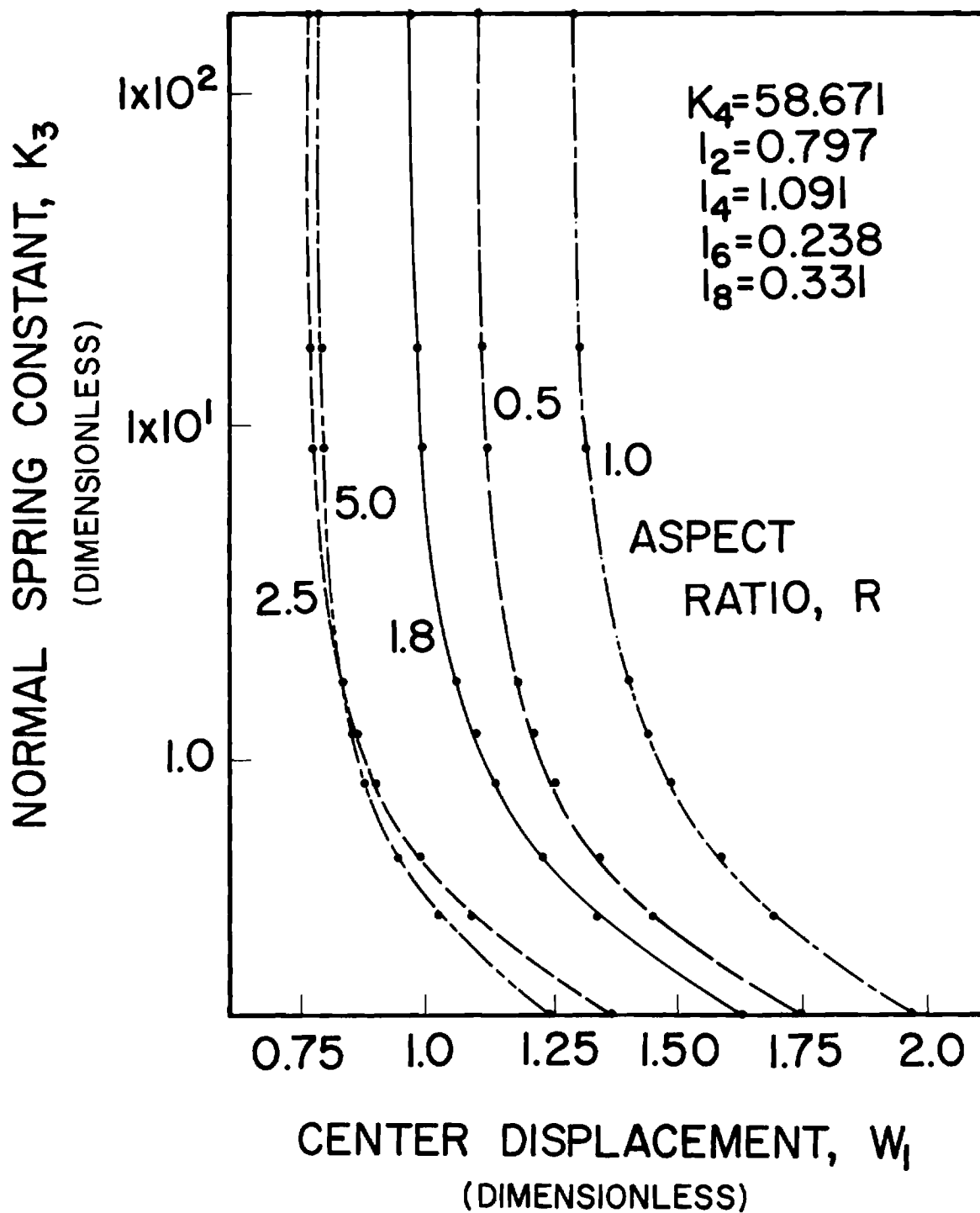


Figure 3.3-7 Static Displacement of Panels with Uniform Loading Versus Normal Spring Constant on Boundary



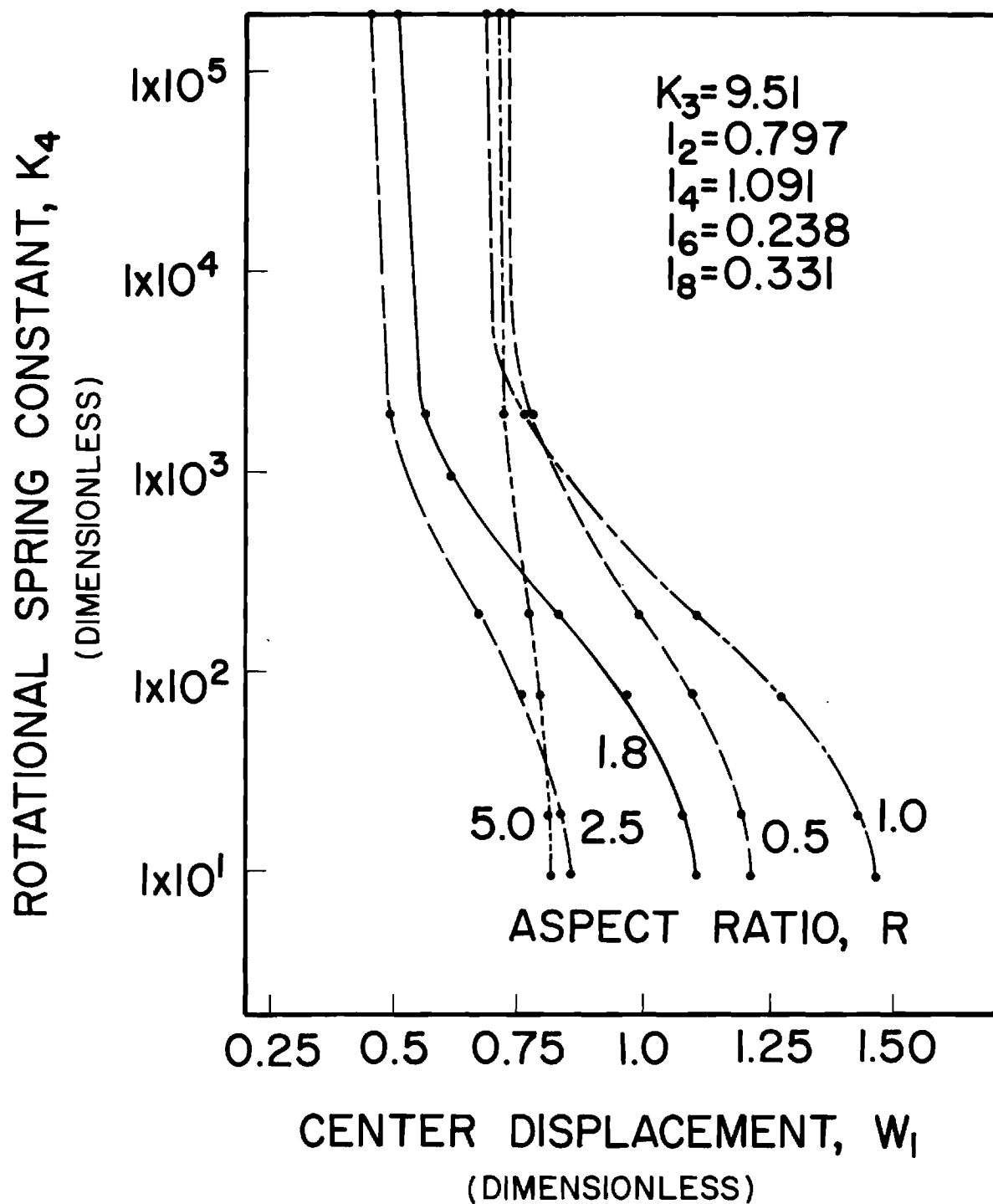


Figure 3.3-8 Static Displacement of Panels with Uniform Loading Versus Rotational Spring Constant on Boundary

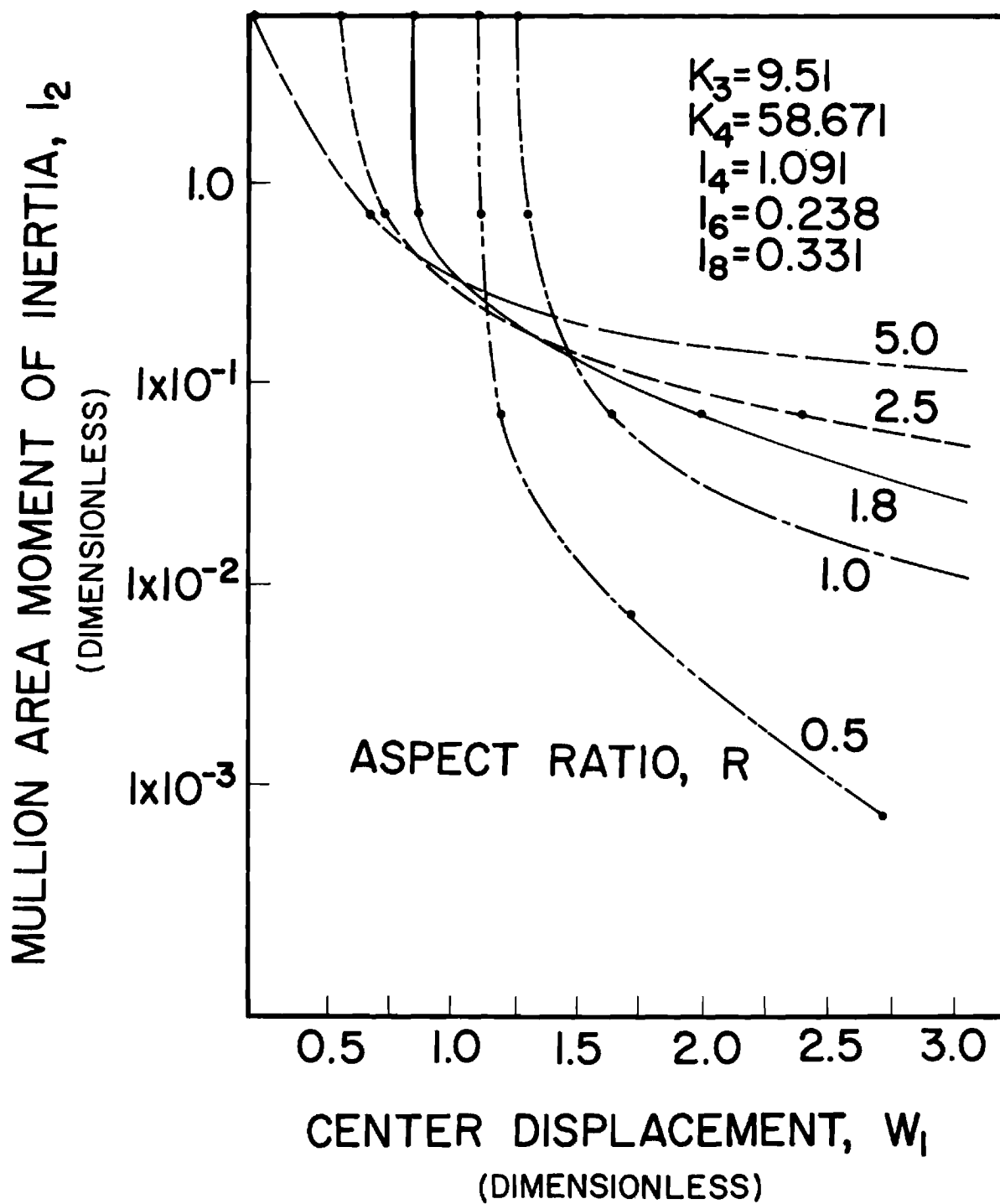


Figure 3.3-9 Static Displacement of Panels with Uniform Loading Versus Mullion Moment of Inertia

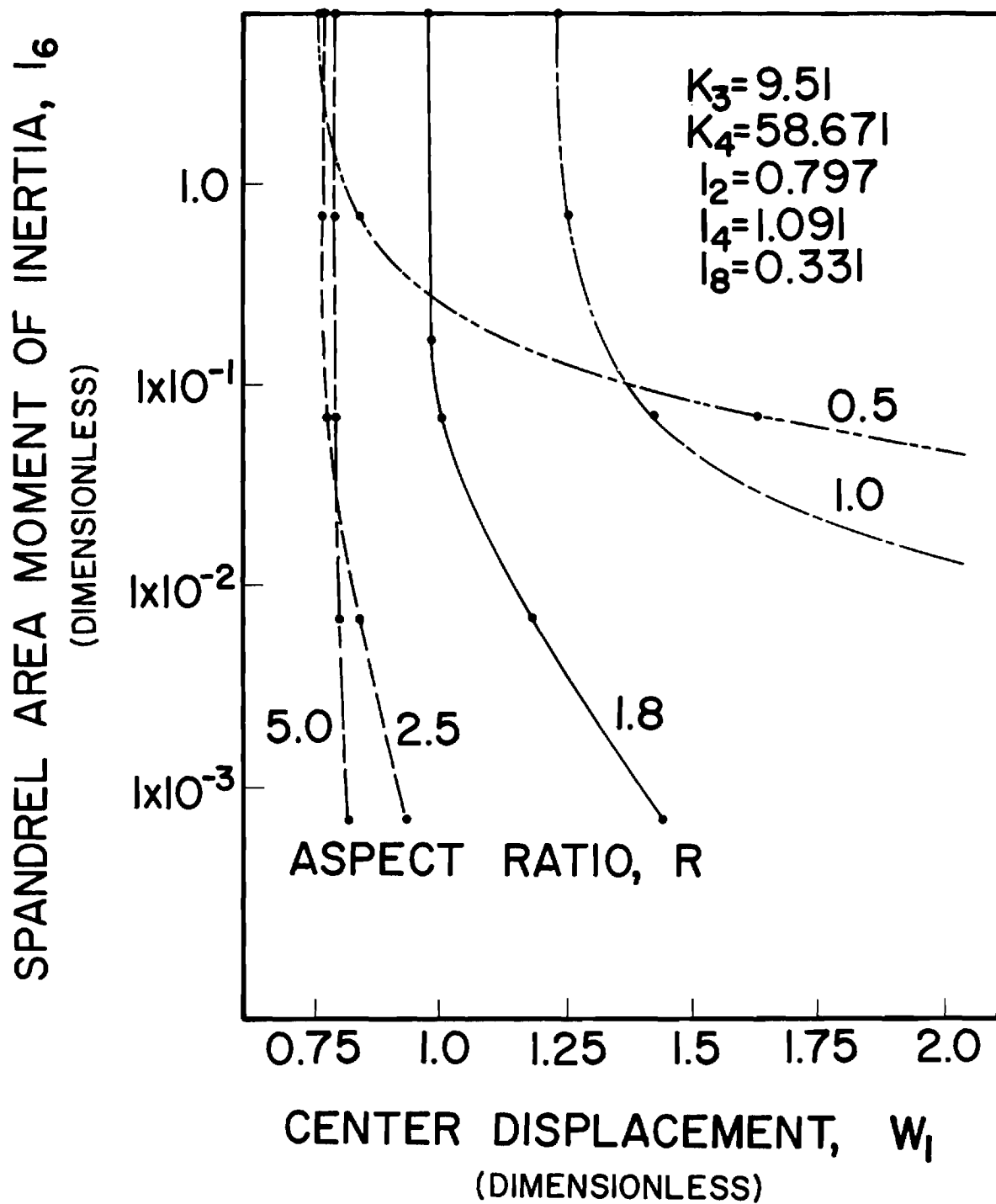


Figure 3.3-10 Static Displacement of Panels with Uniform Loading Versus Spandrel Moment of Inertia

spring stiffness in the range between 1.0 and 10.0 in which  $K_3$  was measured for the prototype structure. Panel response does vary with panel aspect ratio, however, with the largest displacement ratio  $W_1$  recorded for square windows ( $R = 1.0$ ).

All structures except those of aspect ratio 5.0 are sensitive to change in rotational spring constant  $K_4$  as shown in Figure 3.3-8. As the boundary rotational spring stiffness increases, the static displacement at the center of the window decreases. An increase in  $K_4$  is most effective for square ( $R = 1.0$ ) and nearly square ( $R = 0.5, 1.8$ ) structures. When  $K_4$  approaches 1000, panel displacement drops below that of a simply-supported equivalent plate for all aspect ratios.

Figures 3.3-9 and 3.3-10 show structure displacement versus area moment of inertia for mullion and spandrel elements, respectively. Mullion elements which are stiff in flexure reduce static displacements of high aspect ratio panels most effectively since mullion elements are parallel to the long side of the panel (Figure 3.3-9). Figure 3.3-10 demonstrates that spandrel elements are most effective on low aspect ratio windows for the same reason. In both cases, square window-frame structures experience the largest center displacement relative to a simply-supported plate with the same properties, dimensions, and loading.

#### 3.3(d) Dynamic Analysis Results

Dynamic sensitivity of the model is displayed in Figures 3.3-11 through 3.3-18 in which the effect of variation of model properties on fundamental frequency is recorded. Figure 3.3-11 demonstrates that fundamental frequency is insensitive to change in normal spring stiffness except for very low values of  $K_3$ . At the same time, an increase in rotational spring

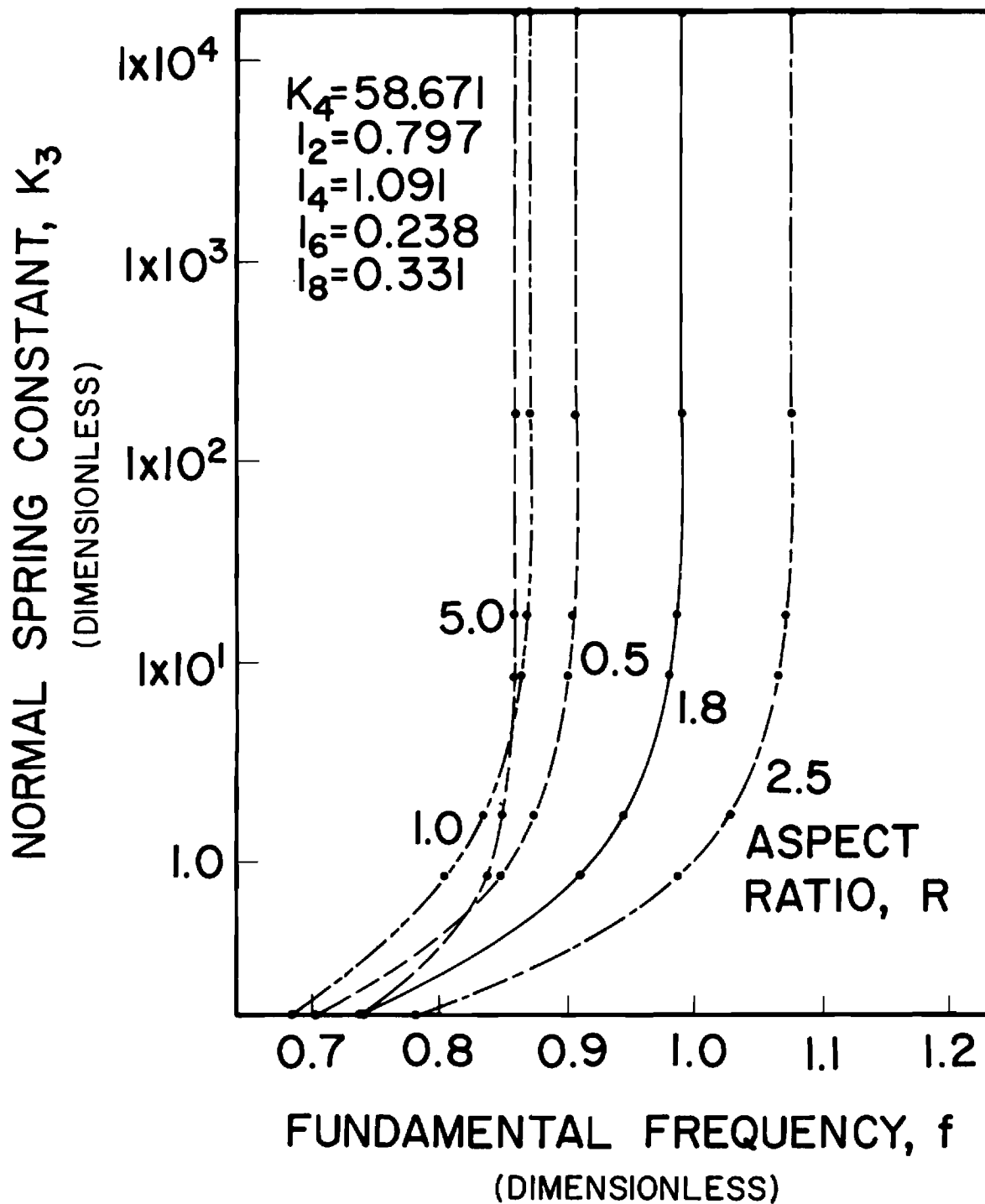


Figure 3.3-11 Panel Fundamental Frequency Versus Normal Spring Constant  $K_3$  on Boundary

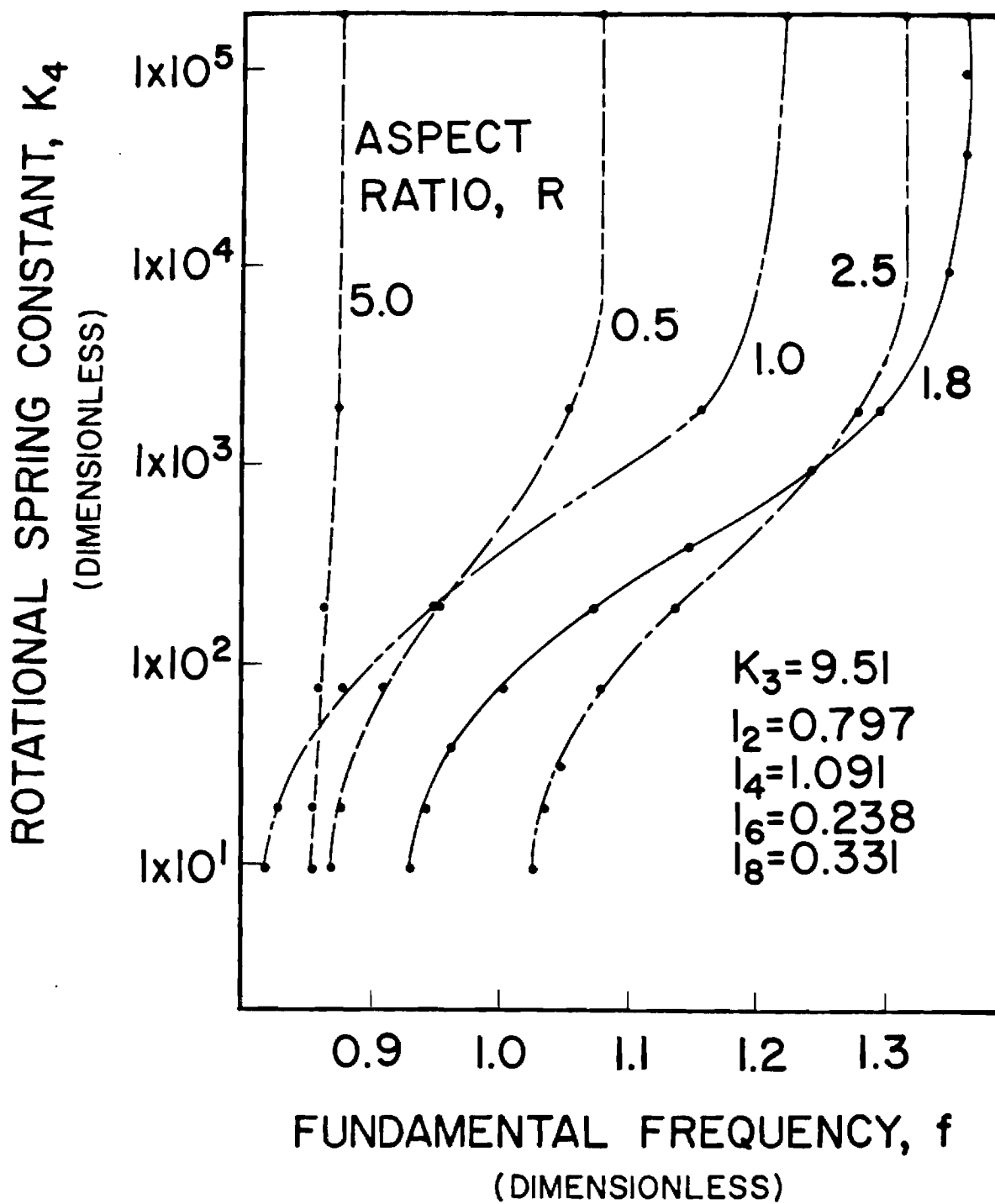


Figure 3.3-12 Panel Fundamental Frequency Versus Rotational Spring Constant  $K_4$  on Boundary

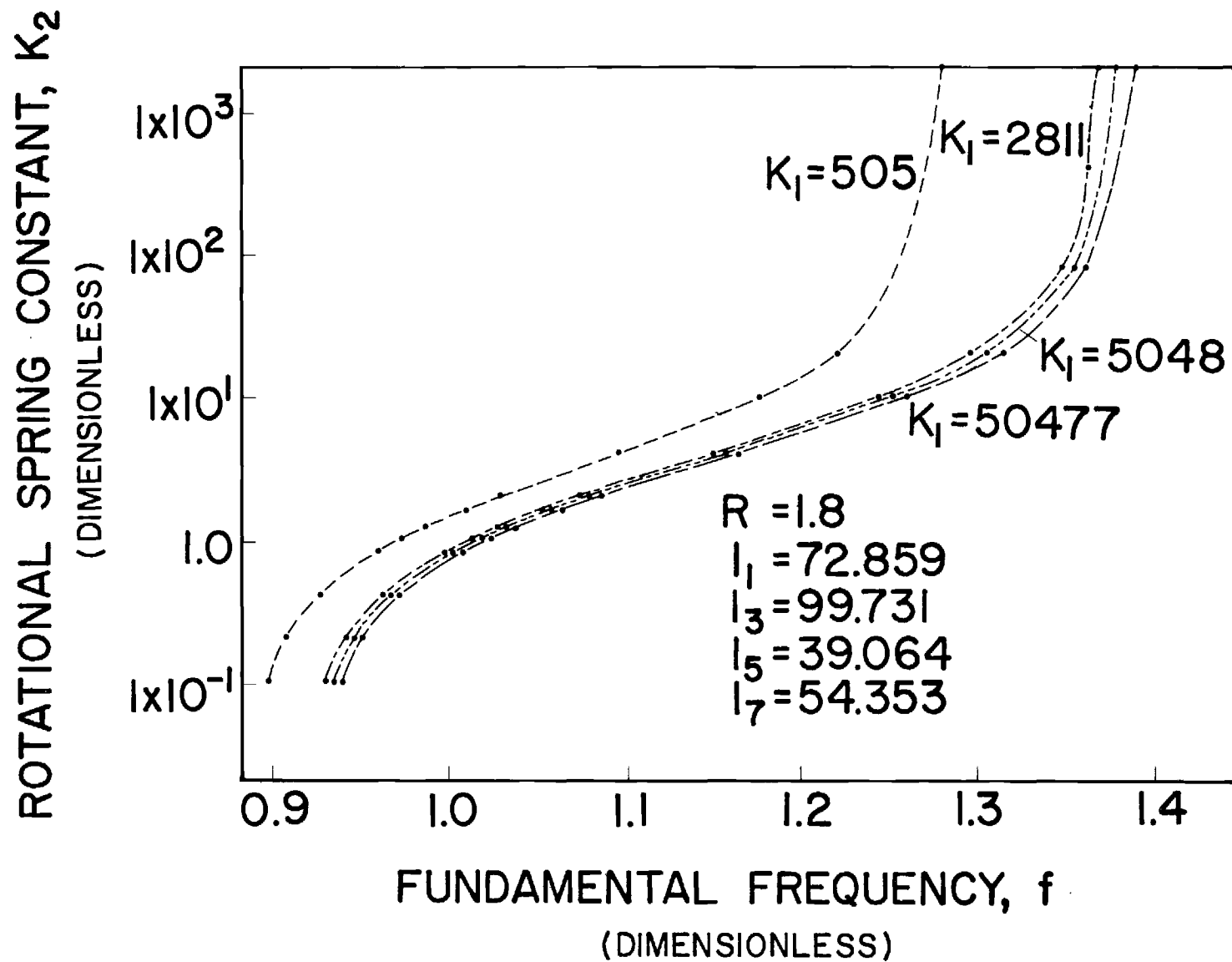


Figure 3.3-13 Panel Fundamental Frequency Versus Rotational Spring Constant  $K_2$  on Boundary

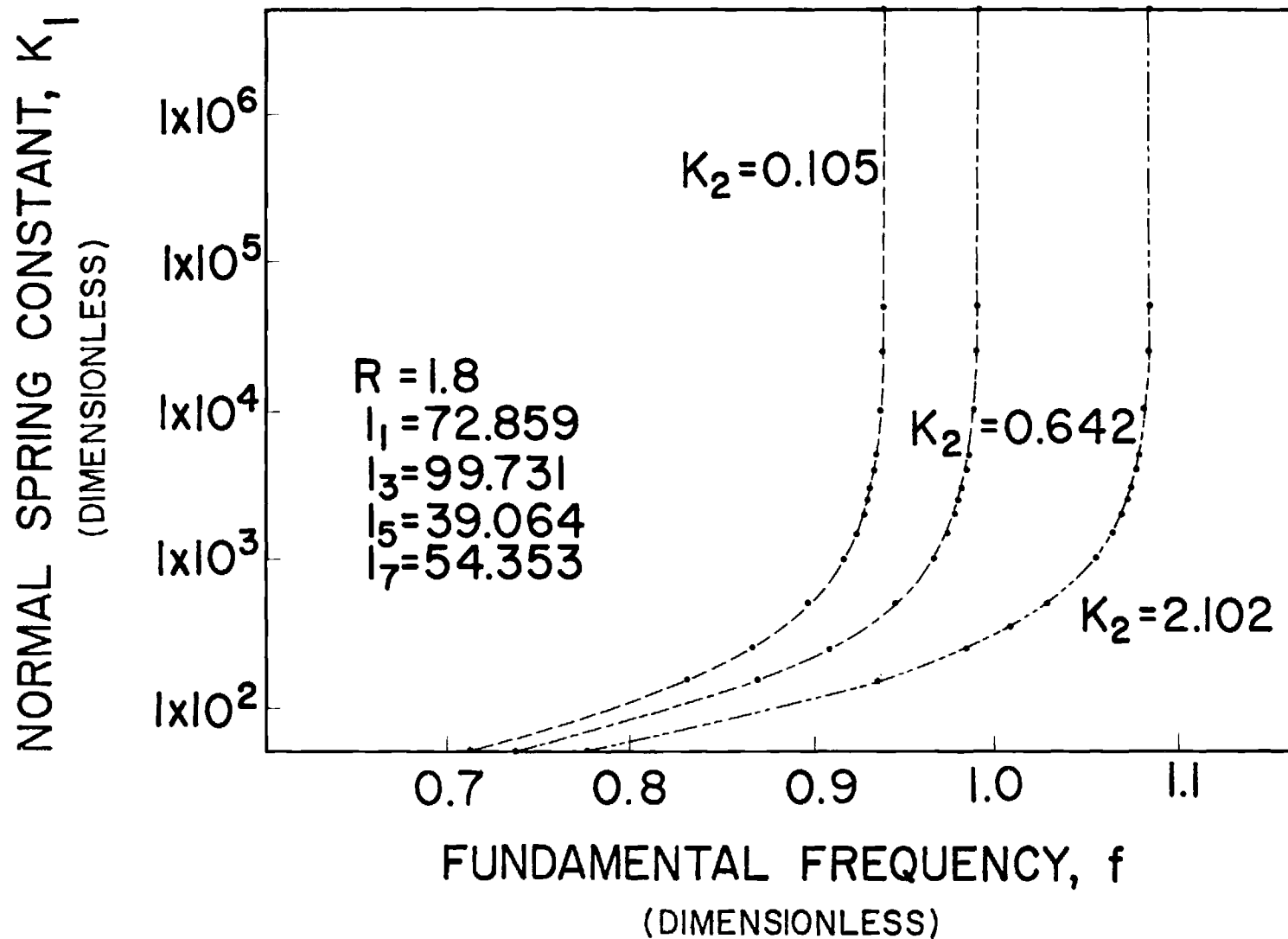


Figure 3.3-14 Panel Fundamental Frequency Versus Normal Spring Constant  $K_1$  on Boundary



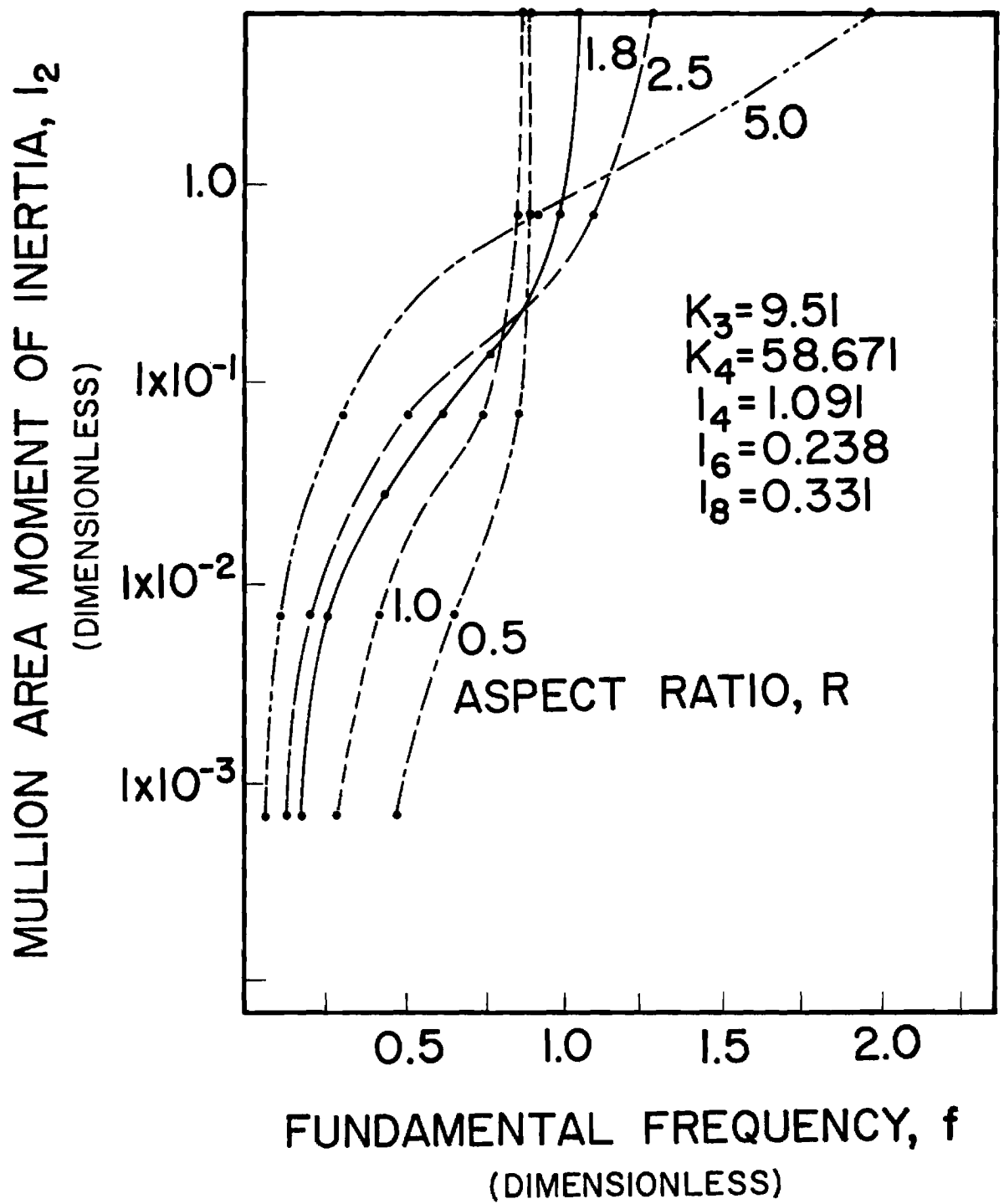


Figure 3.3-15 Panel Fundamental Frequency Versus Mullion Moment of Inertia

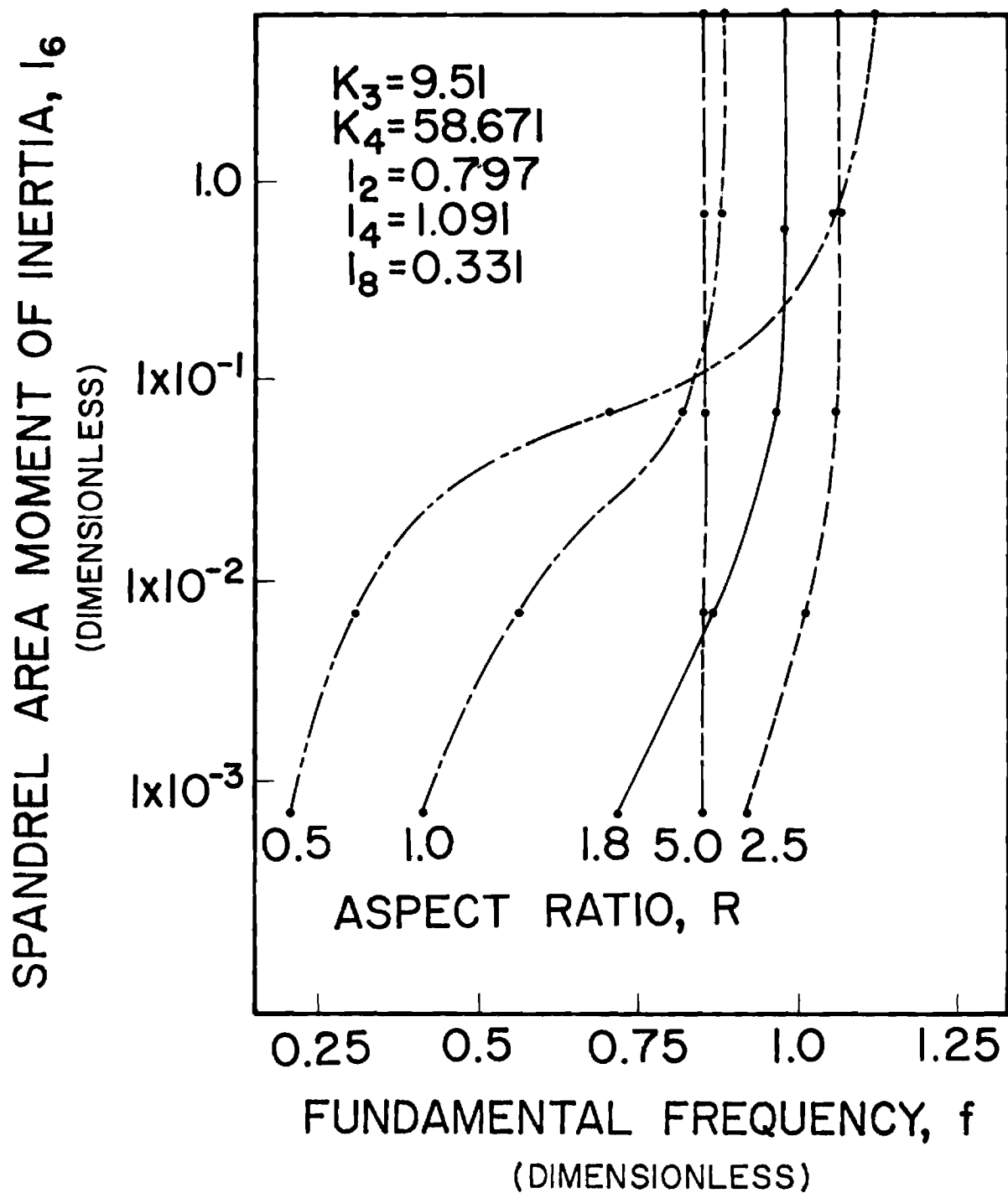


Figure 3.3-16 Panel Fundamental Frequency Versus Spandrel Moment of Inertia

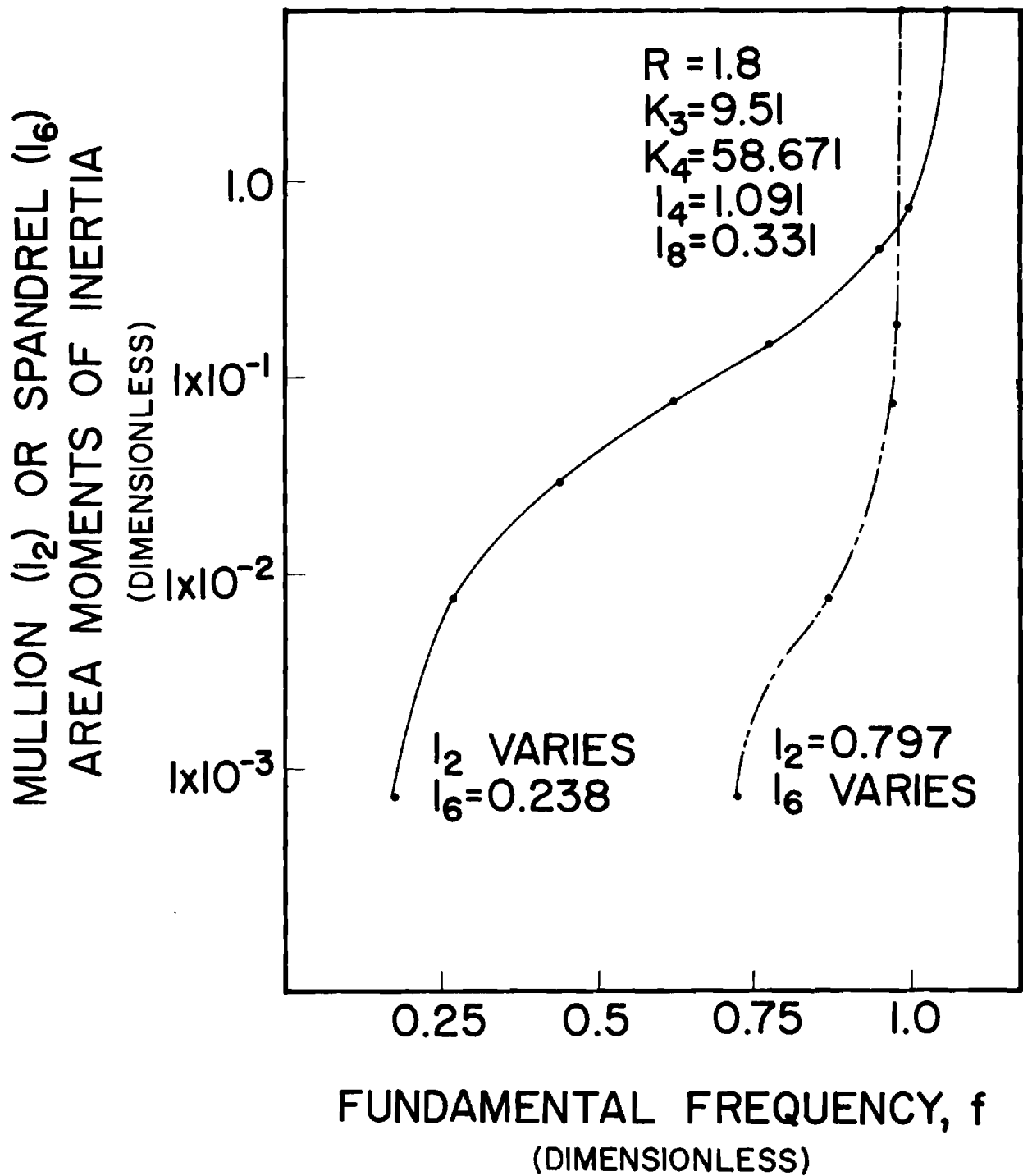


Figure 3.3-17 Panel Fundamental Frequency Versus Mullion and Spandrel Moments of Inertia

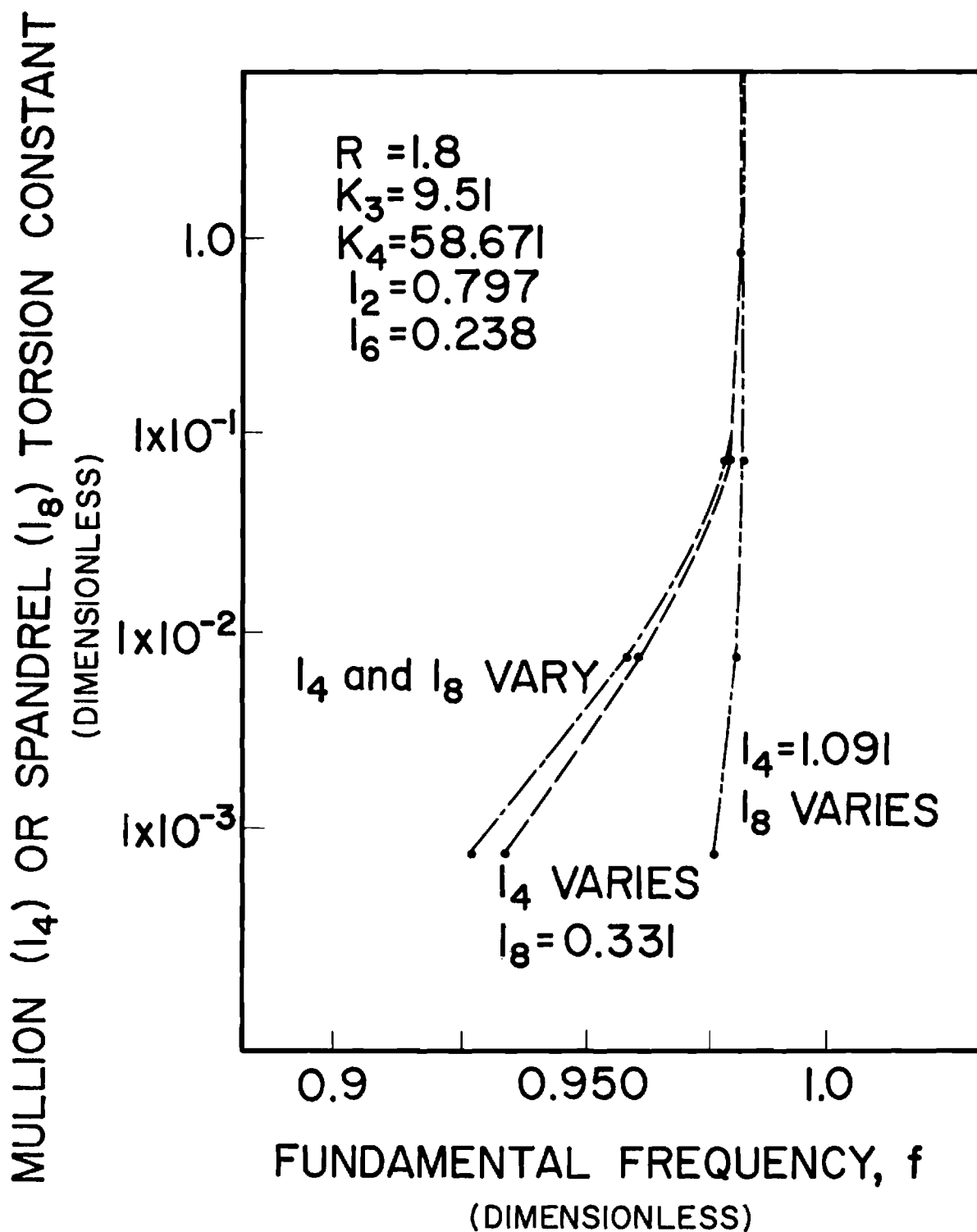


Figure 3.3-18 Panel Fundamental Frequency Versus Mullion or Spandrel Torsion Constant

stiffness produces a notable increase in fundamental frequency with  $R = 1.8$  panels experiencing the greatest variation as shown in Figure 3.3-12. However, as observed earlier in Figure 3.3-8 for the static case,  $R = 5.0$  panels are virtually unaffected by a change in  $K_4$ . The greater sensitivity of frequency to rotational rather than normal spring stiffness is also shown in Figure 3.3-13 for the  $R = 1.8$  case. Variation in  $K_2$  is much more effective in changing  $f$  than a similar change in  $K_1$  and the influence of  $K_1$  diminishes rapidly at higher values. However, Figure 3.3-14 shows that increasing normal spring stiffness parameter  $K_1$  in the region below  $K_1 = 1 \times 10^4$  results in greater increase in frequency for larger values of rotational spring constant  $K_2$ . Figure 3.3-14 applies to the  $R = 1.8$  case only, for the numerical values of frame inertia parameters defined in the figure.

Figures 3.3-15 and 3.3-16 display change in fundamental frequency versus mullion and spandrel area moments of inertia, respectively. Mullion flexural stiffness which increases with value of  $I_2$  especially augments frequencies for high aspect ratio panels where the mullion borders the long side of the panel. In Figure 3.3-16, stiff spandrel elements increase the frequency of low aspect ratio panels for the same reason. Finally, in Figure 3.3-17, the effect of change in either mullion ( $I_2$ ) or spandrel ( $I_6$ ) area moment of inertia on fundamental frequency of  $R = 1.8$  panels are shown on the same plot. For this aspect ratio and the particular window-frame model used (see Figure 3.3-1), variation in  $I_2$  is much more effective in changing structure frequency than variation in  $I_6$ .

Figure 3.3-18 displays the effect of mullion and spandrel torsional stiffness on fundamental frequency of vibration for  $R = 1.8$  structures.

While torsional stiffness of frame elements has very little effect on frequency, mullion stiffness ( $I_4$ ) is relatively more important than spandrel stiffness ( $I_8$ ) for this aspect ratio.

Dynamic Response: The dynamic response of the window-frame symmetry model in Figure 3.3-1 was computed for a number of forcing functions, one of which is presented here. An exhaustive study encompassing many different types of excitations was not made although a wide range of loading types can be handled by the analytical model. As an example, panel response to in-plane racking distortions was determined and is presented in the next section. In general, the limited number of cases studied served to confirm results presented in the dimensionless plots above; for example, the model's sensitivity to gasket rotational stiffness was evident in the dynamic response studies.

The pressure signal in Figure 3.3-19 was recorded in the laboratory during dynamic testing and was assumed to be a representative pressure fluctuation-time history which could be used to investigate the dynamic sensitivity of the system. The pressure was assumed to be uniformly distributed over the surface of the panel. Attention was confined to equivalent single plate models with properties listed in Table 3.3-1, thickness of 0.298 in. (0.76 cm),  $R = 1.8$ , and  $L_e = 0$ .

As stated above, the model proved to be more sensitive to rotational spring stiffness than to normal spring stiffness as Figures 3.3-20 and 3.3-21 show. In these figures, dimensionless normal spring constant  $K_1$  and rotational spring constant  $K_2$  are plotted versus dimensionless maximum dynamic response  $W_2$  at the center of the panel due to the excitation in Figure 3.3-19. The maximum dynamic response of a twelve-finite element

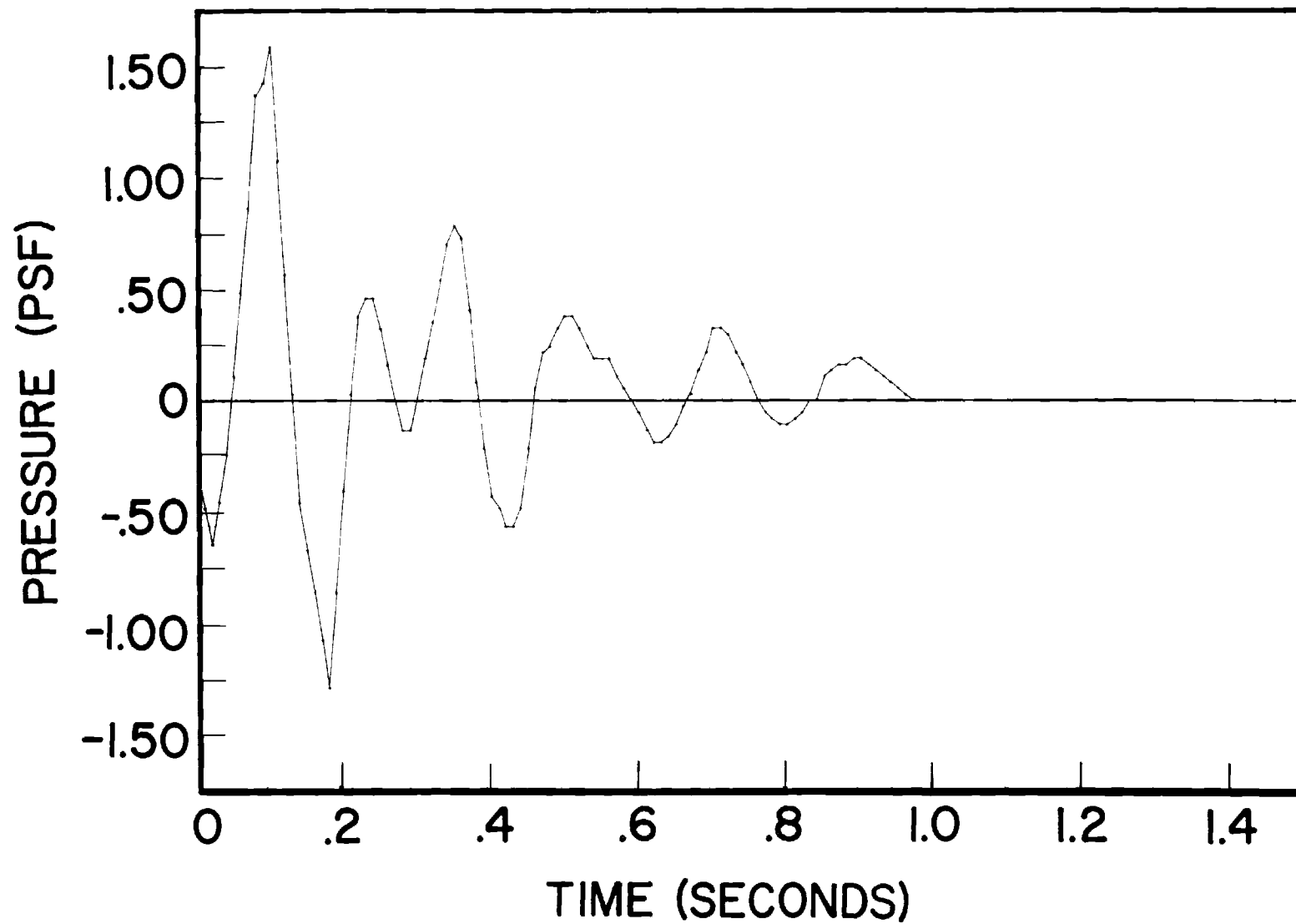


Figure 3.3-19 Uniform Pressure Excitation

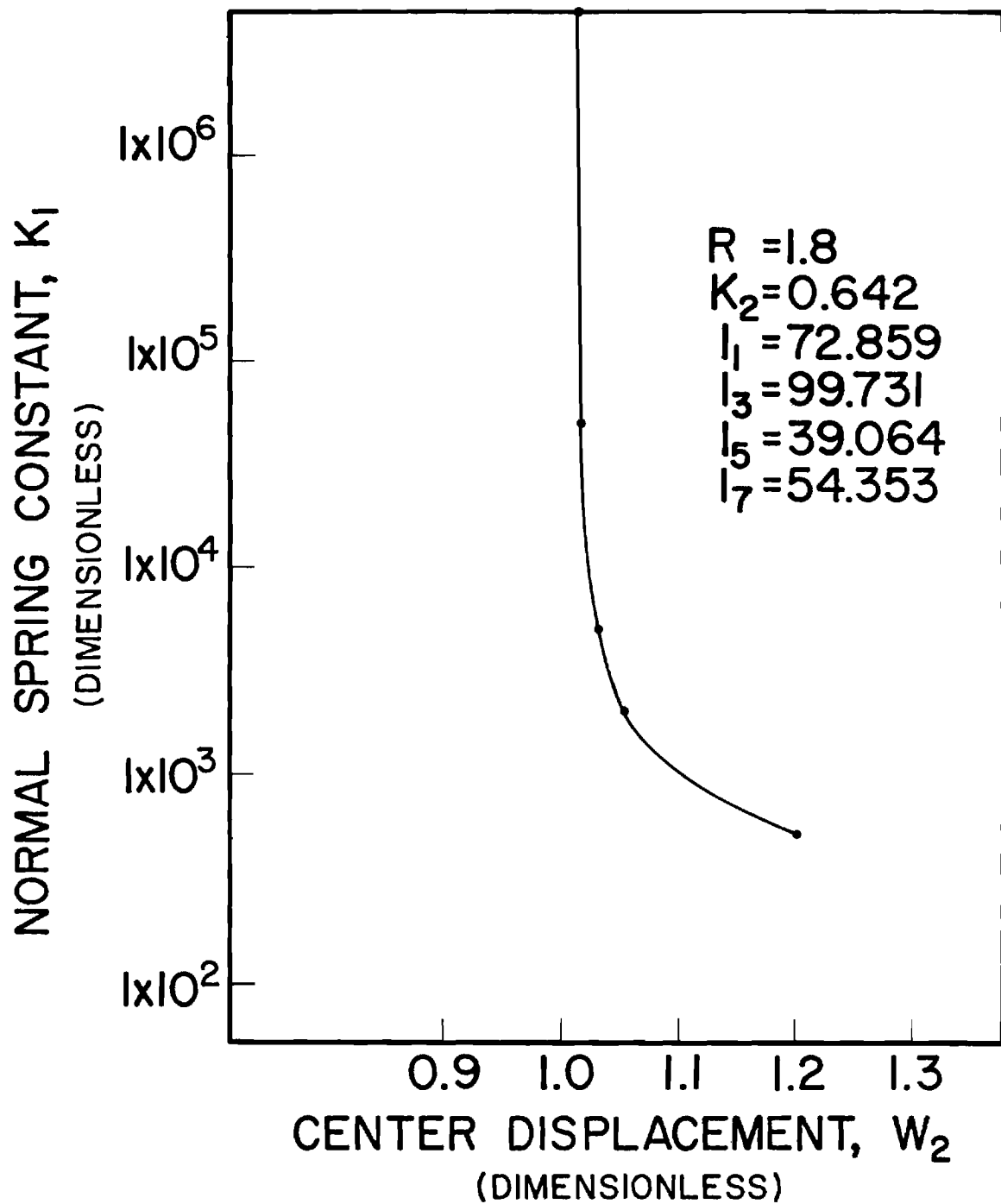


Figure 3.3-20 Maximum Dynamic Response to Pressure Excitation Versus Normal Spring Constant on Boundary



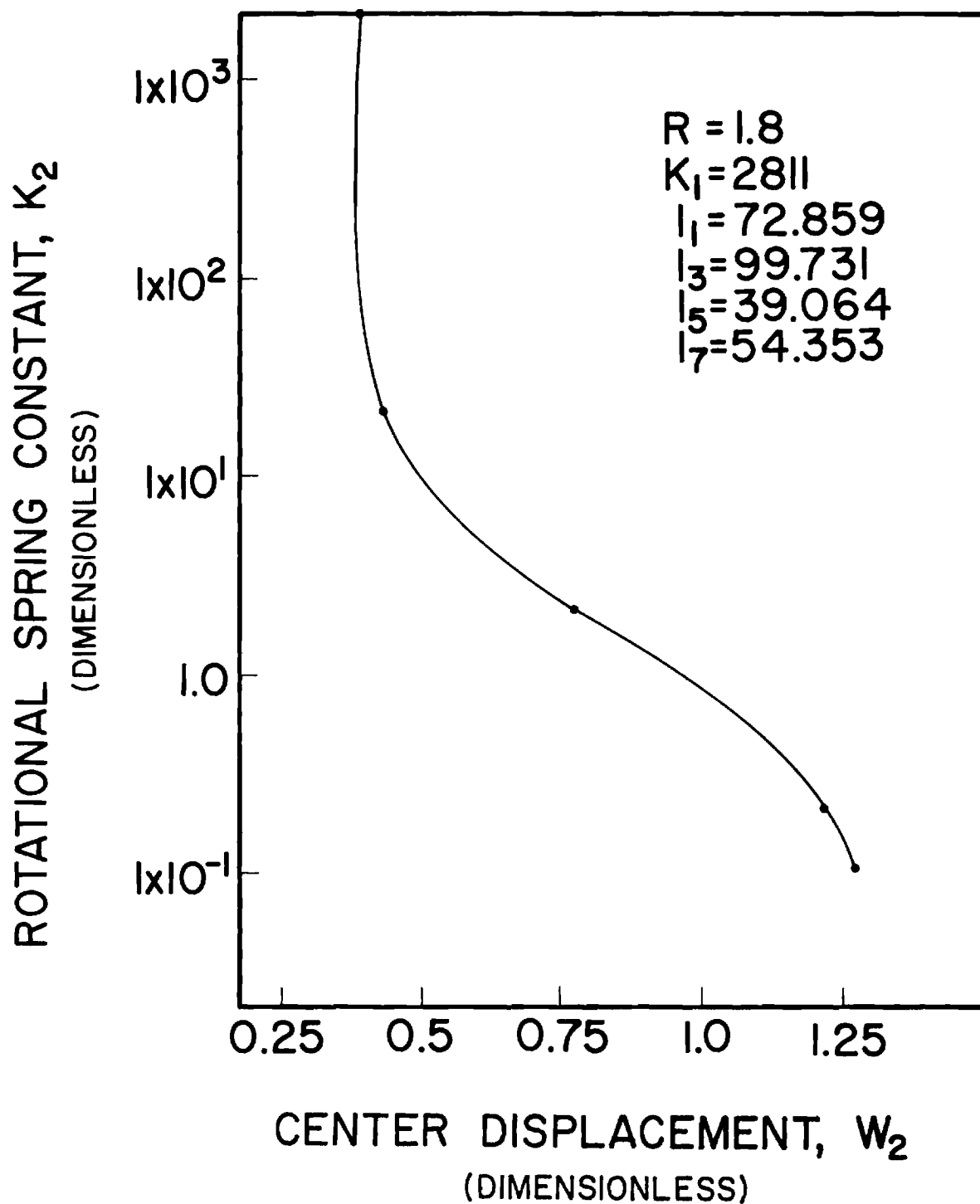


Figure 3.3-21 Maximum Dynamic Response to Pressure Excitation Versus Rotational Spring Constant on Boundary

model of a simply-supported quarter plate (similar to Figure 3.3-1) to this excitation was used to nondimensionalize these plots; its response-time history is recorded in Figure 3.3-22.

The effectiveness of soft ( $k_R = 50 \text{ in-lb/radian/in} = 222 \text{ N}\cdot\text{m/radian/m}$ ) and hard ( $k_R = 1 \times 10^6 \text{ in-lb/radian/in} = 4.45 \times 10^6 \text{ N}\cdot\text{m/radian/m}$ ) rotational springs in reducing panel response at the center is also shown in time history plots, Figures 3.3-23 and 3.3-24. Figures 3.3-25 and 3.3-26, however, show that soft ( $k_N = 100 \text{ lb/in/in} = 690 \text{ kN/m/m}$ ) and hard ( $k_N = 1 \times 10^6 \text{ lb/in/in} = 6.9 \times 10^6 \text{ kN/m/m}$ ) normal springs result in response-time histories which differ only slightly in comparison.

### 3.3(e) Panel Racking Studies

Differential motions of adjacent floors in a highrise building induced by wind or earthquake introduce in-plane racking distortions into the window-framework structure. The analytical model presented in this report can be used to study the linear small displacement behavior of the system. The finite element model of the full panel and frame between floor levels shown in Figure 3.3-27 was developed to study displacement response but a finer mesh would be needed for detailed stress investigations. The support boundary conditions are modified slightly compared to the symmetry model of Figure 3.3-1 to more closely represent actual points of support in the prototype structure. The frame is simply-supported at lower corner points and at the end of mullion element extensions of length  $Y_m$ . Two values of  $Y_m$  are considered to provide a range of natural frequencies of vibration for in-plane motion:  $Y_m = 5 \text{ in. (127 mm)}$  represents the laboratory test rig support conditions and  $Y_m = 50 \text{ in. (1270 mm)}$  approximates the support conditions in the field study building.

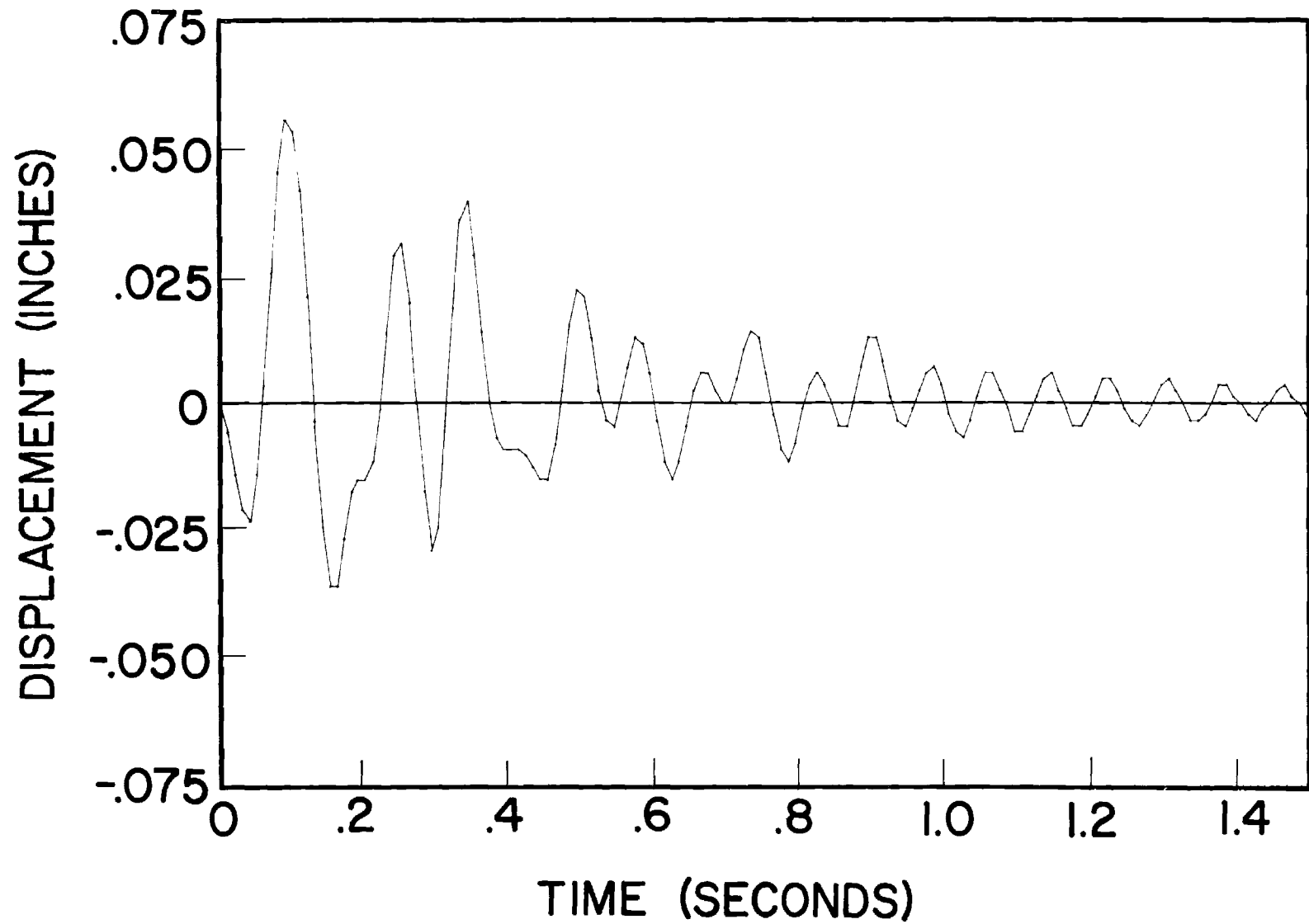


Figure 3.3-22 Dynamic Response of Simply-Supported Plate to Uniform Pressure Excitation

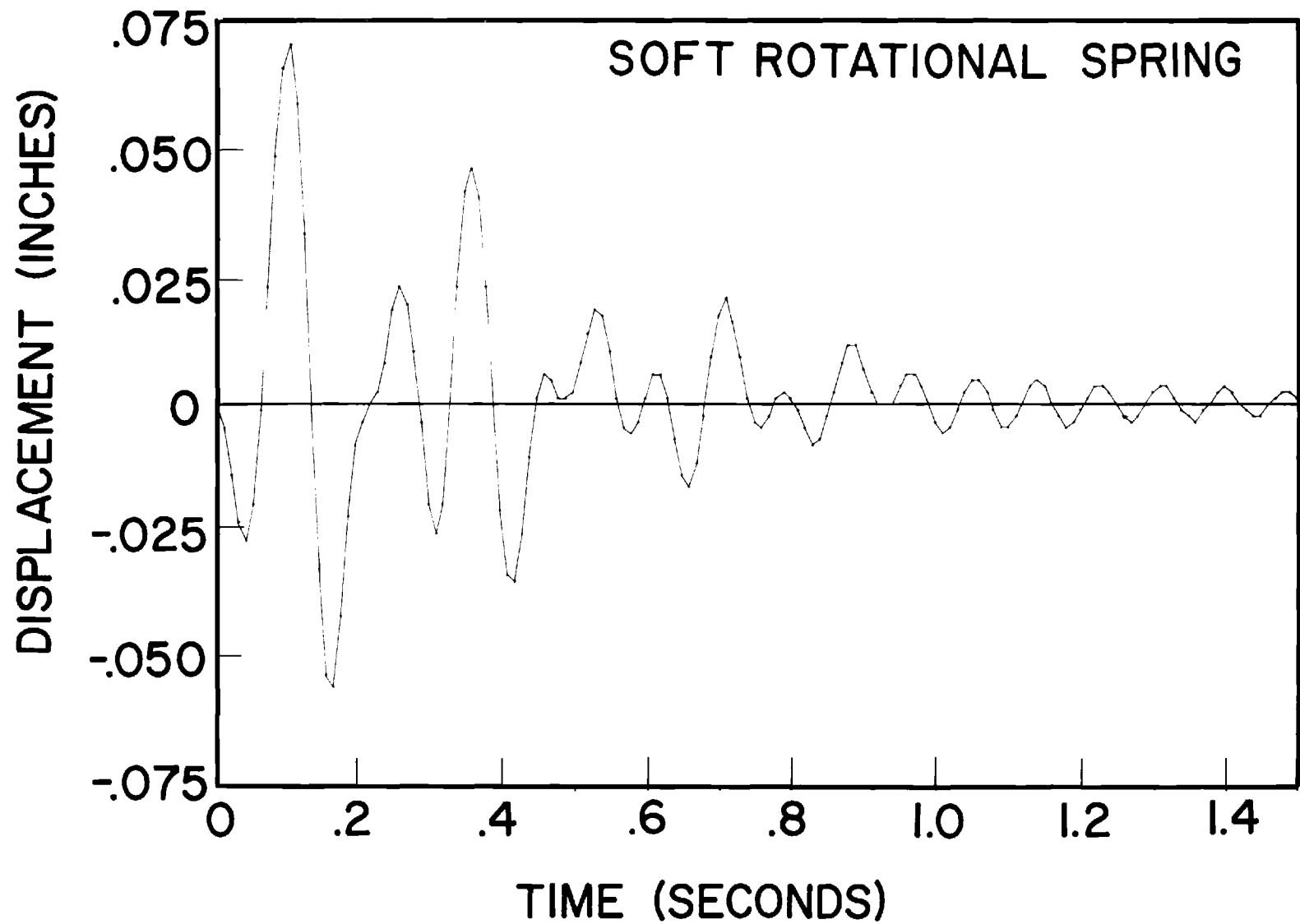


Figure 3.3-23 Dynamic Response of Panel to Uniform Pressure Excitation with Soft Rotational Spring on Boundary

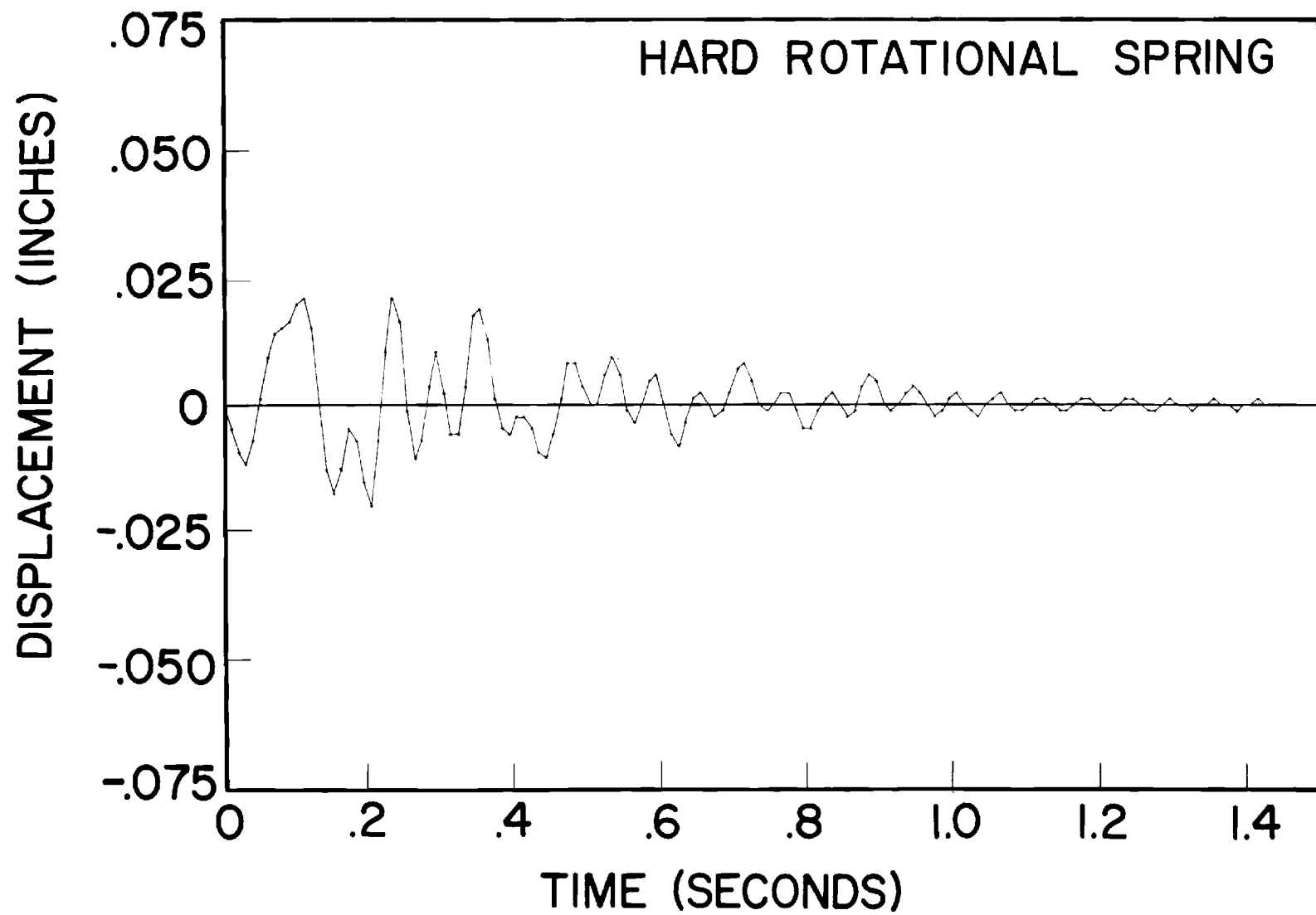


Figure 3.3-24 Dynamic Response of Panel to Uniform Pressure Excitation with Hard Rotational Spring on Boundary

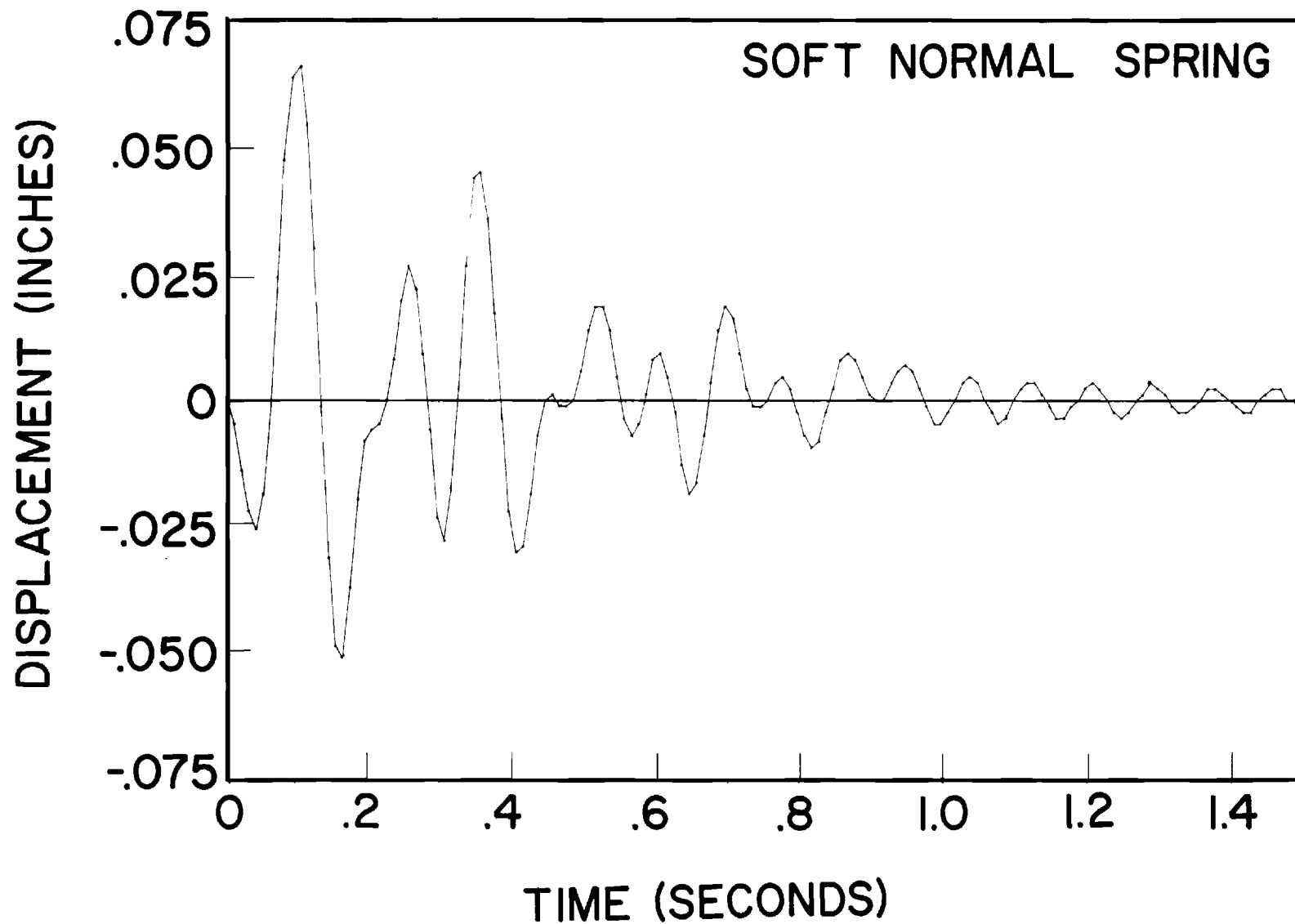


Figure 3.3-25 Dynamic Response of Panel to Uniform Pressure Excitation with Soft Normal Spring on Boundary

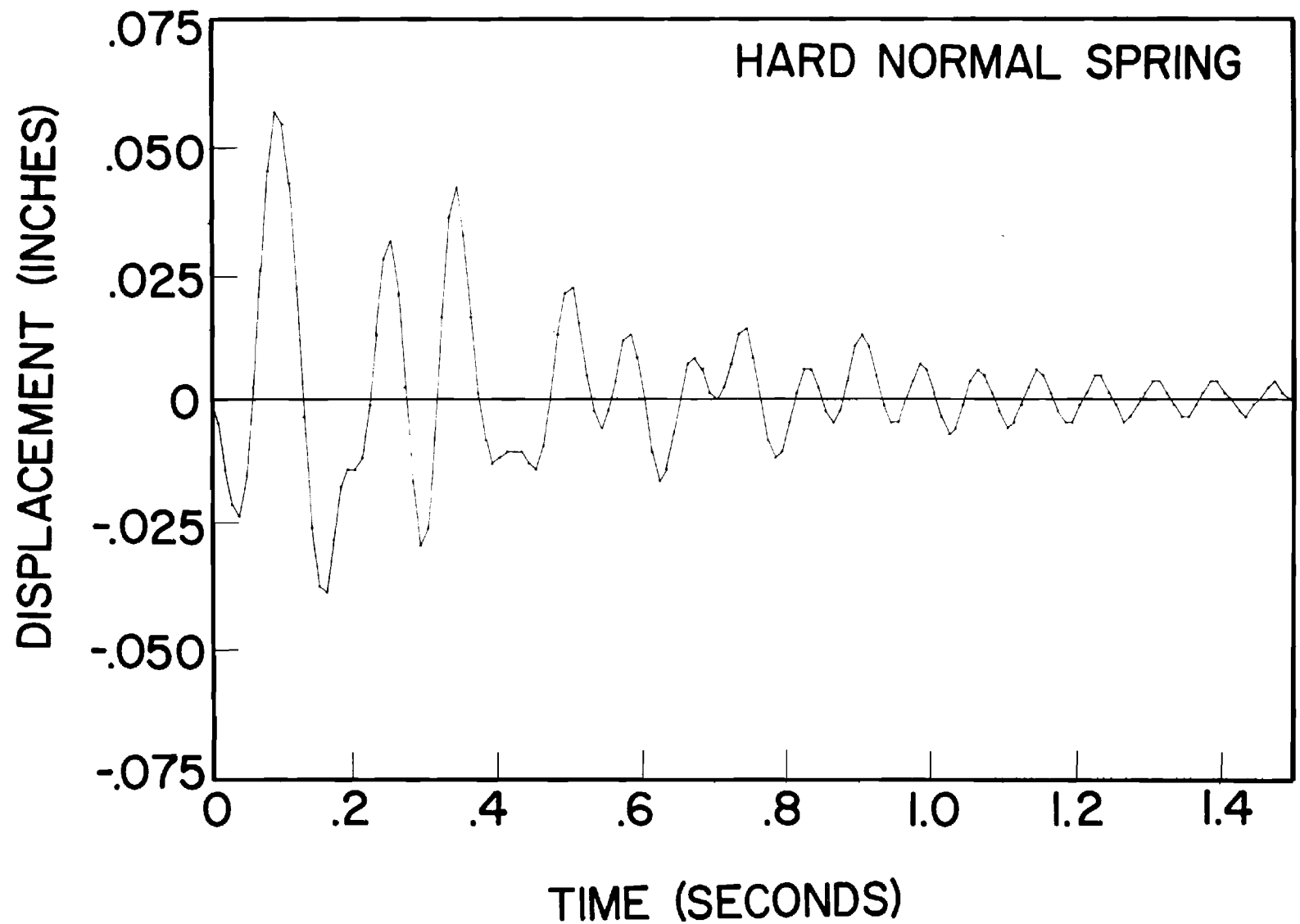


Figure 3.3-26 Dynamic Response of Panel to Uniform Pressure Excitation with Hard Normal Spring on Boundary





An equivalent single panel of thickness 0.298 in. (0.76 cm) and the remaining frame and panel properties in Table 3.3-1 were used for all studies. A consistent mass formulation of the structure mass matrix was chosen and frequencies of vibration determined for the two values of  $Y_m$  listed above. Frequencies and mode shapes for the first three modes are summarized in Figures 3.3-28, 3.3-29, and 3.3-30. First and third mode frequencies and vibration shapes are seen to depend on the value of mullion extension length  $Y_m$  while second mode results are unaffected by  $Y_m$ . All three modes are boundary spring and frame deformation modes and involve very little panel deformation.

To demonstrate the capabilities of the analytical model and to provide some indication of the level and type of response to be expected during racking motions, building response recorded in a moderate earthquake was used to excite the model. The supports at the upper floor level only of the model with  $Y_m = 50$  in. (1270 mm) were assumed to move in accordance with the horizontal displacement and acceleration traces (Figure 3.3-31) recorded on the sixteenth floor of the Alexander Building in the S81W direction during the 1957 San Francisco, California, earthquake. This earthquake of Richter magnitude 5.3 produced a peak floor acceleration of 0.04 g and peak floor displacement of 0.61 in. (1.55 cm). Direct integration of equation (3.1-59) was used to determine the total response of the model assuming simple modal viscous damping at 2% of critical in all modes.

Thirty seconds of displacement response at model degrees of freedom 1, 17, 18, 24, and 33 (see Figure 3.3-27) are shown in Figure 3.3-32. Peak response of 0.107 inches (0.271 cm) is seen to occur at degree of freedom 33 at the top corner of the panel. Comparison of response at

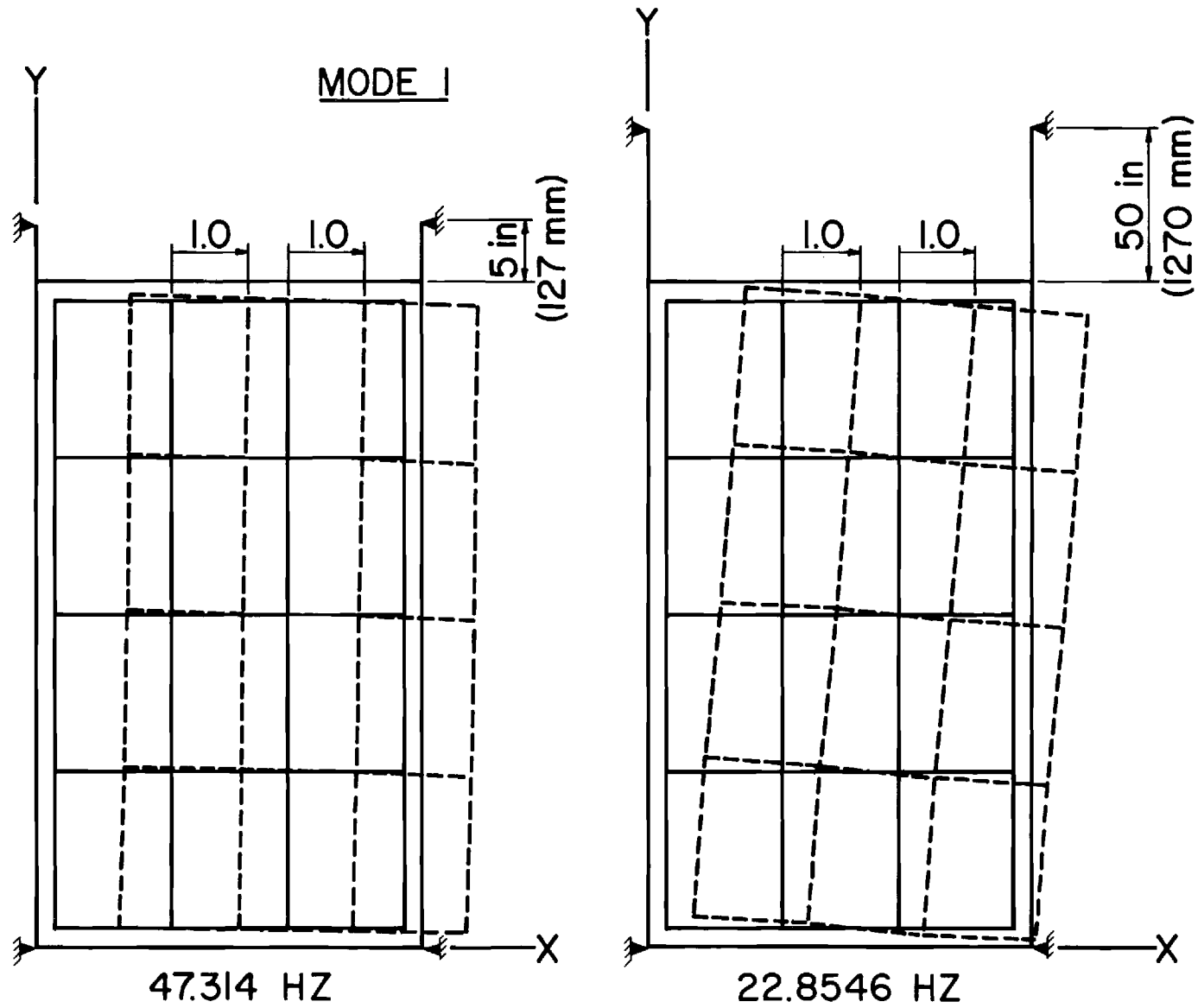


Figure 3.3-28 In-plane Model Frequencies and Mode Shapes, Mode 1

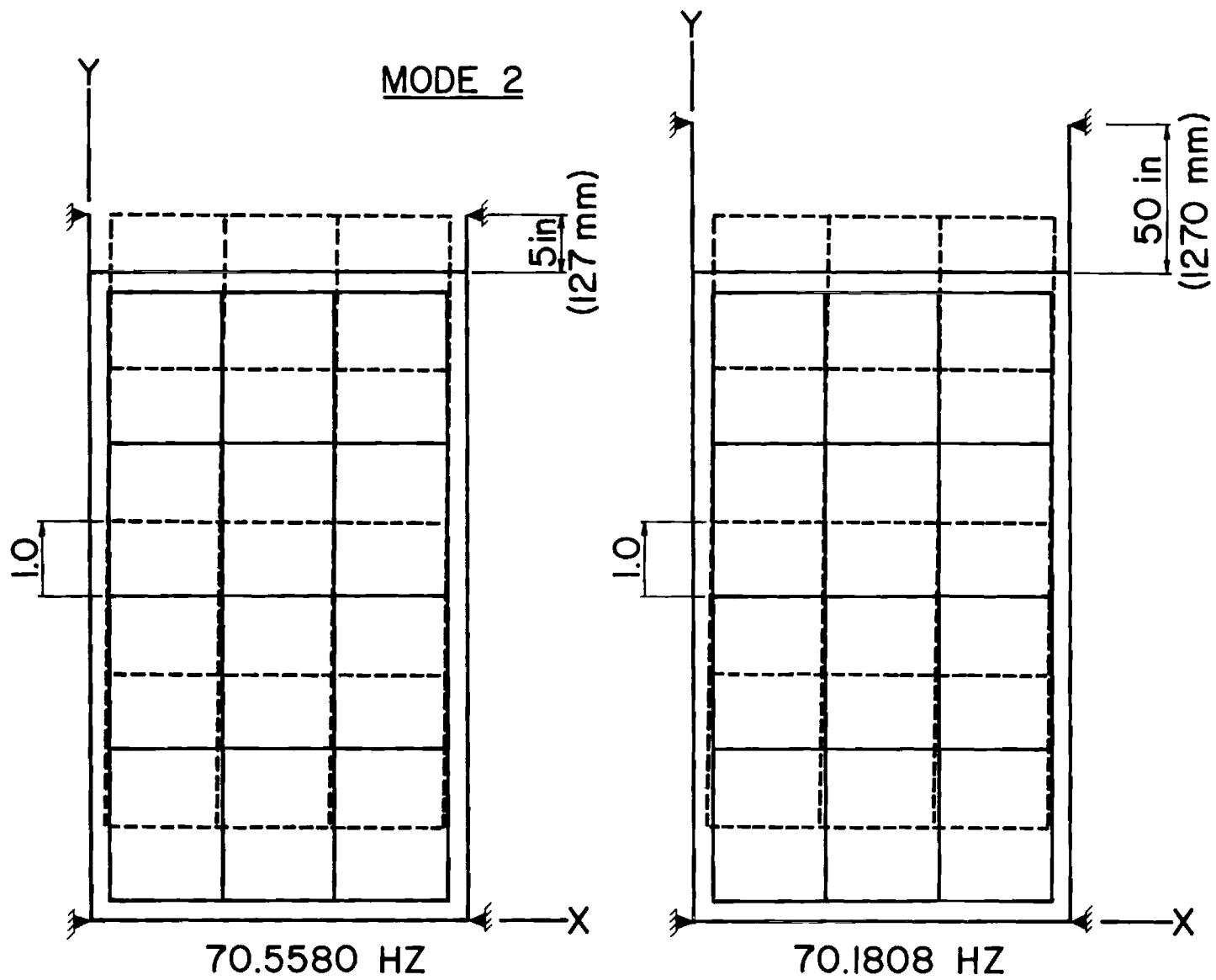


Figure 3.3-29 In-plane Model Frequencies and Mode Shapes, Mode 2

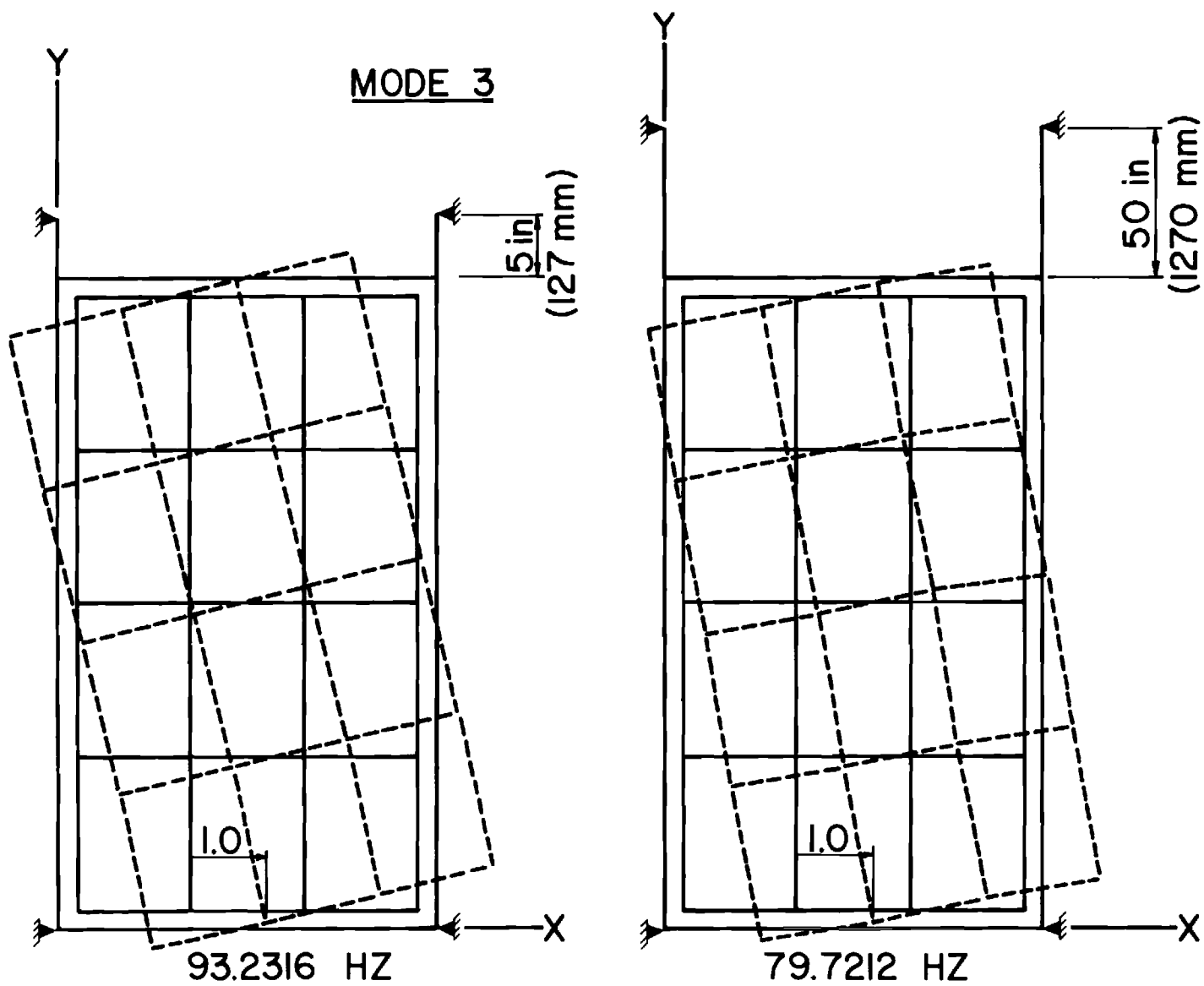


Figure 3.3-30 In-plane Model Frequencies and Mode Shapes, Mode 3

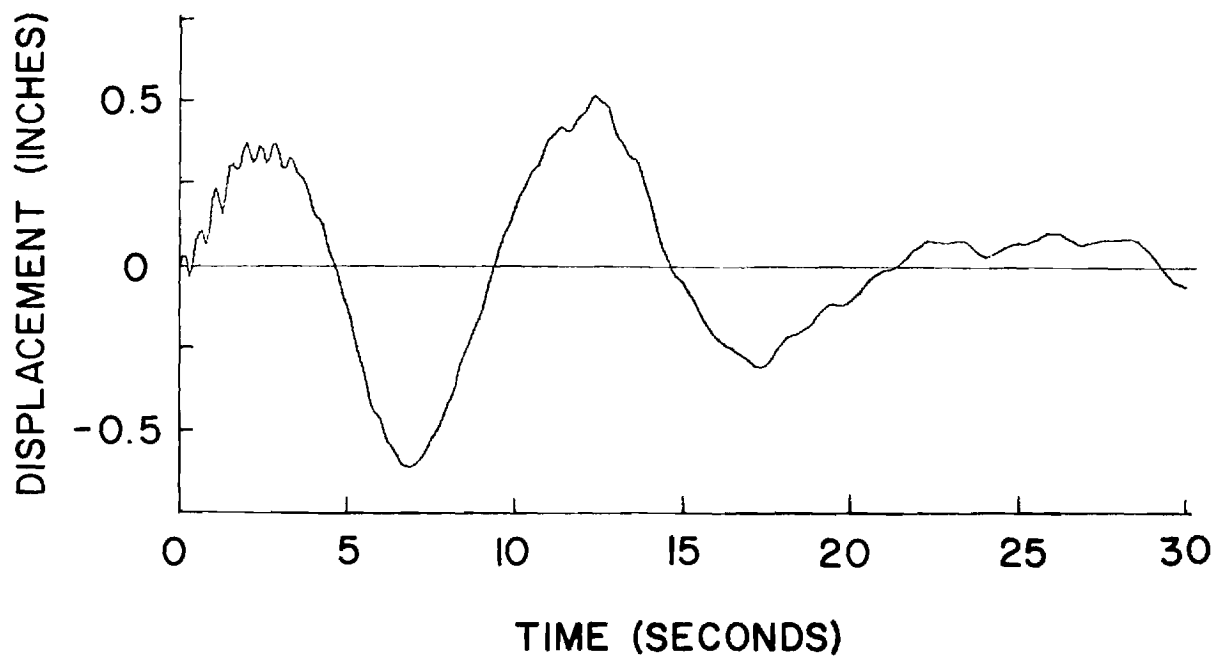
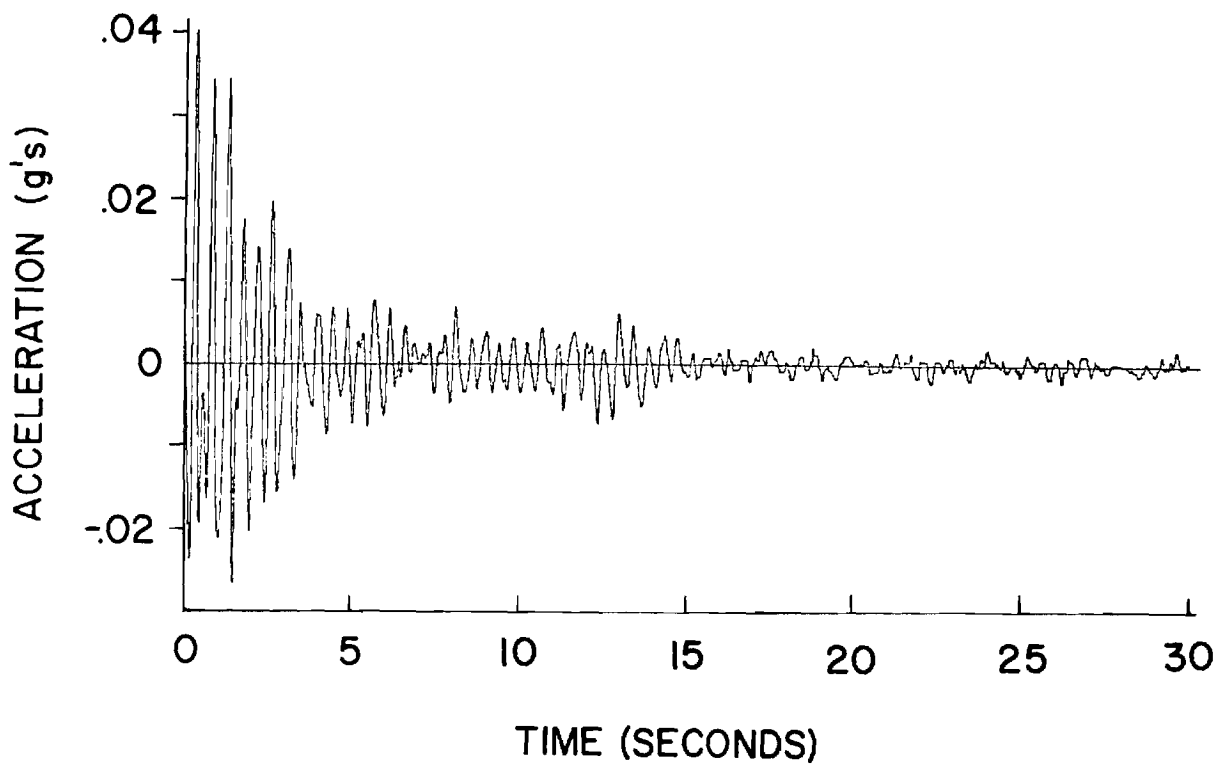


Figure 3.3-31 Acceleration and Displacement Records on 16th Floor of Alexander Building in S81W Direction, 1957 San Francisco, California, Earthquake

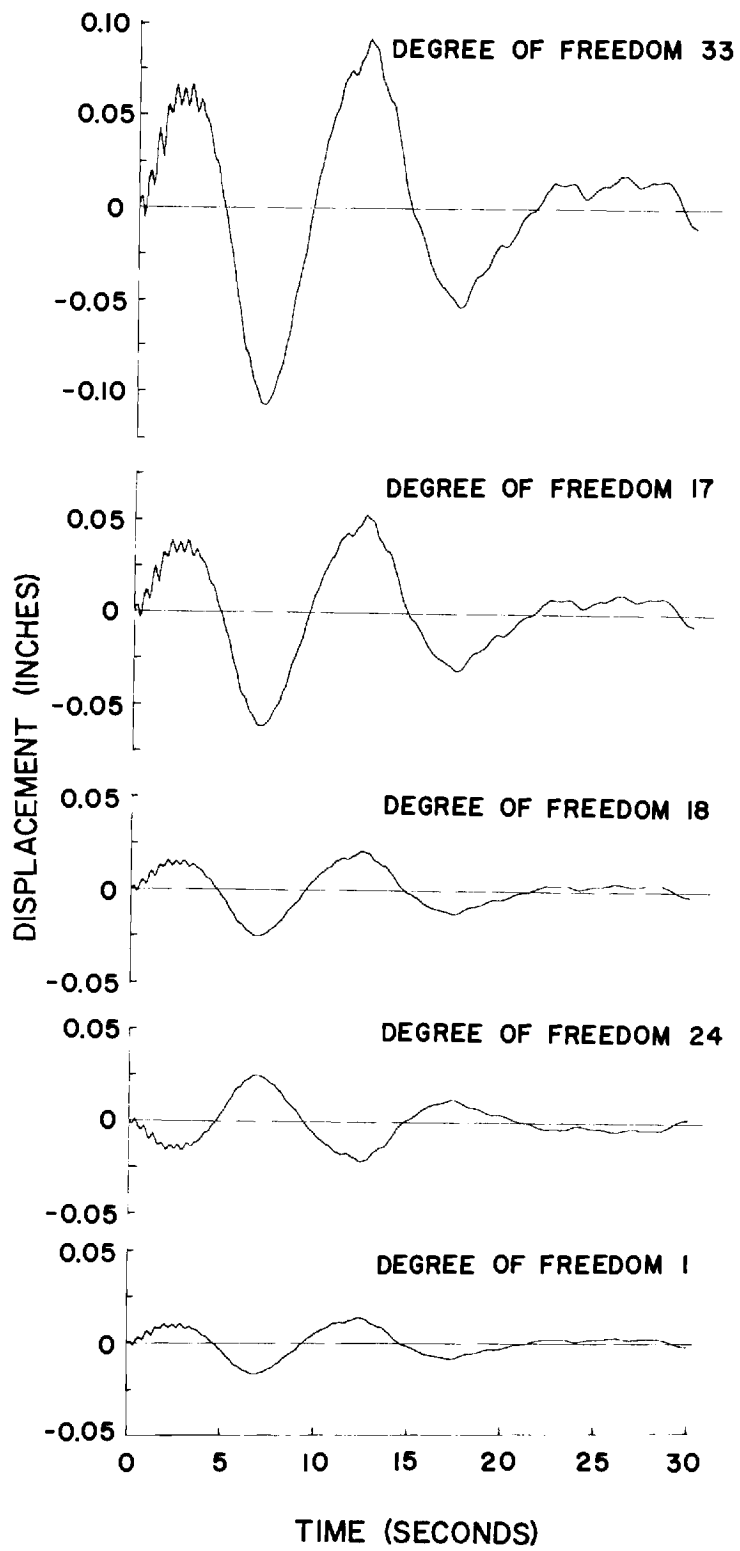


Figure 3.3-32 In-plane Model Displacement Response at Master Degrees of Freedom 1, 17, 18, 24, and 33

degrees of freedom 18 and 24 demonstrates that panel rocking within the frame and gasket assembly is the primary mode of behavior. Panel stress-time histories were not computed.

### 3.4 Double Panel Models

An equivalent single plate model was developed for the static and dynamic sensitivity studies reported in section 3.3. Actual insulating windows in the prototype structure however are composed of two panes of glass separated by a 0.5 inch (1.27 cm) sealed air space. A double plate continuum model and a parallel-plate finite element model were developed to investigate the out-of-phase dynamic properties of the double panel system.

#### 3.4(a) Continuum Model

The double plate continuum model shown in Figure 3.4-1 is based on the assumption that the glass unit is simply-supported within the boundary frame. The metal spacer bars on the periphery of the panel are assumed to introduce rotational coupling between the plates represented by a distributed rotational spring with spring constant  $K_{\theta}$  in units of lb-in/radian/in ( $\text{N}\cdot\text{m}/\text{radian}/\text{m}$ ). A detail of the boundary connection is presented in Figure 3.4-2. A range of values of  $K_{\theta}$ , expressed in terms of nondimensional parameters

$$\xi_B = 2K_{\theta} b/D, \quad \xi_A = \xi_B a/b \quad (3.4-1)$$

was used to approximate varying degrees of interplate connection on the boundary.

In modes of vibration in which volume change occurs, the air gap was modeled as a distributed linear spring with spring constant  $K_{\delta}$  expressed in units of lb/in/in<sup>2</sup> ( $\text{N}/\text{m}/\text{m}^2$ ), i.e. spring stiffness per unit

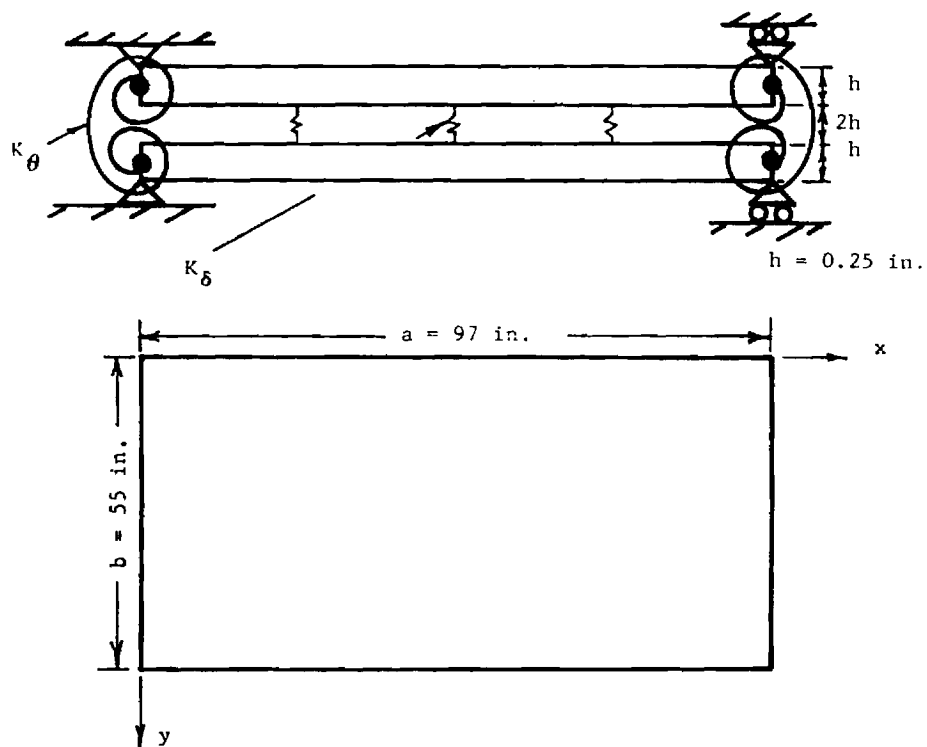


Figure 3.4-1 Continuum Model of the Double Plate System

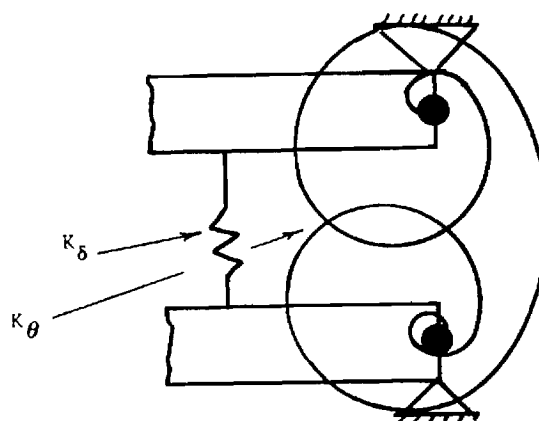


Figure 3.4-2 Details of End Connection



area. To obtain estimates of the natural frequencies, the distributed spring stiffness was taken to be a different constant in each mode of vibration. The stiffness  $K_\delta$  for the entrapped air was determined by assuming adiabatic compression of the volume and by neglecting resonances of the air mass. A linear relationship between volume change and pressure change was assumed to evaluate  $K_\delta$ . The linearized pressure-volume relationship for an adiabatic process can be written as:

$$\frac{\Delta p}{p_o} = -\frac{\gamma \Delta V}{V_o}, \quad \gamma = 1.4 \quad (3.4-2)$$

where  $\Delta p$  is the pressure change,  $p_o = 14.7$  psi (101.4 kPa),  $\gamma$  is the adiabatic exponent,  $\Delta V$  is the volume change, and  $V_o$  is the initial volume of the air space which is equal to 2661 in<sup>3</sup> ( $4.36 \times 10^{-2}$  m<sup>3</sup>). Constant  $K_{\delta mn}$  was then computed for the  $mn^{\text{th}}$  mode as the pressure change per unit modal amplitude, namely,

$$K_{\delta mn} = \frac{\Delta p_{mn}}{a_{mn}} = \frac{\gamma p_o}{V_o} \frac{\Delta V_{mn}}{a_{mn}} \quad (3.4-3)$$

where  $a_{mn}$  is the amplitude in the  $mn^{\text{th}}$  mode. For the (1,1) mode, for example,  $K_{\delta 1,1}$  is equal to 19.88 lb/in/in<sup>2</sup> (4328 N/m/m<sup>2</sup>). This value applies to a rotationally restrained system with  $K_\theta$  equal to 256.2 lb-in/radian/in (1140 N·m/radian/m) which corresponds to  $\xi_B = 2$  in equation 3.4-1.

With  $K_\theta$  and  $K_\delta$  defined, the governing differential equations for the double plate model of Figure 3.4-1 can be written as follows, assuming the two plates are identical:

$$\begin{aligned}
D\nabla^4 w_1 + \rho h \ddot{w}_1 &= -K_\delta (w_1 - w_2) \\
D\nabla^4 w_2 + \rho h \ddot{w}_2 &= -K_\delta (w_2 - w_1)
\end{aligned}
\tag{3.4-4}$$

where  $w_1$  and  $w_2$  are the displacement functions of the upper and lower plate, respectively.

After separation of variables

$$D\nabla^4 W_1 - \omega^2 \rho h W_1 = -K_\delta (W_1 - W_2) \tag{3.4-5a}$$

$$D\nabla^4 W_2 - \omega^2 \rho h W_2 = -K_\delta (W_2 - W_1) \tag{3.4-5b}$$

where

$$W_i = W_i(x, y), \quad i = 1, 2$$

The appropriate boundary conditions are:

$$\begin{aligned}
y = 0 \quad W_1 = 0 \quad D \frac{\partial^2 W_1}{\partial y^2} &= K_\theta \frac{\partial (W_1 - W_2)}{\partial y} \\
W_2 = 0 \quad D \frac{\partial^2 W_2}{\partial y^2} &= K_\theta \frac{\partial (W_2 - W_1)}{\partial y} \\
y = b \quad W_1 = 0 \quad -D \frac{\partial^2 W_1}{\partial y^2} &= K_\theta \frac{\partial (W_1 - W_2)}{\partial y} \\
W_2 = 0 \quad -D \frac{\partial^2 W_2}{\partial y^2} &= K_\theta \frac{\partial (W_2 - W_1)}{\partial y} \\
x = 0 \quad W_1 = 0 \quad D \frac{\partial^2 W_1}{\partial x^2} &= K_\theta \frac{\partial (W_1 - W_2)}{\partial x} \\
W_2 = 0 \quad D \frac{\partial^2 W_2}{\partial x^2} &= K_\theta \frac{\partial (W_2 - W_1)}{\partial x} \\
x = a \quad W_1 = 0 \quad -D \frac{\partial^2 W_1}{\partial x^2} &= K_\theta \frac{\partial (W_1 - W_2)}{\partial x} \\
W_2 = 0 \quad -D \frac{\partial^2 W_2}{\partial x^2} &= K_\theta \frac{\partial (W_2 - W_1)}{\partial x}
\end{aligned}
\tag{3.4-6}$$

The solution of equations 3.4-5 satisfying equations 3.4-6 can be simplified if in-phase and out-of-phase vibration modes are considered separately.

For in-phase motions, no volume change occurs and equations 3.4-5 reduce to

$$D\nabla^4 W_S - \omega^2 \rho h W_S = 0 \quad (3.4-7)$$

where

$$W_S(x, y) = W_1(x, y) + W_2(x, y) \quad (3.4-8)$$

and boundary conditions in equation 3.4-6 become

$$\begin{aligned} y = 0, \quad b \quad W_S = 0 \quad \frac{\partial^2 W_S}{\partial y^2} = 0 \\ x = 0, \quad a \quad W_S = 0 \quad \frac{\partial^2 W_S}{\partial x^2} = 0 \end{aligned} \quad (3.4-9)$$

The eigenvalues and eigenvectors of equation 3.4-7 are

$$\omega_{mn} = \pi^2 \sqrt{\frac{D}{\rho h}} \left\{ \left( \frac{m}{a} \right)^2 + \left( \frac{n}{b} \right)^2 \right\} \quad (3.4-10a)$$

$$W_{S_{mn}} = \frac{2}{\sqrt{\rho a b}} \sin \frac{m\pi x}{a} \sin \frac{n\pi y}{b} \quad (3.4-10b)$$

where

$$\rho \int \int W_{S_{mn}}^2 dx dy = 1 \quad (3.4-10c)$$

For out-of-phase motions, the equations of motion 3.4-5 reduce to

$$D\nabla^4 W_O - (\omega^2 \rho h - 2K_\delta) W_O = 0 \quad (3.4-11)$$

where

$$W_O = W_1(x, y) - W_2(x, y) \quad (3.4-12)$$

If

$$\Omega^2 = \omega^2 \rho h - 2K_\delta \quad (3.4-13)$$

then equation 3.4-11 becomes

$$D\nabla^4 W_O - \Omega^2 W_O = 0 \quad (3.4-14)$$

The boundary conditions governing out-of-phase vibrations are

$$\begin{aligned}
y = 0, \quad b \quad W_o = 0 \quad \frac{\partial^2 W_o}{\partial y^2} &= \pm \frac{\xi_B}{b} \frac{\partial W_o}{\partial y} \\
x = 0, \quad a \quad W_o = 0 \quad \frac{\partial^2 W_o}{\partial x^2} &= \pm \frac{\xi_A}{a} \frac{\partial W_o}{\partial x}
\end{aligned} \tag{3.4-15}$$

where

$$\xi_B = \frac{2K_\theta b}{D}, \quad \xi_A = \xi_B a/b \tag{3.4-1} \text{ repeated}$$

Using assumed eigenfunctions of the form

$$W_o = \sum_m^M \sum_n^N a_{mn} X_m(x) Y_n(y) \tag{3.4-16}$$

the resulting eigenvalue for the  $mn^{\text{th}}$  mode is shown to be

$$\Omega_{mn}^2 = \frac{1}{b^2} \sqrt{\frac{D}{\rho h}} \left[ \left( \frac{b}{a} \right)^4 \alpha_m^4 + \alpha_n^4 + 2 \left( \frac{b}{a} \right)^2 \phi_{mm} \phi_{nn} \right]^{1/2} \tag{3.4-17}$$

where

$$\phi_{mn} = \frac{\alpha_m^2 \left[ \alpha_m (B_m^2 + 1) + 2A_m (B_m - 1) \right]}{\alpha_m \left[ 2A_m^2 - B_m^2 + 1 \right] + 2A_m (B_m + 1)} \tag{3.4-18a}$$

$$\phi_{nn} = \frac{\alpha_n^2 \left[ \alpha_n (B_n^2 + 1) + 2A_n (B_n - 1) \right]}{\alpha_n \left[ 2A_n^2 - B_n^2 + 1 \right] + 2A_n (B_n + 1)} \tag{3.4-18b}$$

Further details are provided in reference 40.

Calculated values of the natural frequencies of in-phase motion up to 55 Hz, and corresponding values for out-of-phase motion are shown in Table 3.4-1. Equations 3.4-10a and 3.4-17 were evaluated to obtain in-phase and out-of-phase values respectively. The natural frequencies for out-of-phase motions are evaluated at extreme and intermediate values of  $\xi_B$  where  $\xi_B = 0$  implies no plate interconnection at the boundaries and  $\xi_B = \infty$  indicates that the plates are rigidly connected.

Table 3.4-1. Natural Frequencies of the Double Plate System,  
Continuum Model

Mode Shape		Natural Frequency, $f_n$ in Hz			
m(x-ax) (1)	n (2)	In Phase S. S. BC's (3)	Out of Phase, $K_\theta$ & $K_\delta$ included		
			$\xi_B = 1 \times 10^{-10}$ (4)	$\xi_B = 2.0$ (5)	$\xi_B = 1 \times 10^{10}$ (6)
1	1	10.46	120.02	129.35	160.14
2	1	18.09	18.09	20.04	28.44
3	1	30.81	75.60	81.98	112.46
1	2	34.21	34.21	36.86	52.10
2	2	41.84	41.84	44.32	59.80
4	1	48.61	48.61	50.41	62.03
3	2	54.56	54.56	56.85	72.80

### 3.4(b) Discrete Element Model

A parallel-plate finite element model was developed for comparison with the continuum model described above. The model shown in Figure 3.4-3 includes the entire frame and panel but only the upper plate is represented due to symmetry. The boundary frame is simply-supported at all frame joints and the panel is permitted to vibrate freely within the framework. Discrete springs with spring constant  $k_N$  equal to twice  $K_{\delta 1,1}$  (due to symmetry) times tributary panel area are attached to interior panel nodes to account for air gap stiffness in the first breathing mode. The value of  $K_{\delta 1,1}$  of 19.88 lb/in/in<sup>2</sup> (4328 N/m/m<sup>2</sup>) reported earlier is used to compute  $k_N$ . Normal springs on the panel boundary, representing the normal stiffness of the spacer bar between the glass plates, have a spring constant  $k_{NB}$  which varies between soft ( $8.85 \times 10^3$  lb/in =  $1.55 \times 10^3$  kN/m) and hard ( $5 \times 10^7$  lb/in =  $8.76 \times 10^6$  kN/m) values so that a range of breathing mode frequencies could be determined. The rotational spring on the panel boundary representing the rotational stiffness of the spacer bar has spring constant  $k_R$  equal to twice  $K_\theta$  (due to symmetry) times tributary boundary length. The value of  $K_\theta$  was set at 256.2 lb-in/radian/in (1140 N·m/radian/m) as in the continuum model.

Discrete model results were generated for the first breathing mode only but for soft and hard springs  $k_{NB}$  and for soft, intermediate, and hard springs  $K_\theta$  (expressed in terms of  $\xi_B$  in equation 3.4-1). Frequencies of vibration are tabulated in Table 3.4-2. The flexible spring boundary used in the discrete model results in first mode frequencies which are up to 10% lower than continuum model values, but hard spring results are within 5% of continuum model frequencies.

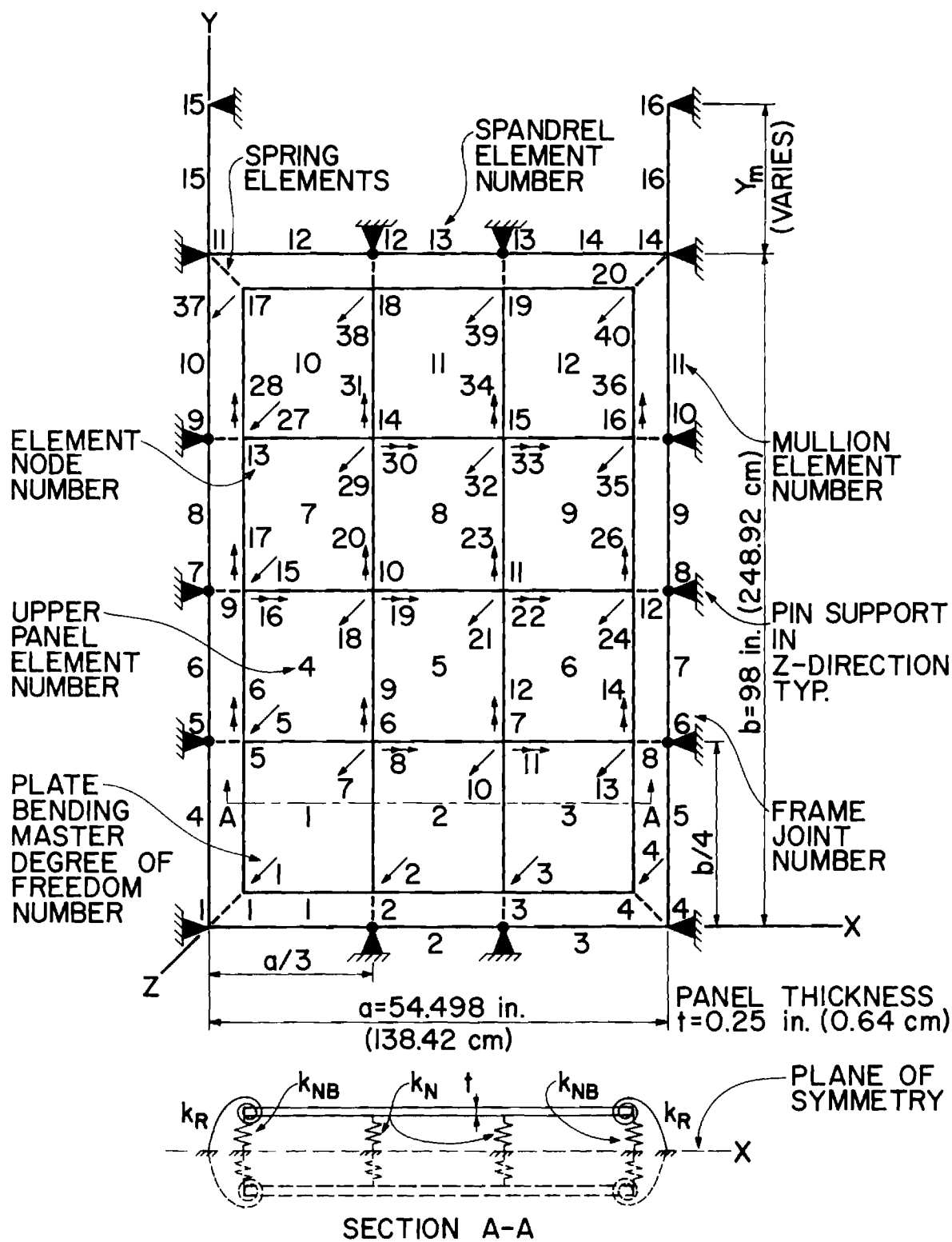


Figure 3.4-3 Finite Element Model of Double Plate System

Table 3.4-2 Natural Frequencies of Double Plate System,  
First Breathing Mode

Model  (1)	Natural Frequency in Hertz		
	$\xi_B = 1 \times 10^{-10}$ (2)	$\xi_B = 2$ (3)	$\xi_B = 1 \times 10^{10}$ (4)
Continuum <sup>a</sup>	120.02	129.35	160.14
Discrete			
--soft spring <sup>b</sup>	116.93	119.91	144.68
--hard spring <sup>c</sup>	118.44	122.28	156.89

<sup>a</sup>See Table 3.4-1.

<sup>b</sup> $k_{NB} = 8.85 \times 10^3 \text{ lb/in } (1.55 \times 10^3 \text{ kN/m})$

<sup>c</sup> $k_{NB} = 5.0 \times 10^7 \text{ lb/in } (8.76 \times 10^6 \text{ kN/m})$



### 3.5 Overall Structure Cladding Model

The laboratory, field, and computer studies described in this report have sought to provide information on dynamic loadings and response mechanisms which affect the structural integrity of individual window-frame units such as those found in the prototype structure. Attention has been focused on a small part of the structure's curtain wall and factors affecting interaction of the exterior cladding with the structural frame have not been considered in these initial studies. However, a number of recent failures particularly of lightweight cladding elements demonstrate that the building curtain wall forms a first line of defense for the structure. In fact the overall cladding system may actually provide a considerable amount of resistance to low level excitations and help to control interstory drift and building motion affecting occupant comfort.

The stiffness of expensive cladding elements can, and should, be used to advantage by the structural designer. In the past, the exterior curtain wall has been treated as nonstructural and isolated from the motions of the primary structure. However, analytical models of the total structure including both the basic frame and exterior curtain wall can be developed using a multistory building model, such as the tier building substructure model (93,94,95), and an appropriate cladding model of the type described in this report. For example, a finite element representation of a portion of the curtain wall can be assembled and inserted into the framing substructure between floor levels A and B as shown in Figure 3.5-1. In this way, all elements contributing stiffness to the overall structure can be accounted for in a rational manner and cladding-structure interaction forces determined for use in design of the curtain wall. Follow-on studies of this kind are currently underway.

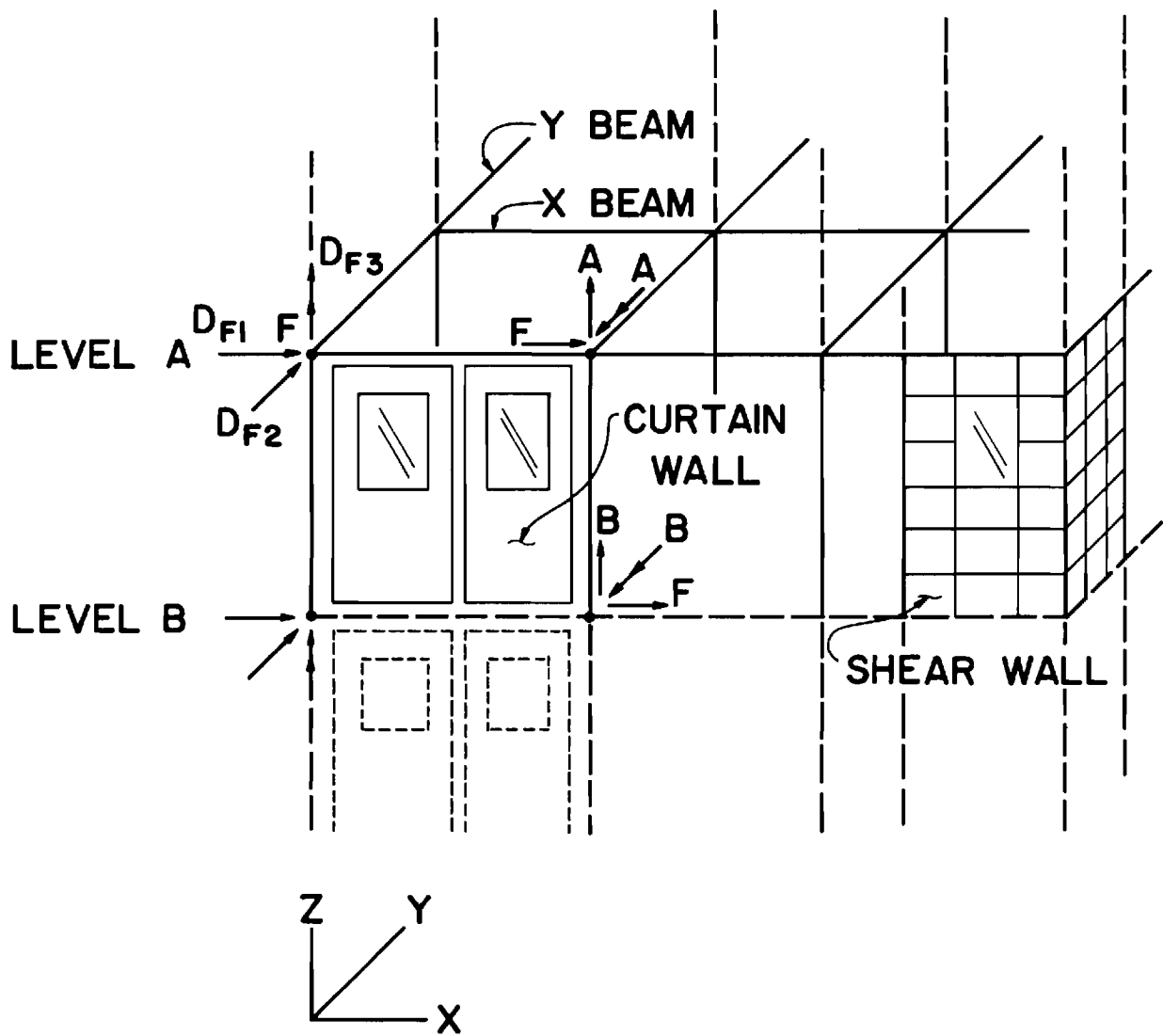


Figure 3.5-1 Tier Building Substructure with Curtain Wall Model

### 3.6 Summary

The initial analytical studies of the static and dynamic response of glass curtain wall systems described in this chapter have focused attention on the properties and behavior of a small segment of the exterior curtain wall encompassing the frame, gaskets, and a single glass panel. A frame-finite element model and computer program were developed to represent the cladding system and model properties adjusted until a good correlation between laboratory and computer response was established. The analytical model is general in scope and permits a wide range of cladding configurations and loading types to be considered. Panel response to static and dynamic pressure loadings and racking distortions introduced by differential motion of adjacent floors were some of the loadings considered.

Sensitivity studies were conducted to determine the influence of model properties on the static and dynamic response of the panel-frame assembly. Response was shown to be sensitive to gasket rotational stiffness and frame flexural stiffness for selected panel aspect ratios. Results of the analytical studies are presented in a series of nondimensional plots of selected properties versus normalized panel displacement or natural frequency of vibration.

The analytical cladding model is compatible with existing three-dimensional multi-story building models and is currently being used to study cladding-structure interaction in highrise buildings.

#### 4.0 LABORATORY FACILITY

##### 4.1 Background

A comprehensive review of the cladding loads and response problem was presented in Chapter 1, and a field study to both qualitatively and quantitatively explore the problem in a uniquely shaped highrise building was then described in Chapter 2. The data collected from the field study consisted largely of differential pressure measurements taken across the cladding at numerous locations and for various wind conditions. These data were used to construct a probabilistic description of the time and frequency domain aspects of the localized pressure, and by comparison to measured window response data, to determine the causality between dynamic wind pressure loading and response over a range of frequencies. While the measurements provided invaluable data under service conditions, a number of problems were encountered in both refining and generalizing this information. First, it was not possible to control the manner of loading or the test conditions; for example, a 15-20% variation in dynamic properties as a result of mounting, support framing, etc., was observed over 12 windows. In addition even though sophisticated remote sensing transducers and portable instrumentation were used, it was not practical to make extensive measurements of the outer pane response. For obvious reasons, as well, it was not possible to geometrically or structurally alter the windows and associated cladding.

The full scale, laboratory grade dynamic loading facility described in this chapter provides the means for resolving many of the problems encountered in the field study, and in addition, for the first time

offers the general capability for simulating a variety of dynamic loading conditions under controlled conditions. The facility was designed and built under the present research project and is currently in operation in the School of Civil Engineering (96).

Present design procedures for windows consider an essentially static pressure load and rely upon full scale laboratory or on-site tests to verify performance. In this case, static retention of the lights and resistance to water penetration are the prime considerations. As a result, most of the tests have involved extremely simple arrangements for providing a static pressure while at the same time spraying intense streams of water at the test assembly. Dynamic response mechanisms have not generally been considered in the experimental work, and the reason can be traced to two problems: (a) the dearth of information as to the precise nature of the dynamic loads themselves, and (b) the lack of adequate test facilities for producing dynamic loads. Chapter 2 addressed the first point; this chapter deals with the latter.

The concept of a full scale wind load simulation facility is not new. The fixture described by Frownfelter (97) was in operation in 1959 and is capable of producing pressure loads of up to 40 psf.(1.9 kPa) It can accomodate a 16 ft by 16 ft (4.9m by 4.9m) specimen, thus making it possible to test several sections of a curtain wall and anchor system as installed on a real structure. The specimen wall forms one side of a sealed chamber which can be partially evacuated using a high volume blower. Suction was controlled with a valve on the blower exhaust and measured with a direct reading pressure gage. Deflection of the system under load was measured with dial

indicators and stresses were measured with resistance wire strain gages. Although the importance of dynamic loading was recognized and the test apparatus could produce pulsating pressures, no details were given on the method used to drive the system and no dynamic load data or dynamic performance data were included in the publication.

A more elaborate Japanese test facility was completed in 1965 and can produce time varying pressure loads representative of a typhoon or other severe weather (98). The typical specimen size is approximately 10 ft by 11.5 ft high (3m by 3.5m) As with the apparatus described by Frownfelter, the specimen forms one side of a pressurized chamber, but two blowers are used, one to create positive pressure and the other for suction. The two separate supply pipes used to connect the blowers to the pressure chamber each contain a butterfly damper valve. The butterfly dampers are linked together and operate 90° out of phase so that one is fully opened while the other is closed. The apparatus was designed to produce peak pressures of approximately  $\pm 100$  psf ( $\pm 4.8$  kPa) with a flat frequency response up to 0.3 Hz. An analog computer was used to control the butterfly servo mechanism with feedback signals obtained from a potentiometer on the damper linkage and a secondary feedback transducer monitoring chamber pressure. An optical curve tracer converted the desired pressure excursions from analog paper charts to voltage which was then used as primary input to the computer. Except for some leakage problems that were discussed, the equipment appeared to perform as designed.

The British Standards Institute at Hemel Hempstead has constructed a wind load test facility which was reported on in 1974 (99). The apparatus is similar in concept to the equipment previously discussed.

It consists of a double chamber arrangement driven by an air blower whose output is boosted by a compressed air reservoir. A typical test cycle consists of raising the internal pressure on one side of the chamber to a maximum in approximately 7 seconds holding at this level for 3 seconds, and finally bleeding down to ambient over 6 seconds, to complete a 16 sec. cycle. The cycle can be immediately repeated, and the proposed test sequence calls for 5 such cycles at each test pressure, to be followed by similar treatment in the opposite direction. The general emphasis in this facility is more towards standards testing than experimental research. The Paisley Institute of Technology boasts a similar facility and photographs of the Paisley equipment appear in (104).

All of the above facilities were designed primarily to check the integrity of windows or curtain wall structures when subjected to severe steady state or low frequency ( $< 1$  Hz) wind loads. In addition, some have the capability of spraying the specimen with water to check for resistance to rain leaks. The facility described here has been designed to extend the simulation capabilities to a higher frequency range more representative of sharp, short duration transient loads such as might be produced by severe localized flow separation or perhaps by sonic booms. No provision has been made for water injection.

#### 4.2 Design of the Facility

A peak static pressure capability of 75 psf (3.6kPa) was chosen as being a conservative maximum based upon the existing body of experimental data on wind pressure loads. The maximum dynamic component is not so readily specified, since relatively few studies have been

directly concerned with the frequency content of localized wind pressures, or have employed pressure transducer systems with flat response above about 5 or 10 Hz. The desire to simulate sharp-edged transient pressures coupled with mechanical compliance and driving power limitations above about 20 Hz combined to fix the dynamic specification at  $\pm 25$  psf ( $\pm 1.2$  kPa) at a 20 Hz. cyclic rate. This figure is consistent with wind tunnel measurements of fluctuating pressures (101), recent full-scale measurements (102), and the results of the present field study presented in Chapter 2. The range is adequate to provide significant driving power for the lower vibration modes of many components of typical cladding systems, particularly windows with areas greater than about 20 sq. ft. ( $2\text{m}^2$ ).

The principle of mounting the cladding specimen to the side of an airtight box which can be pressurized has been amply justified by previous experiments and was chosen as the basis of this design. Further, the existence of a robust hardback structure of suitable dimensions fixed the maximum specimen size to be 8 ft. by 15 ft. high (2.5m by 4.5m). The main frame was constructed of heavy structural steel tubing, and as no water spraying or other bulky equipment is required and response of the apparatus is important, the chamber depth was limited to the dimension of the steel section. This minimizes the chamber volume for best dynamic performance as will be shown later.

Two system concepts were considered for producing the time-varying pressures in the test volume: (a) a flow control design similar to that in (98), and (b) a positive displacement design in which



moving "pistons" are used to produce volume changes and hence pressure fluctuations. The positive displacement approach was ultimately adopted, primarily for its better frequency response above 5 Hz, and because of the ready availability of high pressure hydraulic servo components with low moving mass and high force capability.

In order to estimate the volume change required to produce the design pressure levels and frequency response, several assumptions were made:

- (1) Chamber volume would be approximately 80 cu. ft ( $2.3\text{m}^3$ ).
- (2) Test panel compliance would provide a maximum volume change equivalent to the volume under the deflected shape of a 3/4 in. (19.1mm) glass (or aluminum) plate covering the entire fixture.
- (3) The effect of chamber leaks on the dynamic response is equivalent to increasing the fixture compliance.
- (4) The chamber pressure varies isothermally.

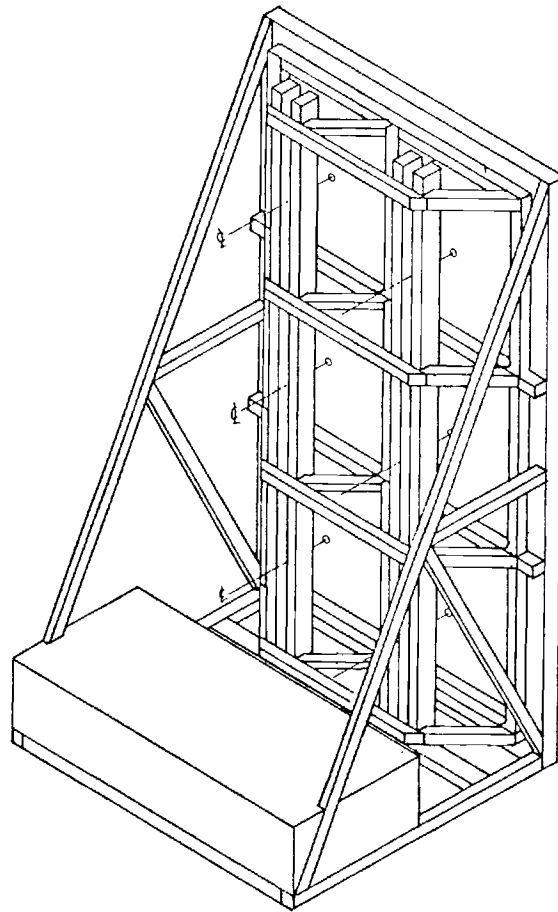
The maximum test panel compliance when considered over the design pressure change will produce a 3 cu. ft. (85 l.) change in the chamber volume. An additional 3 cu. ft. (85 l.) displacement is used to account for the effects of leaks; the figure represents roughly 40% of the total displaced volume and is reasonable based upon the results of leak tests carried out in (98). Finally, assuming an otherwise rigid test chamber, the isothermal volume charge that is required to produce the design limit psf pressure change can be readily calculated at approximately 2 cu. ft. (57 l.). Thus, the positive displacement system must be capable of producing a gross dynamic volume change of 8 cu. ft. ( $0.23\text{m}^2$ ) at a 20 Hz rate. This amounts to a 10% volume

change in 0.05 seconds.

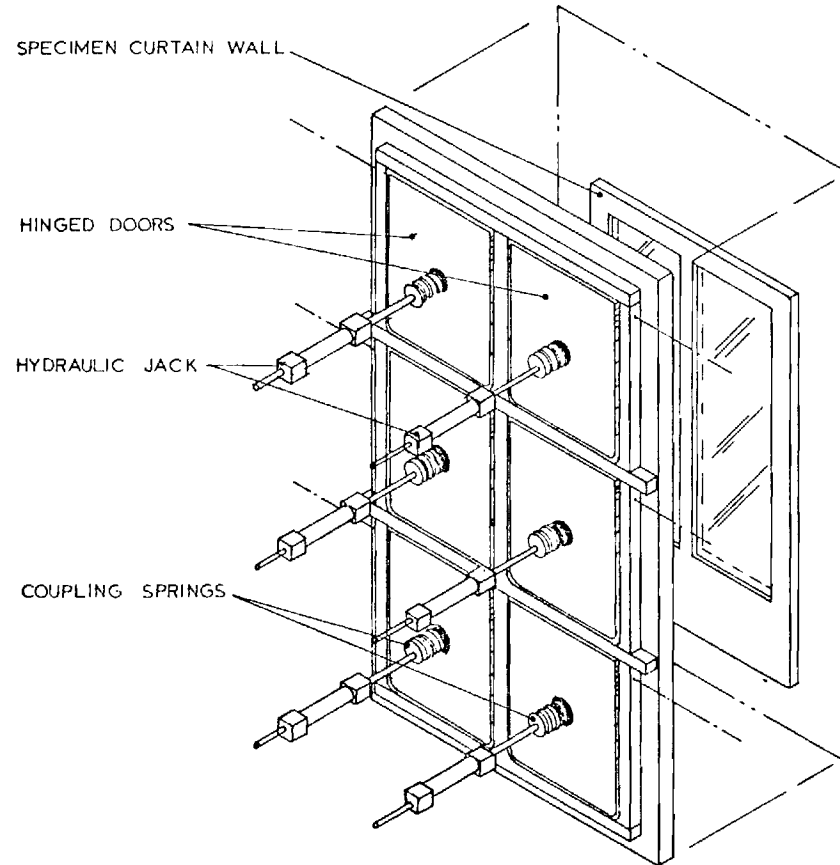
The geometry of the final positive displacement type design is shown in Fig. 4.2-1. The design consists of an upright welded steel test fixture approximately 15 ft. by 8 ft. (2.5 by 4.5 m) which supports the hydraulically driven positive displacement actuators mounted behind the test panel. The test panel(s) are attached over the front of the fixture and enclose the 80 cu. ft. ( $2.3\text{m}^3$ ) chamber volume. The displacement system consists of 6 large "doors" hinged at their top edge and pneumatically sealed on all edges to the test chamber by means of a flexible seal. The doors are each 15 sq. ft. ( $1.4\text{m}^2$ ) in area and cover the entire rear surface of the chamber. Each door is driven at its center by a servo-controlled hydraulic actuator which receives electronic command signals from a feedback control system. The actuators are supported by a steel framework structure which forms part of the rear of the test fixture and which is capable of handling both the static as well as the dynamic actuator reactions. Static loads are obtained by either pressurizing or evacuating the chamber using a blower system, and the door displacements produce the prescribed dynamic component of loading. The fixture is fitted with a large capacity pressure relief valve (actually, a simple weighted poppet valve 6 in. (15cm) in diameter) to prevent inadvertent static or dynamic over-pressures.

Major design considerations which ultimately led to this system configuration can be summarized as:

- (1) Components necessary for a one moving piece displacement system were too large to handle and prohibitively expensive.
- (2) Off-peak conditions are more readily handled by six smaller



(a) Structure



(b) Geometry

Figure 4.2-1 Test Fixture Structure and Geometry

components; it is not necessary to drive all six doors all the time.

- (3) Future simulation of pressures with spatial variation is possible with a split system.
- (4) Hydraulic servo valves of suitable size were immediately available.

The completed fixture with two of the six doors installed and operational is shown in Fig. 4.2-2 and 4.2-3. A large double glazed window (approximately 40 sq. ft. ( $3.7\text{m}^2$ )) is installed in the test area. The servo-control electronics and minicomputer data acquisition are also shown. The rear view shows a more detailed picture of the high pressure hydraulic components required to operate the doors.

#### 4.3 Performance Evaluation

The number of active doors, their mass and area, the hydraulic actuator size, and the hydraulic flow rates are all related and a particular combination is required to meet the specified performance. The principle constraints were the available hydraulic driving power and the maximum hydraulic flow rate to each door limited by the available servovalves). Increasing the number of doors decreases the flow rate to each, while increasing the door area reduces the required actuator stroke and hence the flow rate as well. The effective piston area of the hydraulic jack fixes the maximum force output and is thus related to the mass of the moving doors and the design frequency and pressure objectives. The door is by far the most critical component in the system design because of the mass and strength considerations. After several studies, the concept of ultra-lightweight doors was deemed technically feasible but was rejected because of the large

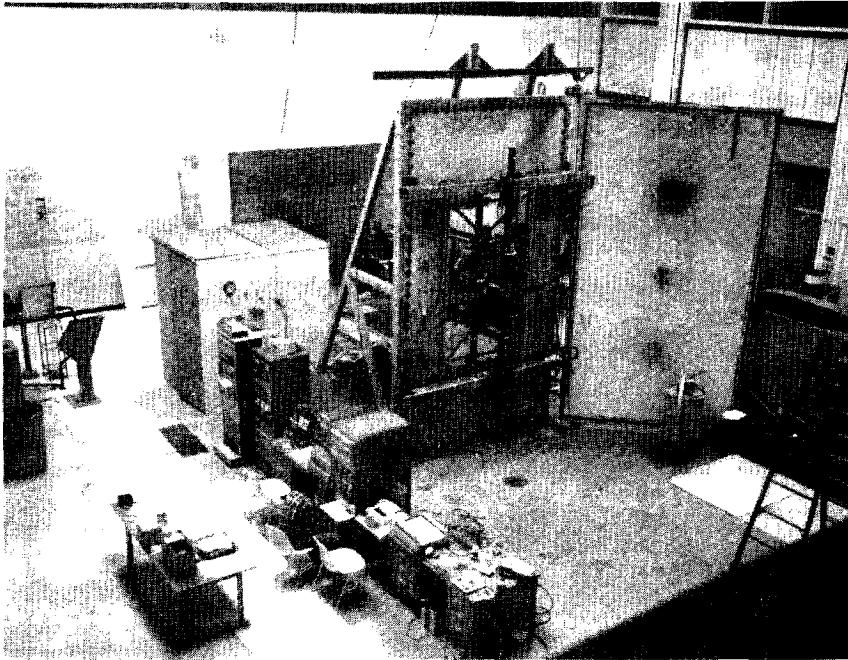


Figure 4.2-2 Test Fixture and Instrumentation.

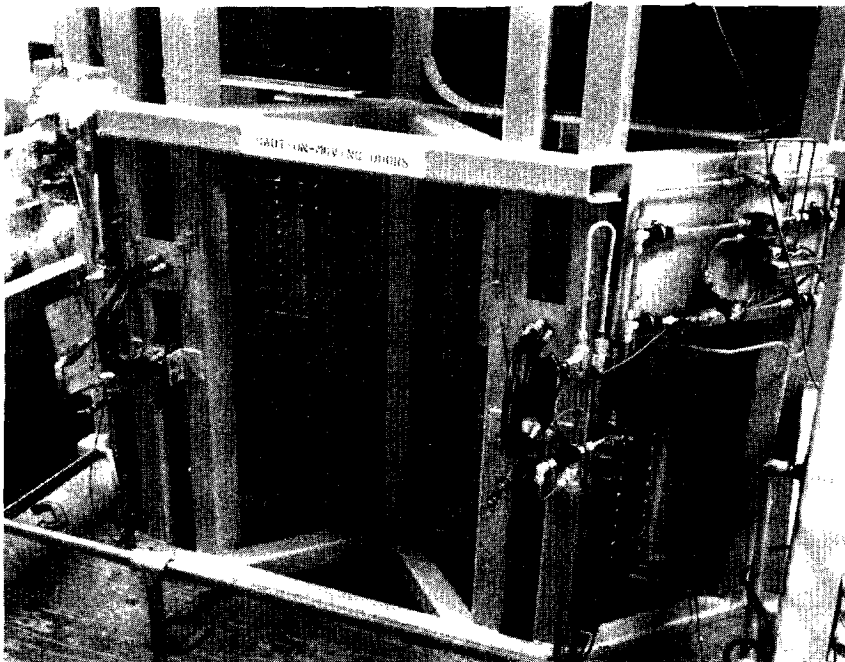


Figure 4.2-3 Moving Doors and Hydraulic Controls.

expense and the time required to build them. The present doors were designed primarily to be rugged and easily assembled, and secondarily, to be as light as possible.

The doors are driven through adjustable coupling springs by small, 1 1/8 in. (28.6 mm) bore hydraulic actuators mounted to the support structure. The role of the coupling springs will be described later. The hydraulic flow to the actuator is modulated by an electronically controlled high-response servovalve manifolded directly to the actuator.

A schematic of the driving system and the feedback control required to operate it is shown in Figure 4.3-1. The springs are lumped parameter models for the specimen compliance ( $k_w$ ), the chamber compressibility including leaks ( $k_a$ ), and the coupling spring ( $k_s$ ), while the door is represented by a lumped mass. It is necessary to use a feedback type control system as shown in order to accurately produce the time-varying pressure load. Mean pressures are supplied statically but the dynamic component is produced entirely by door movement. Primary control of the door motion in response to a forcing function,  $p(t)$ , is therefore required. However, direct control of the actuator in response to a prescribed pressure forcing function is not desirable because this places a resonant mechanical system ( $m$ ,  $k_s$ ,  $k_a$ ,  $k_w$ ,  $m_w$ ) in the loop and requires extensive compensation for stable operation. Rather, the actuator is operated under direct stroke control with stability compensation provided by a  $\Delta P$  transducer and rate control in  $A_2$ . The setpoint for this inner loop is provided by a second loop consisting of a door motion transducer and controller,  $A_1$ . The command signal,  $p(t)$ , can be furnished by a signal generator, a noise generator, or from a tape recording of actual wind pressures.

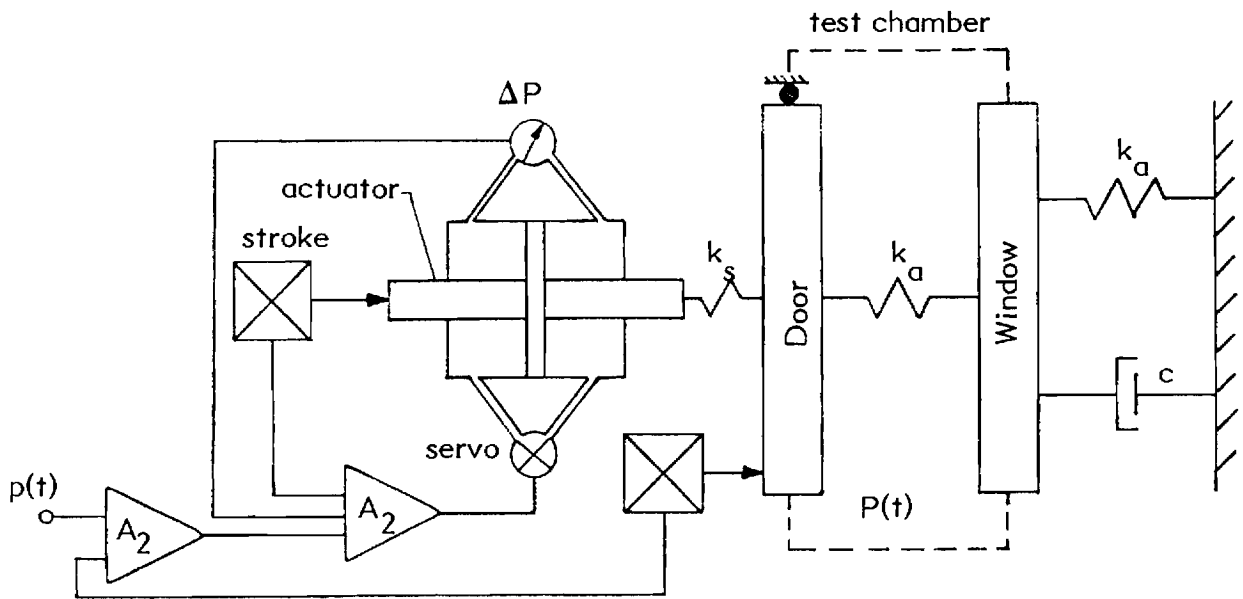


Figure 4.3-1 Feedback Control System

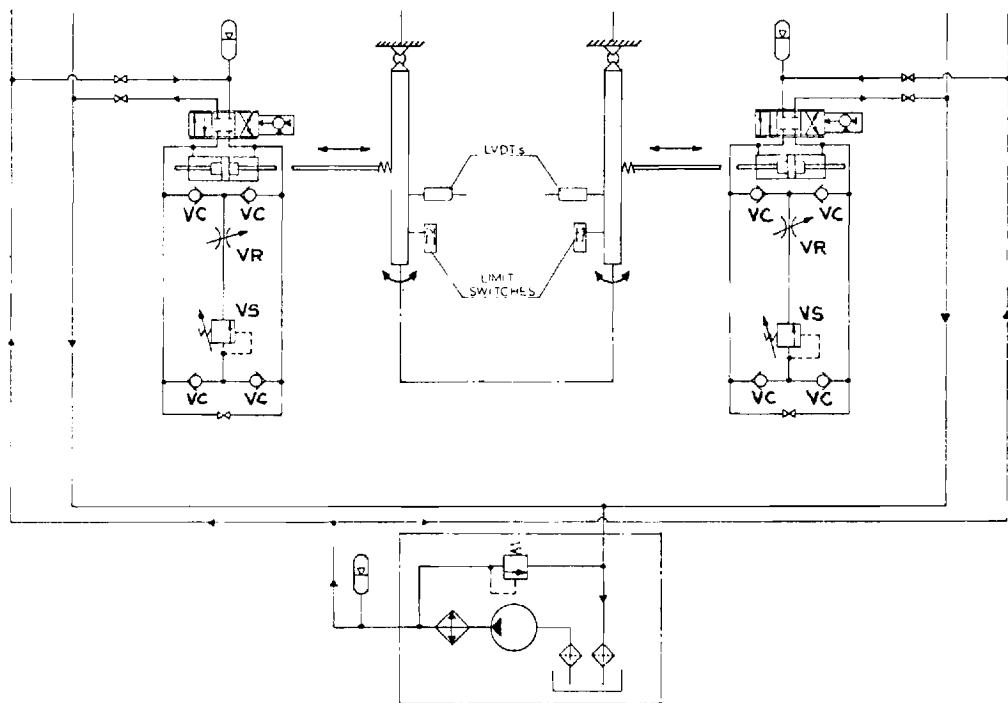


Figure 4.3-2 Hydraulic Schematic Showing Actuator Overpressure Circuits (2 doors)

The lumped parameter values and component sizes are fixed by the dynamic response requirements. For example, to a first approximation, the compressibility effects in the chamber as felt by a single door can be represented by a linear spring with stiffness:

$$k_a = k_o A_d^2 N$$

where  $k_o = \Delta P / \Delta V$  is the ratio of pressure change to the isothermal volume change required to produce it,  $A_d$  is the door area, and  $N$  is the number of doors operating together. The effect of specimen compliance and leakage can be considered similarly as a series stiffness:

$$k_w = k_e A_d^2 N$$

where  $k_e = \Delta P / \Delta V_e$  is the ratio of pressure change to chamber volume change due to test panel flexure and fixture leakage. The overall fixture stiffness as seen by each door is, then, the series combination of  $k_a$  and  $k_w$ , and given the door masses, the system natural frequencies can be determined from this. An interesting feature is that the stiffness, and hence natural frequencies, are dependent upon the numbers of driven doors, so that system dynamics may be readily adjusted by varying this value. For example, with the present door weight of 64 lbs (29 kg), when all six doors are driven the natural frequency is 14 Hz, but when only two doors are used this drops to 6 Hz.



If the system were driven directly by a prescribed actuator force, large and uncontrollable door motions and pressure amplitudes would result from the resonant behavior. However, when driven with servo-displacement control of the door and actuator stroke, a properly tuned feedback controller will suppress the resonance and provide uniform response. Two limitations exist that reduce the ultimate performance:

- (1) Hydraulic flow restrictions (10gpm (69 l/m)) in the servovalve and actuator effectively limit the maximum piston velocity and hence stroke at a given frequency.
- (2) Actuator force is limited by the maximum hydraulic pressure available (3000 psi (20.7 MPa)) and the total opposing static chamber pressure on each door.

The limitations are all the more severe because they are self-conflicting; that is, piston velocity can be increased by reducing the bore, but this then reduces the maximum force available.

The performance can, however, be improved within these restrictions by two methods:

- (1) By decreasing the door mass or increasing the number of driven doors the natural frequency is raised; this reduces the required force at higher frequencies, especially near natural resonance,

- (2) By adding a coupling spring,  $k_s$ , (Fig. 4.3-1) the door is partially uncoupled from the actuator stroke restriction so that large motion is possible at high frequencies; the system thus operates in a controlled-resonance mode.

For typical windows of less than about 40 sq ft ( $3.7\text{m}^2$ ), the latter method is not required and adequate performance can be obtained simply by increasing the number of driven doors.

When operating under feedback control near resonant conditions, there is always the risk that failure of a control circuit or transducer will result in loss of control and possibly catastrophic instability of the system. A simple pressure relief valve described earlier will automatically limit the maximum test pressures. Hydraulic pressures can, however, exceed allowable levels independently so that additional protection circuits are required. A special hydraulic damping circuit as shown in Fig. 4.3-2 (showing circuits for two doors) is connected across the actuator to provide protection to the equipment in the event of loss of power or control near resonance. The circuit uses an adjustable sequence valve, VS, and four check valves, VC, arranged so that when hydraulic pressure across the actuator piston exceeds a preset value, the sequence valve will open and shunt the actuator parts together through a flow-restricting needle valve, VR. The sequence valve automatically resets at lower pressures and normal operation is restored. In combination with the relief valve, this feature can be used to mechanically limit the maximum allowable test pressures.

In order to verify operation of the facility at design levels,

the system was configured with two driven doors and fixed panels in place of the other four doors. The test area was enclosed with a stiff reinforced steel plate and the chamber volume carefully sealed. A single door, and later both doors together, were driven sinusoidally at fixed amplitudes over the 20 Hz operating range. Actual door motion and the fluctuating chamber pressures were then recorded for several static pressure levels. Control loop compensation was adjusted, first to identify stability limits beyond which large, uncontrollable oscillations occurred, and finally to tune the system for maximum stable dynamic response. The measured single door natural frequencies agreed well with the calculated value of 6.8 Hz for a well-sealed chamber.

Figure 4.3-3 shows the performance envelopes for the two-door system for several different mean pressure levels and fixture compliances. Flow restrictions limit the stroke inversely with frequency, and this can be translated into a maximum cycle pressure limitation which depends, as shown by the two curves, on chamber compliance. The envelop curves fall off at a 6dB/octave rate.

The performance is also limited by the maximum actuator force available to drive the inertial loads (door mass). This restriction depends directly upon the mean chamber pressure since a static hydraulic pressure level must be applied to maintain the door at rest. The maximum available force, and hence chamber pressure, depend inversely on the square of frequency and consequently roll off at a 12 dB/octave rate. At low peak pressures (or stroke) the maximum cyclic rate is alternately limited by inertial loads while at higher pressures the limit is due to flow capacity.

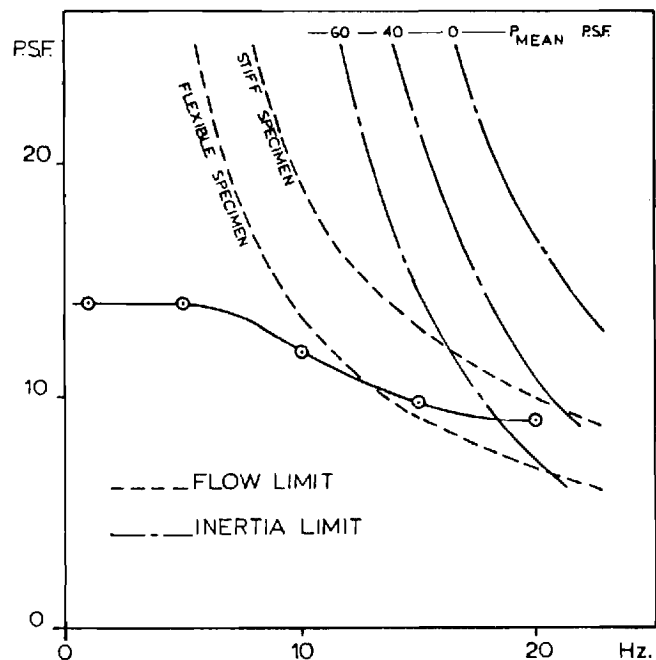


Figure 4.3-3 Performance Envelopes and Maximum System Performance Results for Two Doors

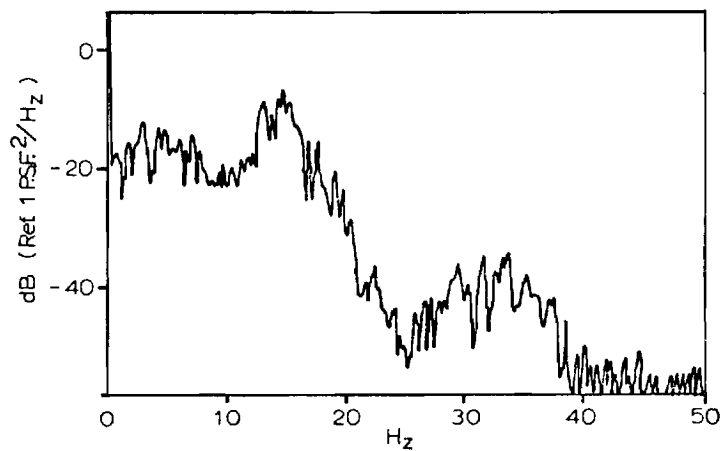


Figure 4.3-4 Typical Simulated Pressure Power Spectral Density

The performance is also limited by the maximum actuator force available to drive the inertial loads (door mass). This restriction depends directly upon the mean chamber pressure since a static hydraulic pressure level must be applied to maintain the door at rest. The maximum available force, and hence chamber pressure, depend inversely on the square of frequency and consequently roll off at a 12 dB/octave rate. At low peak pressures (or stroke) the maximum cyclic rate is alternately limited by inertial loads while at higher pressures the limit is due to flow capacity.

Actual performance values are shown in Figure 4.3-3 for a 14 psf (670 Pa) cyclic load produced with a 15 psf (720 Pa) mean pressure. The curve is essentially flat with frequency until the flow limitation is met. However, it can be seen that this limit is exceeded at 20 Hz. The explanation lies in the use of a finite coupling spring,  $k_s$ , which in this case consisted of a heavy rubber bushing attaching the actuator to the door.

A typical simulated wind pressure power spectrum obtained using an ensemble of 20 pressure records is shown in Fig. 4.3-4. The forcing function in this case was a pseudo-random binary sequence signal generator with a Gaussian amplitude distribution. Low pass filters were used to shape the rolloff to approximate wind pressure characteristics.

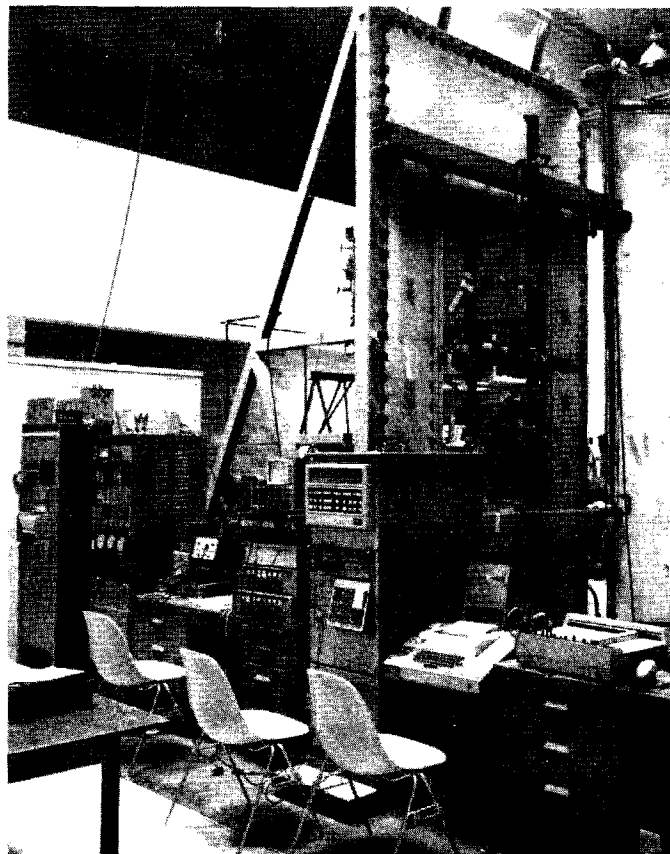
#### 4.4 Static Testing

The full scale test fixture described in Section 4.3 was used for the present work to conduct static and dynamic tests on a section of cladding from the same high rise building used in the field study (Section 2). A single storey section of the cladding, including mullions, muntions, spandrel framing, glazing materials and 4 double

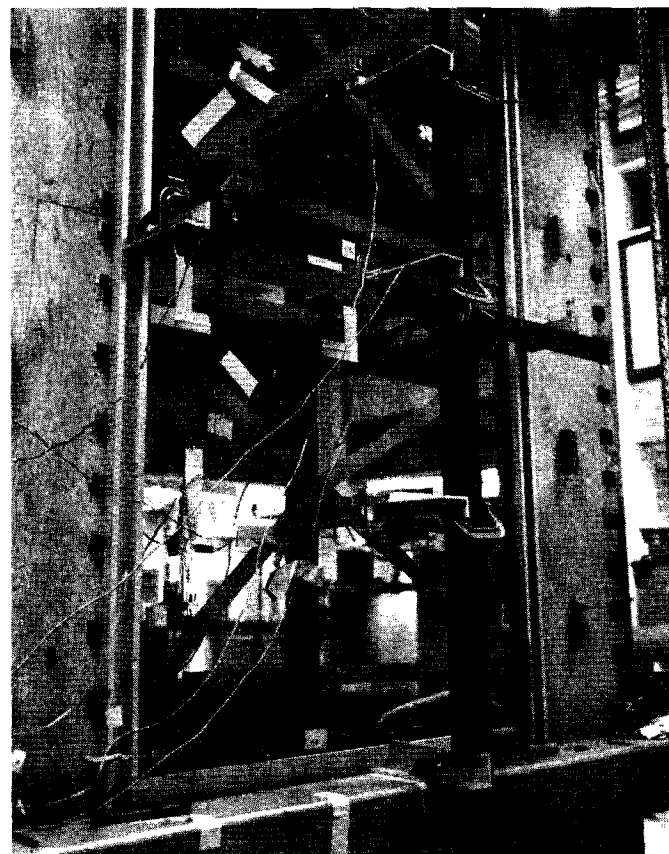
pane vision lights, was obtained from the fabricator and installed in the test fixture as shown in Fig. 4.4 1. In actual service, the 24 ft. (731 cm) mullion sections are pinned to the floor slabs at a 12 ft. (366 cm) spacing with a splice every 2 storeys, but in the present case, only two 12 ft. (366 cm) sections were available so a single storey spliced configuration was constructed. The overall layout along with typical sections are shown in Fig. 4.4-2. The vision light was a 99 x 57 x 1 in. (251.5 x 144.8 x 2.5 cm) insulated silver reflective glass unit composed of two 0.24 in (6.4 mm) thick panes bonded along their edges to a 0.5 in. (12.7 mm) thick square steel spacer. Plywood panels were installed around the window in the same manner to simulate the effect of spandrel sections (single glazing in the field site) and adjacent windows. With the exception of the light itself, the entire test section was sealed internally with a plastic membrane so that internal pressurization could be produced with minimum leakage. The window assembly was oriented with the normal "outside" surfaces facing outward so that building external suction (negative) pressures could be produced. It should be noted, however, that the test fixture is capable of handling negative internal pressures and was limited in these tests to positive pressures only for convenience of sealing. Details of the design of the window assembly and the static test program along with results and a static analysis are given in Ref. 104.

#### 4.4(a) Boundary Properties

A major step in understanding and modelling the behavior of a window must begin with a thorough and precise evaluation of the

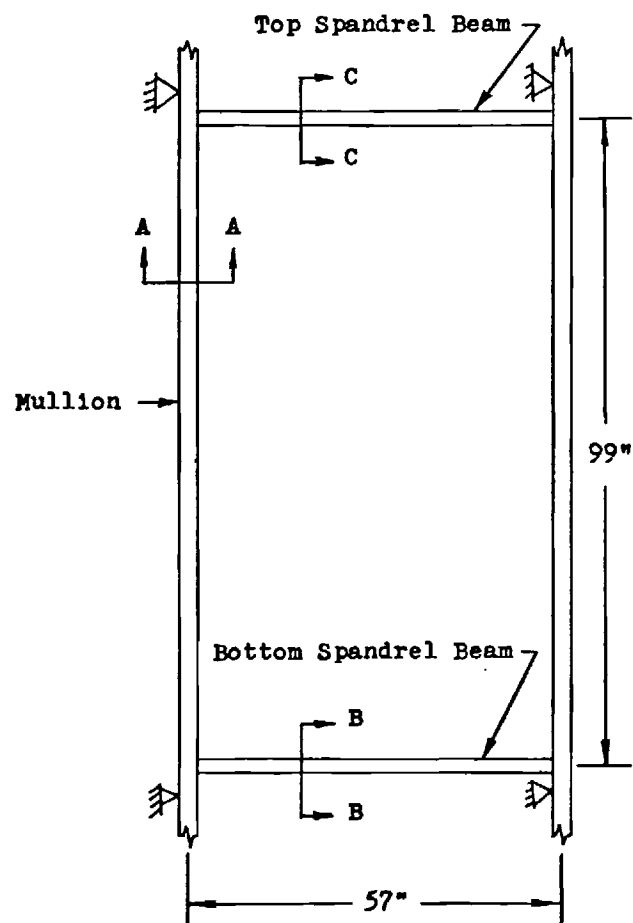


General View

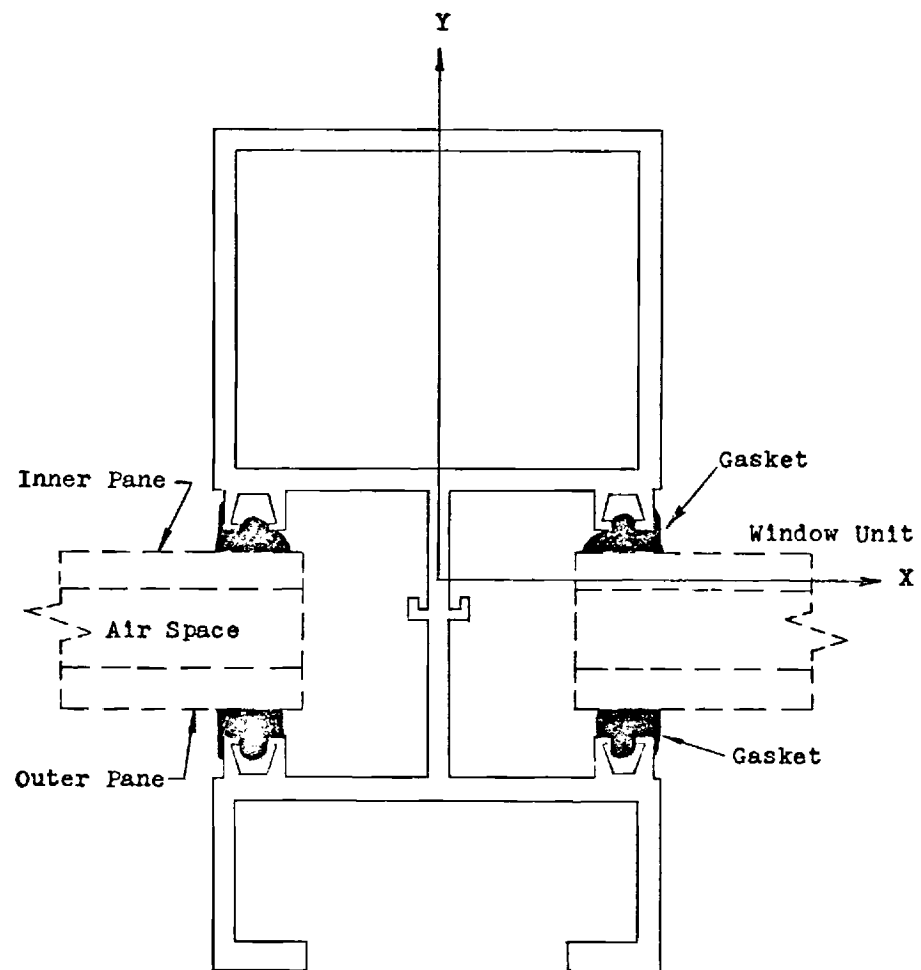


Instrumentation Detail

Figure 4.4-1 Test Window Installed in Fixture.



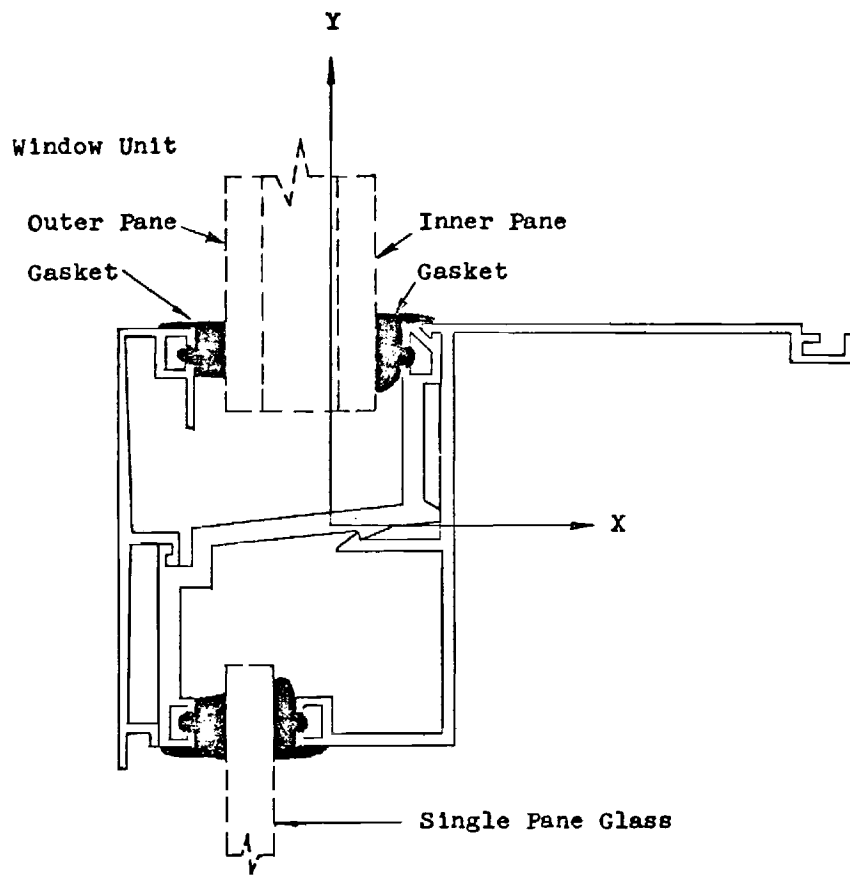
(a) Window Geometry



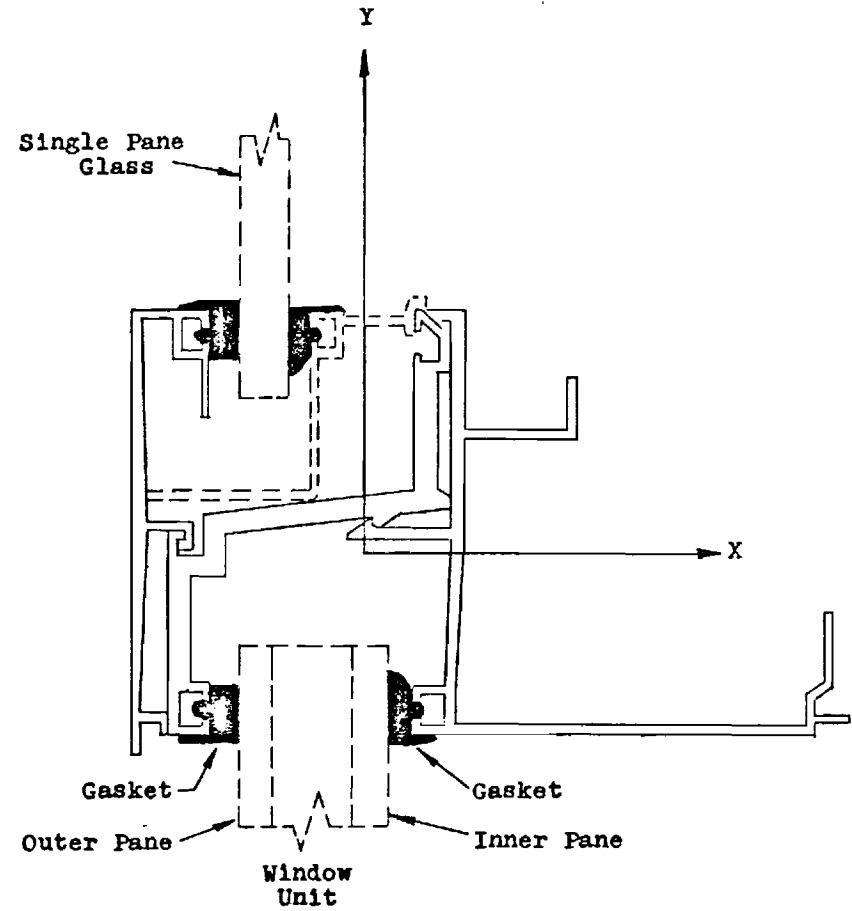
(b) Section A-A

Figure 4.4-2 Test Window Geometry with Framing Section Details





(c) Section B-B

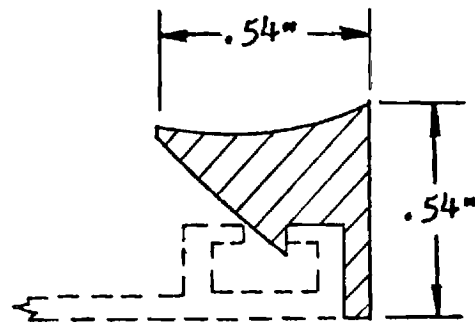


(d) Section D-D

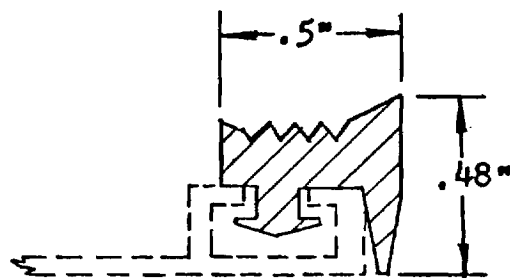
Figure 4.4-2 (continued)

supporting structure, its configuration, dimensions, and material properties. The window/cladding system employed in the present tests is representative of a large class of glazing systems using metal mullion-spandrel sections with neoprene sealing gaskets. As a result, while the reported properties and performance are for a specific configuration, the general behavior is representative of a larger class.

The glazing support structure is a built-up framework consisting of several interlocking aluminum extrusions as shown in Fig. 4.4-2. The glazing gaskets which actually support the glass within the framework are made from special neoprene rubber arranged as shown in Fig. 4.4-3. The "outer" gasket is made from a closed-cell neoprene sponge of medium density while the "inner" gasket is a 75-duro solid neoprene. The outer gasket is located in the frame first, the window is next installed and positioned on two or more glazing blocks, and finally, the inner gasket is wedged into position using a special tool and some force. This latter step can result in differences in support conditions from window to window since invariably the inner gasket installation is accompanied by some longitudinal extension or compression. If the gasket is wedged in with a roller-type of tool moving from corner to corner with any excess cut off, the gasket may be stretched by as much as 10%, but if the gasket is precut and worked in from each corner to the middle of an edge, no stretching will result. In the event that the gasket is precut to a slightly longer length and forced into position, a net compression results. This might be done in an attempt to "strengthen" the glazing, but



**75 Duro-neoprene Gasket**



**Neoprene Sponge Gasket**

Figure 4.4-3 Neoprene Gasket Cross-sections

if exercise forces are used may produce a failure of either the glass or the spacer or both.

In order to characterize the gasket support, a series of tests were devised to determine the elastic properties for four modes of motion at the glass plate boundary. Specifically stiffnesses are:

1. Out-of-plane or normal stiffness which is due directly to the compressive gasket stiffness.
2. In-plane-normal or pull-out stiffness which is due to the shear stiffness of the gasket in a direction perpendicular to a boundary edge.
3. In-plane-longitudinal stiffness which is also due to the gasket shear stiffness but in a direction parallel to a boundary edge.
4. Rotational stiffness or the moment restraint about an axis parallel to the edge which is due to a combination of shear and compression in the gasket.

In general, of course, there are 6 degrees of freedom at the edge and 6 corresponding restraint stiffnesses can be defined. The above 4 are the predominant ones for static deflection or dynamic out-of-plane response. Along with the edge restraints, additional constraint forces are produced between the two glass plates in the window as a result of the spacer. Unfortunately, short of disassembling the vision light itself, there is no simple means for experimentally measuring the associated stiffnesses so that in all analytical work, extreme values have been arbitrarily assumed.

Wey (104) in his thesis described a series of tests to determine the above four stiffnesses and then using both continuous and discrete models for the window and framework compared the measured deflections with the predicted for static pressures up to 20 psf (958 Pa). This section presents a summary of the boundary stiffness results along with several observations on the behavior. The measurements were performed using a small 1 ft. by 1 ft. (30.5 cm by 30.5 cm) architectural evaluation sample of the actual window which, other than its area, was of identical construction. An edge of the test sample was glazed in the actual test fixture for the inplane and rotational stiffness tests. Loads were supplied by weights using cables and pulleys, and deflections were measured with precision dial gages. The gasket compression tests were performed in a Universal Test Machine, again using the glass window sample and a short piece of the mullion section.

#### Out-of-Plane Stiffness:

The out-of-plane gasket stiffnesses were measured by applying a direct compressive load to the gasket and measuring the resulting deformation. The gasket was mounted in a section of mullion to duplicate installed conditions. Figures 4.4-4 and 4.4-5 show typical results for the neoprene sponge and 75 duro-neoprene. In both cases the behavior is nonlinear and hardening although the effect is much less severe than expected. The most significant result is the observation that the sponge material is apparently stiffer as evidenced by the steeper average slope:

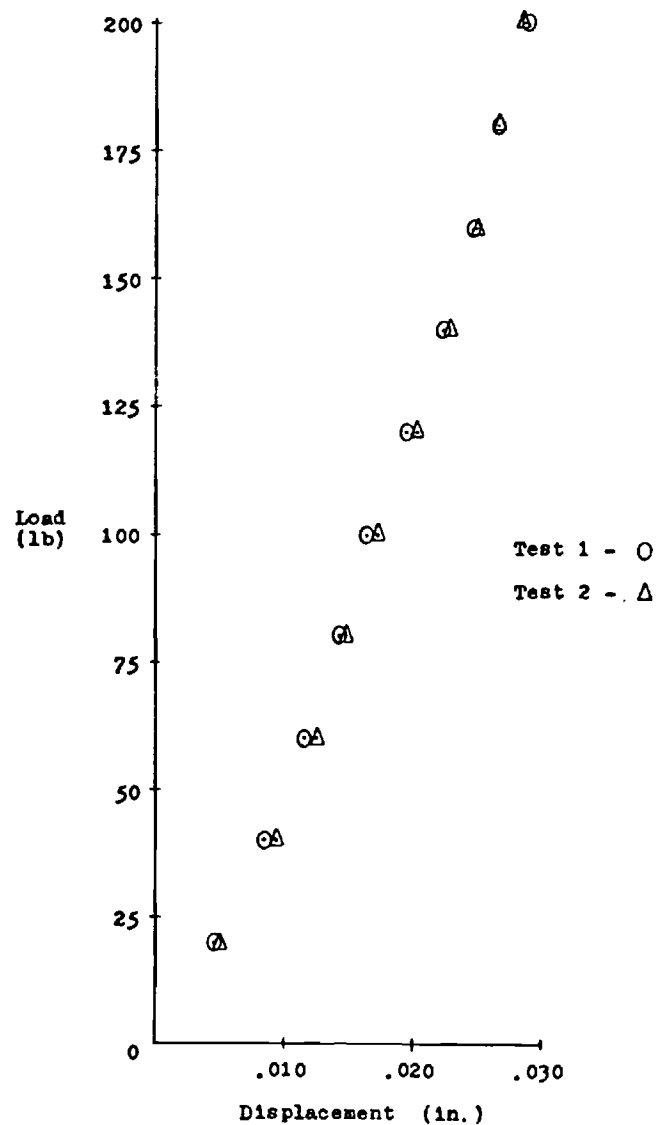


Figure 4.4-4 Gasket Compression Test for Neoprene Sponge Gasket Material

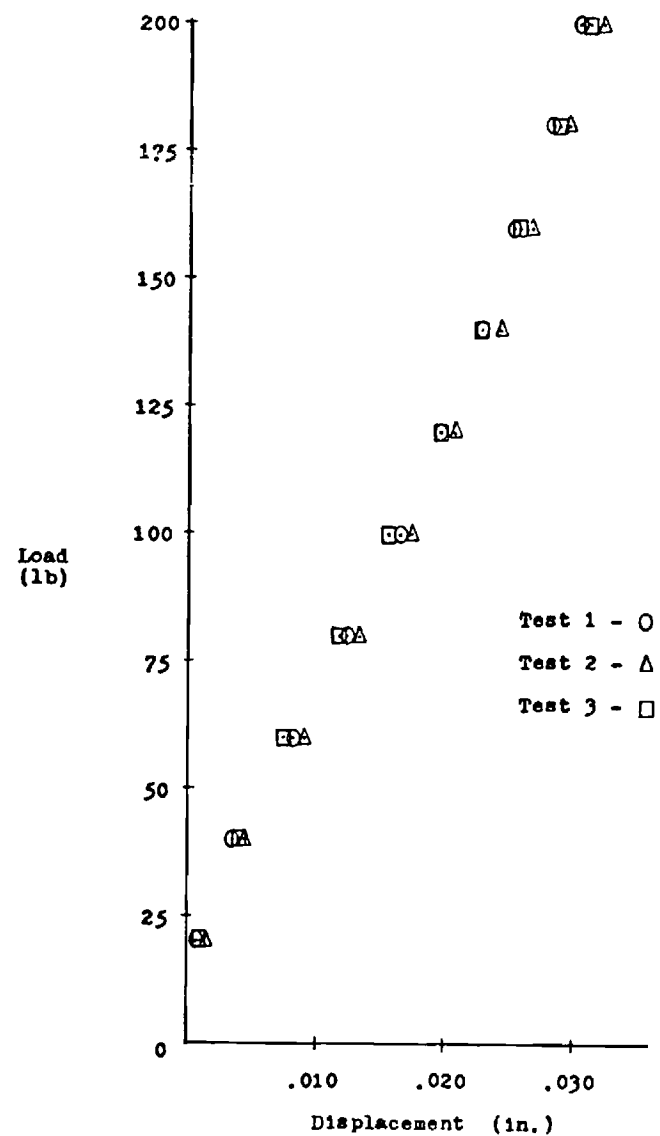
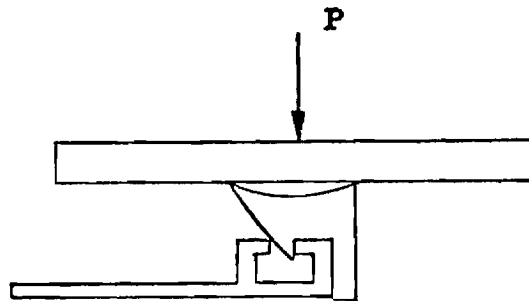
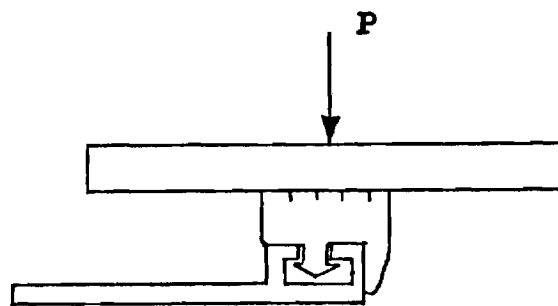


Figure 4.4-5 Gasket Compression Test for 75-duro Neoprene Gasket



**75 Duro-neoprene Gasket Under Load**



**Neoprene Sponge Gasket Under Load**

Figure 4.4-6 Detail of Gaskets Under Normal Compression Load Showing Different Contact Areas

<u>Material</u>	<u>Average Slope, lb/in/in, (kN/m/m)</u>
Neoprene Sponge	335 (2310)
Solid Neoprene (75 duro)	222 (1530)

The explanation for this seemingly contradictory behavior lies in the manner of deformation for each gasket. As shown in Figure 4.4-6, the two gaskets behave quite differently under load. The sponge material is in full contact with the glass at all times so the stress is fairly uniform. The solid neoprene because of its shape is in contact with the glass along two lips and considerable localized bending occurs. Even at a compressive load equivalent to three times the expected service load, a significant cup shape remained. When the material was loaded until full contact was established the expected sharp increase in stiffness was observed, however, this condition was not judged typical.

Normal In-Plane Stiffness: The in-plane stiffness normal to a glazing edge was measured using a small specimen light mounted in the full-scale test rig to simulate actual glazing conditions. As might be anticipated, after the load was increased beyond a certain value the light began to slip and, if the load was maintained, would pull out of the glazing. The pull-out load was found to depend on several parameters that were difficult to control, e.g., installation technique, load history. Experience with actual field installations generally reveals that after a window has been installed for a few months, the glazing gasket becomes tightly stuck to the glass. To simulate this condition, a cyano-acrylate adhesive was used to adhere the glass to the gasket and prevent slipping during the tests. Figure 4.4-7 shows results from several tests. As a result of the initial gasket



configuration after gluing and localized gasket tip behavior a measurable deflection does not occur until a threshold load has been reached. The average stiffness for loads above the threshold is 342 lb/in/in ( 2360 kN/m/m) although as before a noticeable hardening behavior is evident.

In Plane Longitudinal Stiffness: The gasket stiffness parallel to the window edge was determined in a manner similar to that above. Figure 4.4-8 shows typical results which, in contrast to the normal in-plane case, do not exhibit the threshold feature. In this case the gasket tip distortion phenomenon is not present since the load is acting along the edge. The curves show an initial softening followed ultimately by slippage of a maximum load. It is apparent that the inherent hardening characteristic of the rubber is overcome by slippage. The average stiffness for shear loads above 1 lb/in ( 0.2 kN/m) is 318 lb/in/in (2190 kN/m/m).

Rotational-Stiffness: Boundary rotational stiffness was determined in a manner similar to the above two cases, but in this case the loading was applied so as to produce a known edge moment and the resultant rotations were measured. The gasket loading involves a combination of compression and shearing normal to the edge, and as for the in-plane stiffness normal to an edge, the effect of gasket-window slippage is evident with and without slippage. In this case, when the gasket is free to slip, localized lip slippage increases the contact surface and a higher rotational stiffness results. When the gasket is glued, the primary deformation mode involves localized bending of the gasket lips and a considerably lower rotational stiffness results. Figure 4.4-9 shows the measured behavior.

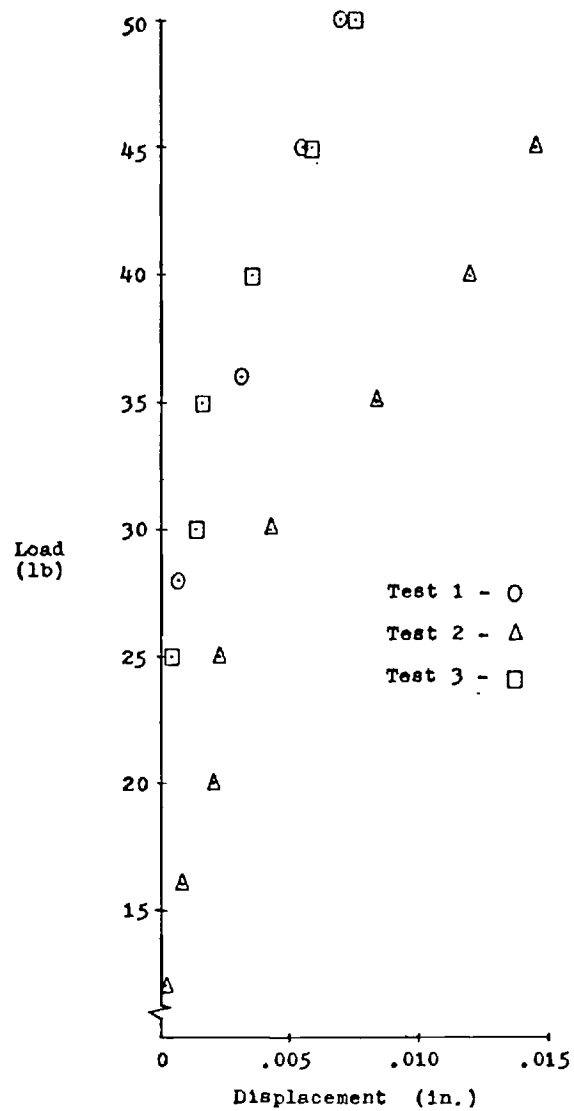


Figure 4.4-7 In-plane Stiffness Normal to the Edge - No Gasket Slippage

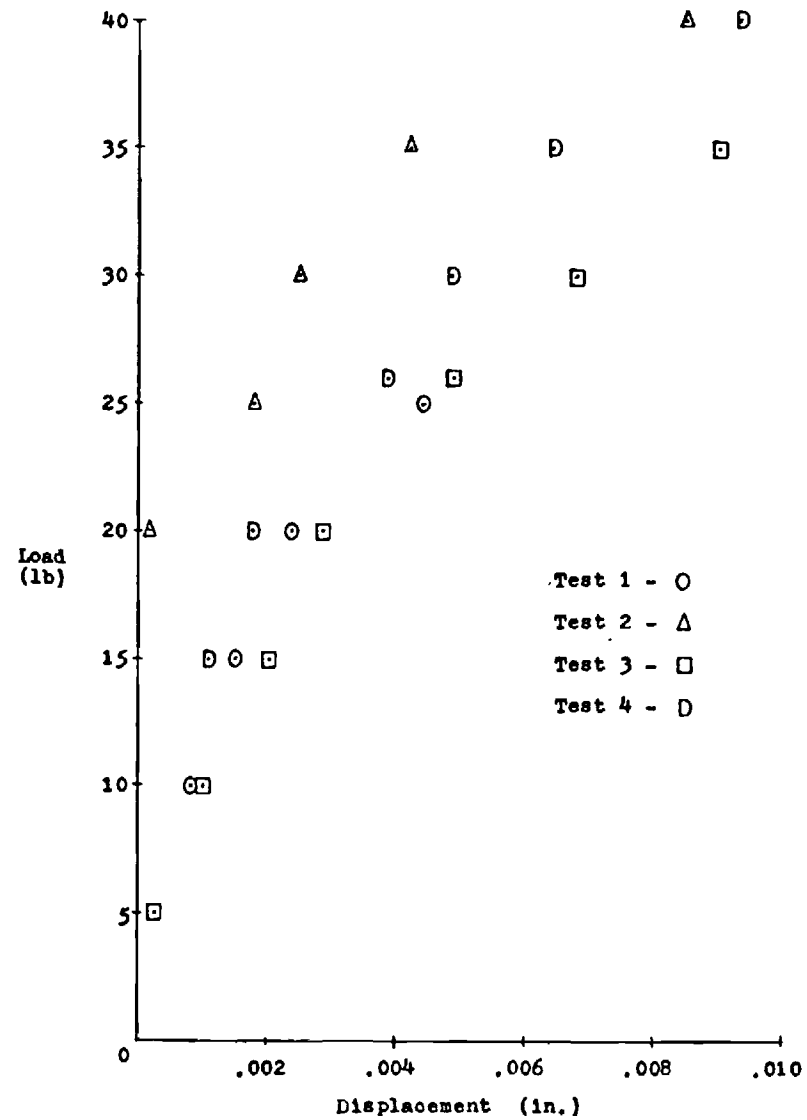
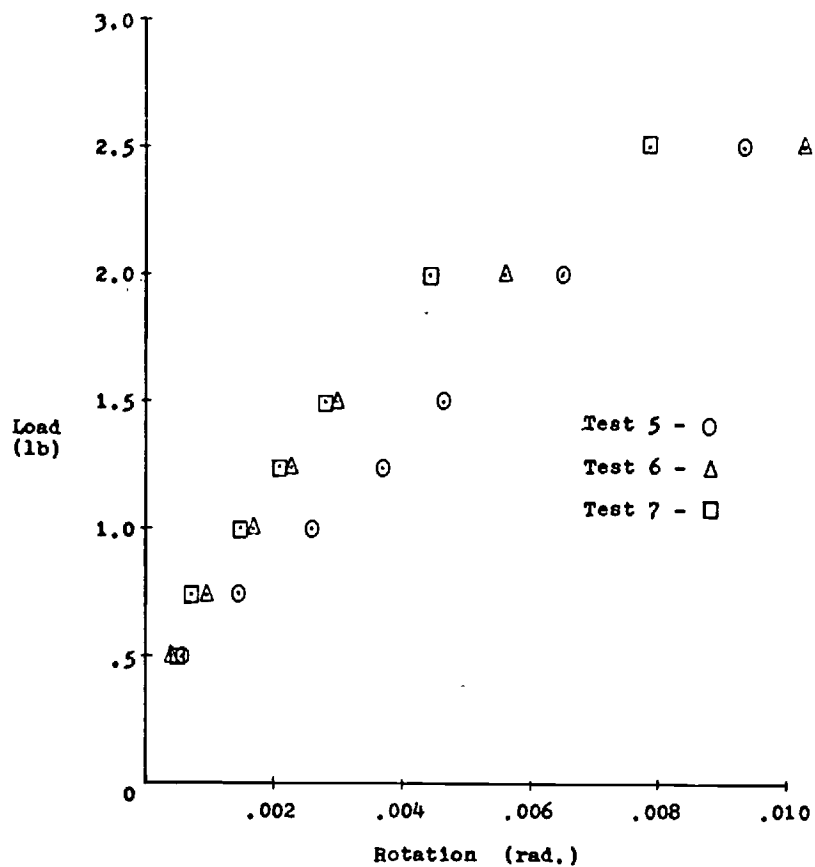
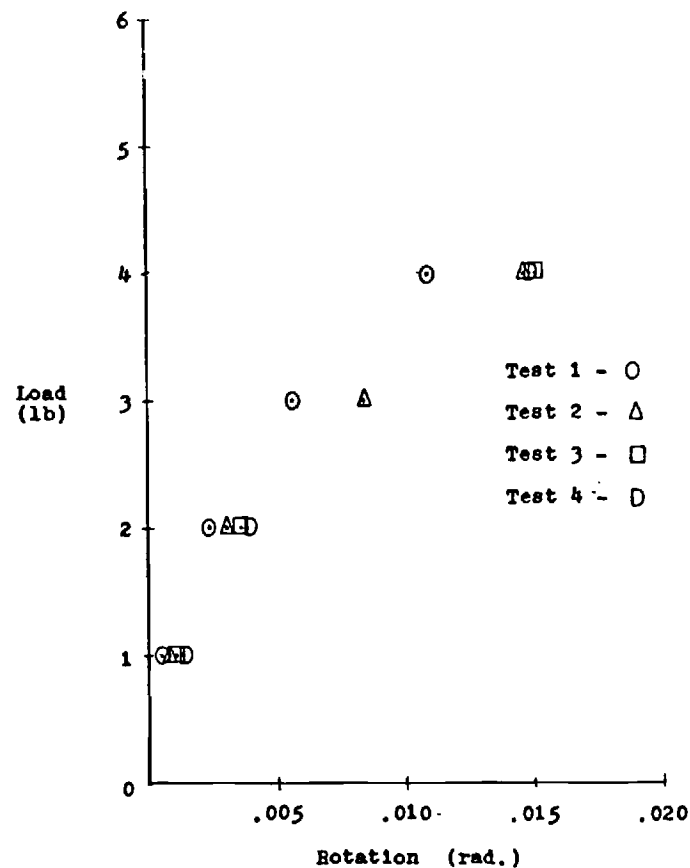


Figure 4.4-8 In-plane Stiffness Parallel to the Edge - No Gasket Slippage



(a) Gasket Free to Slip



(b) Gasket Not Slipping

Figure 4.4-9 Edge Moment Restraint Stiffnesses

The gasket properties as found above were used in the analytical modelling in Chapter 3. They are summarized below:

Mode	Stiffness, lb/in/in, (kN/m/m)
Out-of-Plane	
Neoprene sponge	335 (2310)
Neoprene, 75-duro	222 (1530)
In-Plane, Normal (no slip)	342 (2360)
In-Plane, Longitudinal	318 (2190)
Rotational	
Slip	305 (2100)
No slip	159 (1100)

Of these, the out-of-Plane and Rotational stiffnesses are the most important in their effect on static and dynamic window behavior. The inplane stiffnesses are so small compared to the inplane window stiffness as to be insignificant (as they are supposed to be by design in order to avoid introduction of detrimental inplane loading).

#### 4.4(b) Static Test Results

Static pressure tests were conducted on the basic window unit described in 4.4(a) while it was mounted in the full scale test fixture. As noted, the support conditions were arranged to closely duplicate a typical mounting configuration. The instrumentation consisted of an array of linear variable differential transformers (LVDT's) for the measurement of deflection and numerous electrical resistance foil strain gages bonded directly to the glass and framing

surfaces. The LVDT's were Hewlett-Packard DCDT-series transducers which are dc powered and provide a dc output with from 6 to 30 V/in (15-75V/cm) sensitivity. The units were used on both sides of the window and for reference purposes were mounted to stiff support bars firmly anchored to the test fixture frame. As many as 8 units were used simultaneously. The strain gages were 350 ohm metal foil types with a large gage area selected so as to minimize self-heating at the 6V excitation level used. Sixteen gages were typically used in the tests. The displacement and strain gage layout is shown in Fig. 4.4-10.

An HP9825A calculator-based data acquisition system using an HP3490 voltmeter and HP3495 reed scanner was used to acquire, process and output the test data. Strain readings were computed directly from the microvolt-level Wheatstone bridge outputs and bridge balance was accounted for simply by subtracting the initially measured unbalance voltage at zero load using the calculator program. In this way, a typical test required no more than 15 sec. measurement time at each pressure load.

Test Configurations: The static tests were designed to determine as much as possible about the behavior of a typical insulating glass window when supported in a manner typical of that found in lightweight cladding designs. For reasons of economy, both of time and effort, the reported test program was limited to a particular window and mounting configuration, but extensive measurements were made to ascertain the behavior mechanisms under so-called "standard" installation conditions as well as under "abnormal" conditions. While this approach is not

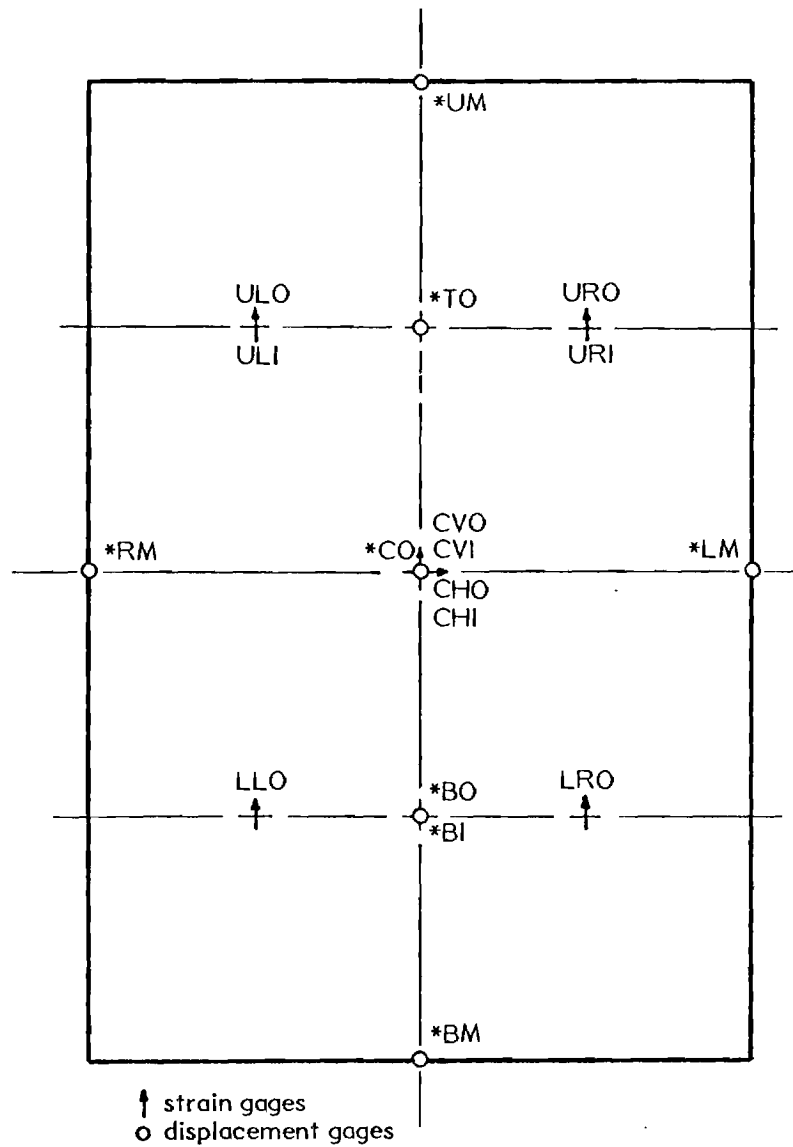


Figure 4.4-10 LVDT Deflection and Strain Gage Transducer Locations on the Test Specimen

as general as might be desired, it does provide certain design sensitivity information for variations in conditions from the standard or nominal configuration. In the present study two major parameters were considered:

- (1) Glazing support conditions - in particular, the effect of increased or decreased glazing support due to precompression of the neoprene glazing gasket.
- (2) Window integrity - the effect on deformation and stresses of maintaining a hermetically sealed volume (or gap) of gas between the two glass panes in an insulating window.

The first case is directly concerned with the effect of gasket support conditions on static response. Rather than consider the more complex case of variable structural geometry (which has been extensively explored in the analytical developments in Chapter 3), the present tests were concerned with the effect on static performance when more or less gasket material is forced in the glazing thus providing varying degrees of gasket precompression. This is typically the result of differences in glazing techniques and for certain designs the effect on performance may be substantial. The second case deals with the case of window performance when the seal between window panes is broken, either through aging or mechanical failure of the spacer and sealant.

The first case was simulated by using a standard glazing configuration in which the gasket length was cut exactly equal to the window perimeter. Then configurations with 10% less and 5% more gasket lengths were used. These values were the practical limits of the

amounts of gasket that could be removed without leaving gaps or the extra length that could be forced into the glazing.

The second case was simulated by boring four 3/8 in (9.5mm) holes through the insulating window spacer and into the internal volume. These were located near each corner and were fitted with removable rubber plugs to allow simulation of either a "vented" or "sealed" configuration. One plug was fitted with a pressure tap to allow measurement of the internal pressure under "sealed" conditions.

The test window with strain gages attached was mounted in the framing and glazed with the specified amount of gasket material. The window was supported on glazing blocks and carefully centered with a uniform penetration in the mullion and spandrel sections of approximately 0.5in (13mm). A period of 48 hours was allowed before testing to provide time for the gasket material to relax and conform to the support framing.

Standard Gasket Length-Sealed Spacer: This represents the nominal design configuration to which the balance of the tests are referenced. Figure 4.4-11 shows the deflections measured at the center of the outer (\*CO) and inner (\*CI) panes as well as at the mid-height of the right mullion (\*RM). Two results are immediately clear:

- (1) The inner and outer panes deflect almost equal amounts,
- (2) The deflection is almost linear with pressure with only a slight hardening evident.

The difference between the inner and outer pane movement is less than 3% when the window spacer seals the enclosed volume. The relative



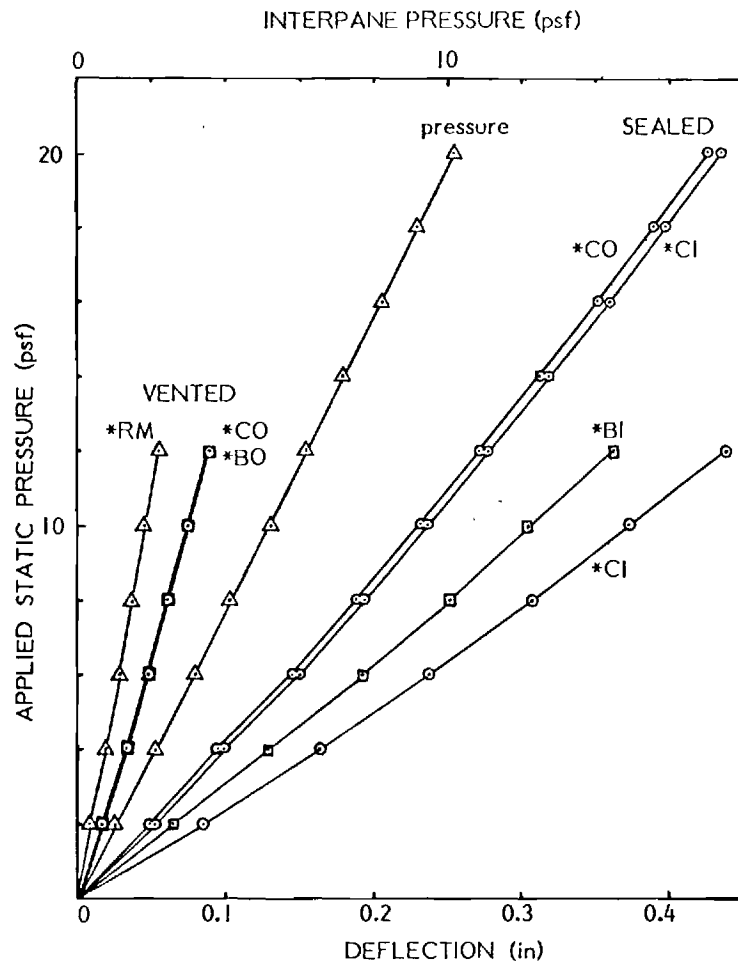


Figure 4.4-11 Deflection and Interpane Pressure Versus Static Pressure Applied Across the Window

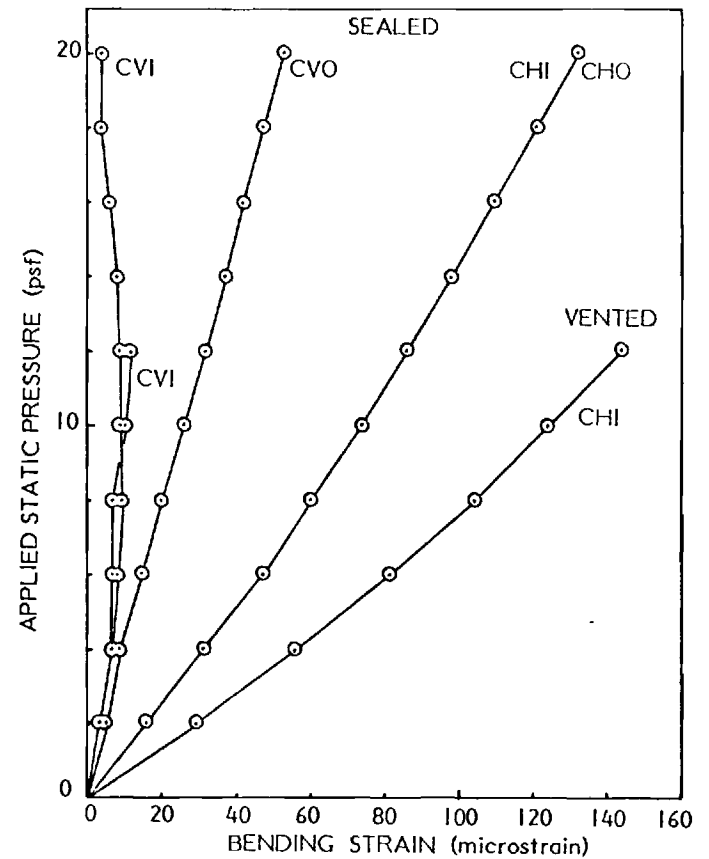


Figure 4.4-12 Glass Strain Versus Applied Pressure for the Standard Gasket Configuration

midpoint deflection (difference between \*CI or \*CO and \*RM) is approximately linear with pressure over a range of deflection that is on the order of the equivalent double-plate thickness, 0.32in (8.1mm), obtained by assuming the two panes act together in bending with the neutral surface located halfway between. Figure 4.4-11 also shows the interpane pressure reference to outside pressure as a function of the applied differential pressure across the complete window. The values are almost exactly one half the total differential pressures which, like the deflection data, indicates that the load is almost equally divided between the two panes.

Figure 4.4-12 shows the variation in strain with static pressure for selected locations. Again as for displacements the same two features are evident in the strain from gages located at the center with horizontal orientation (CHI, CHO). In addition, the data show that the strains in the vertical direction at the center (CVI, CVO) are much smaller. Thus, the deflected shape primarily involves curvature in the horizontal direction as would be expected for a plate of that aspect ratio. The table below which lists the strain for each location bears this out.

Strain Levels at 20 psf (958Pa) - Sealed Spacer

<u>Location</u>	<u>Microstrain</u>	<u>Location</u>	<u>Microstrain</u>	<u>Location</u>	<u>Microstrain</u>
CVO	53	URO	37	LRO	36
CHO	127	URI	-37	LRI	-30
CVI	-3	ULO	36	LLO	36
CHI	-132	ULI	-32	LLI	-32

The peak strains are clearly well below the ultimate values for glass.

Standard Gasket Length - Vented Spacer: This configuration represents the situation where the spacer seal has been compromised so that the pressure load is developed across one pane only. Figure 4.4-11 shows the deflections at the midpoint of the window as well as the quarter-height location (\*B0, \*BI). Pressures were limited to 12 psf (575 Pa) in order to maintain deflections within the linear range (less than about one equivalent plate thickness). It is clear from the plot that the deflection occurs almost entirely in the inner plate across which the pressure is developed, while the outer pane deflection is about that of the mullion midpoint. Similar curves in Figure 4.4-12 show the same behavior for the midpoint strains.

The strain levels at 12 psf (575Pa) are shown below:

Strain Levels at 12 psf (575 Pa) -Vented Spacer

<u>Location</u>	<u>Microstrain</u>	<u>Location</u>	<u>Microstrain</u>	<u>Location</u>	<u>Microstrain</u>
CVO	4	URO	3	LRO	6
CHO	12	URI	-39	LRI	-34
CVI	-1	ULO	2	LLO	4
CHI	-144	ULI	-37	LLI	-37

It is apparent that the outer pane deformation is due to a combination of mullion deflection and edge moment transmitted through the spacer.

Deflected Shape: The measured deflections presented above and reported in detail in (104) show that when either in a sealed or vented configuration, the window behaves like a flat plate of equivalent thickness with its edges supported by boundary frame members.

The analytical representation presented in Chapter 3 when used with member properties for the test setup, predict the behavior with high accuracy. In addition, a classical solution developed in (103) for a more simplified model also agrees well with the measured data.

Comparison of Predicted and Measured Deflection at a  
Pressure Load of 20 psf (958 Pa) - Sealed

Location	Discrete Model (Chapt. 3)	Continuous Model (Ref. 103)	Measured Results
Midpoint	0.460 (11.7)	0.436 (11.1)	0.430 (10.9)
Mullion	0.054 (1.4)	0.049 (1.2)	0.086 (2.2)

Note: deflections are: in(mm)

The above summary shows excellent agreement for midpoint deflections and moderate agreement at the boundary frame.

Gasket Effects: For both the sealed and vented spacer conditions, static tests were run for three different gasket configurations:

- (1) Standard Length Gasket - both the inner and outer gaskets were cut to the exact glazing dimensions and installed accordingly.
- (2) 10% Less Gasket - the 75 duro neoprene gasket was cut 10% under length and during installation was stretched to complete the glazing.
- (3) 5% Extra Gasket - the 75 duro neoprene gasket was cut 5% over length and during installation was compressed to fit the glazing.

These configurations represent the nominal and the extreme conditions likely to be encountered. For (2) the stretching reduces the cross

section dimensions and thus the pressure exerted by the gasket on the window. This reduces the effective moment restraint along the edge. In (3) the axial compression of the gasket that occurs as the extra length is forced in the glazing increases the cross section dimensions and increases the gasket pressure on the window, thus increasing the edge restraint.

The static deformation results for the Standard Length Gasket are shown in Fig. 4.4-13 for both sealed and vented conditions. The data are similar to those in Fig. 4.4-11 and show the same features. Fig. 4.4-14 shows the deformation for 10% Less Gasket and it is apparent that the behavior is close to the same as for the sealed configuration. On the other hand there is a definite increase in the inner pane deflection for the vented configuration. While not shown, a decrease in deflection occurs for the 5% Extra Gasket - vented configuration.

The plots in Fig. 4.4-15 show the effects of gasket properties on the static response more clearly. As noted earlier, the inner and outer pane displacements are essentially the same for the standard gasket configuration when the interpane space is sealed. The same behavior was observed for both the 10% Less Gasket and 5% Extra Gasket. Fig. 4.4-15(a) shows a comparison between the inner pane center displacement for the different gasket configurations when compared to the standard configuration. The abscissa is the difference between the particular inner-pane deflection and the value for the Standard Gasket configuration. It is clear that there is no significant variation with gasket properties over the pressure range tested (the

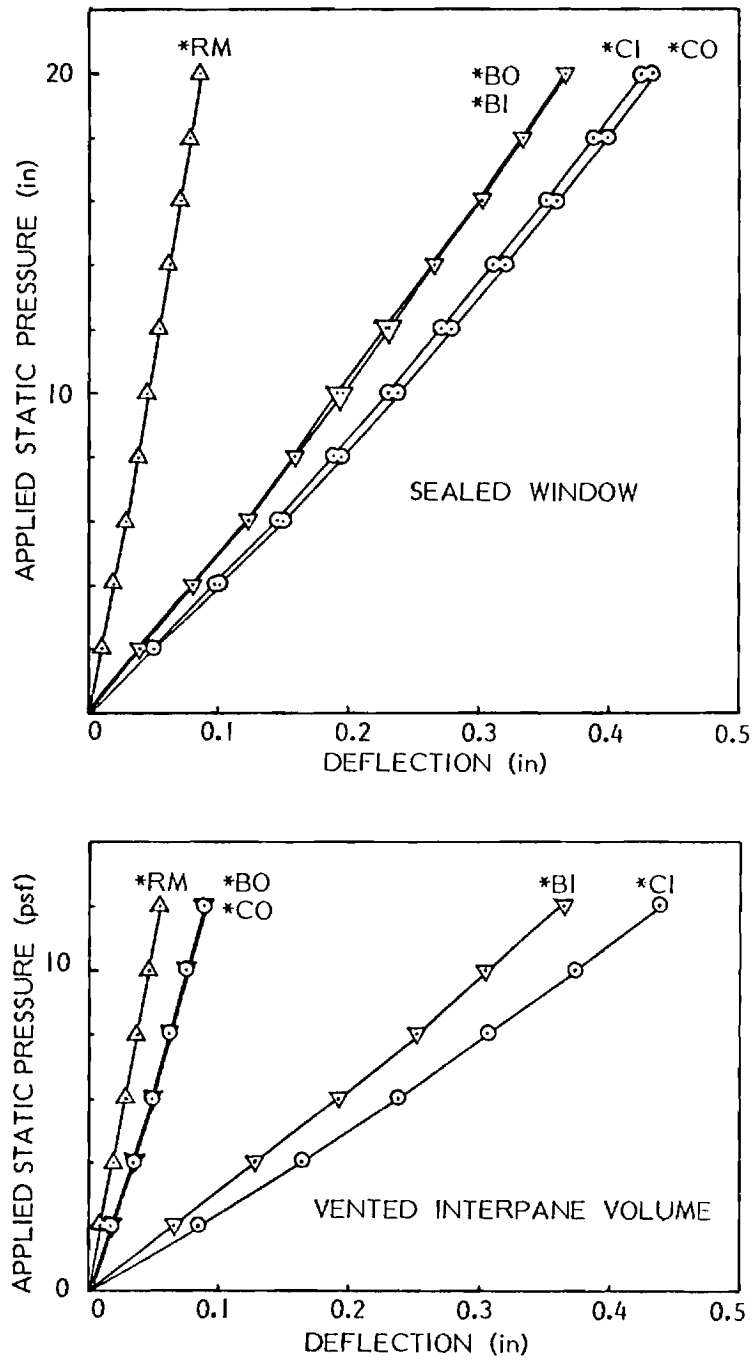


Figure 4.4-13 Static Deformation with Applied Pressure for the Standard Gasket Configuration

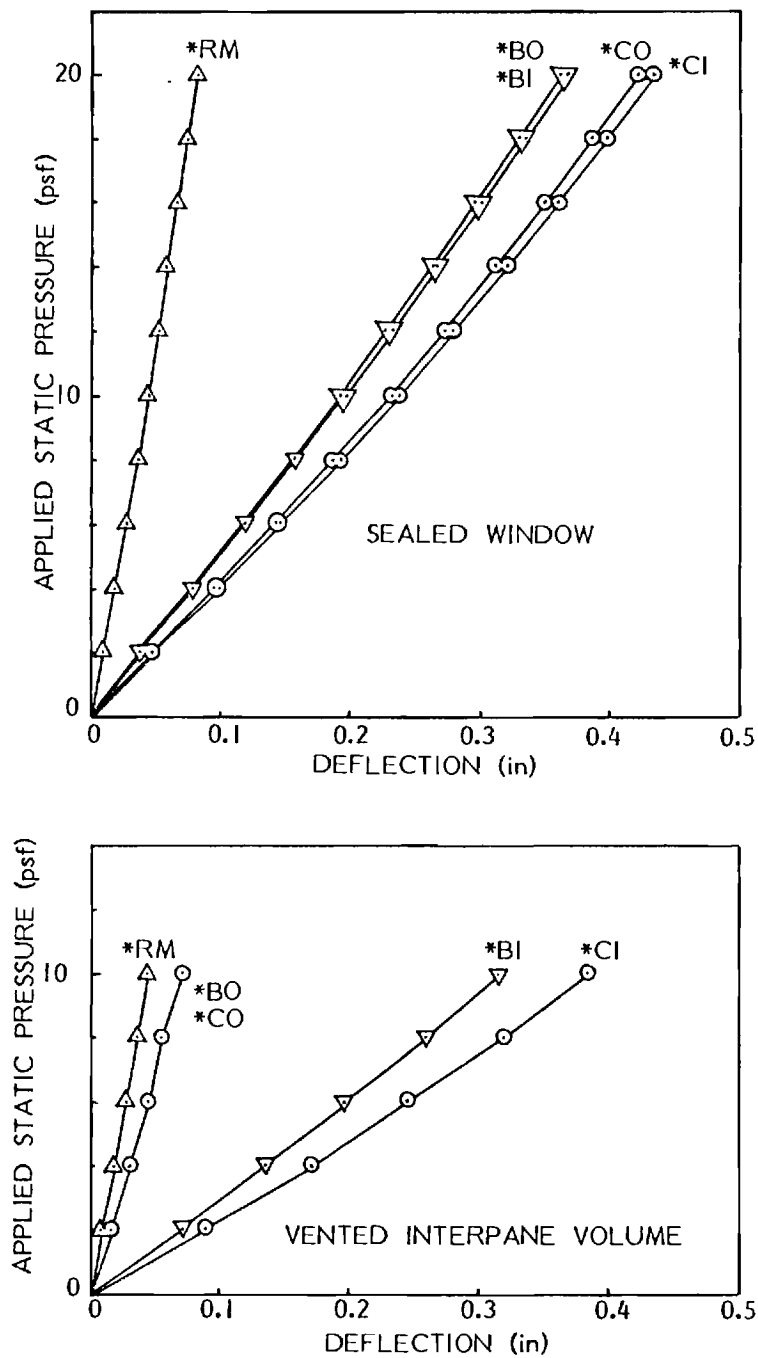
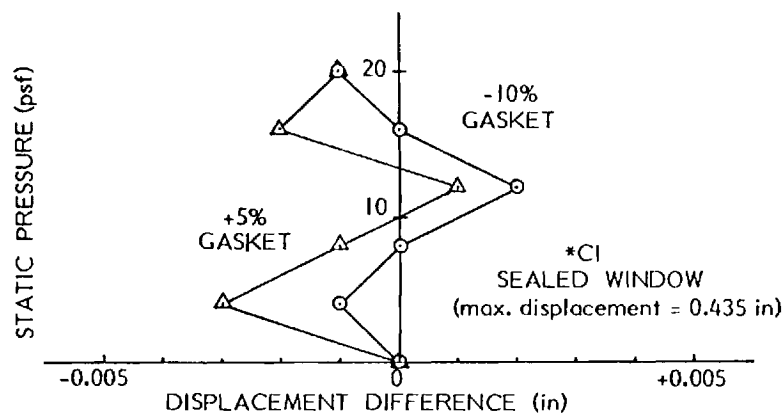
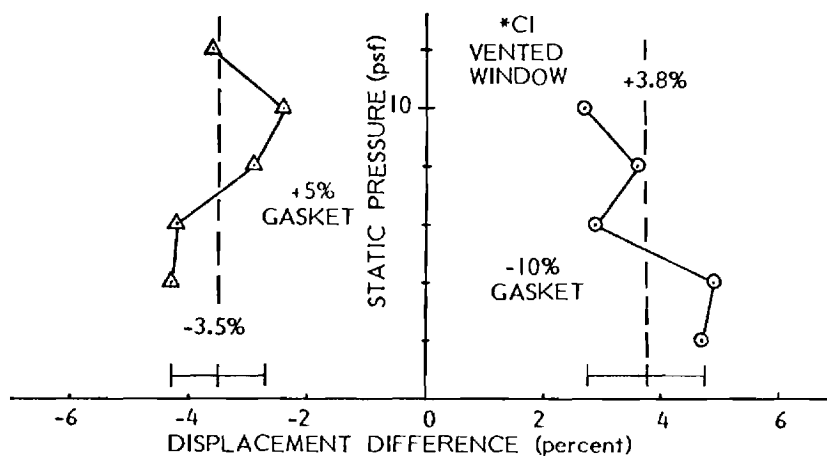


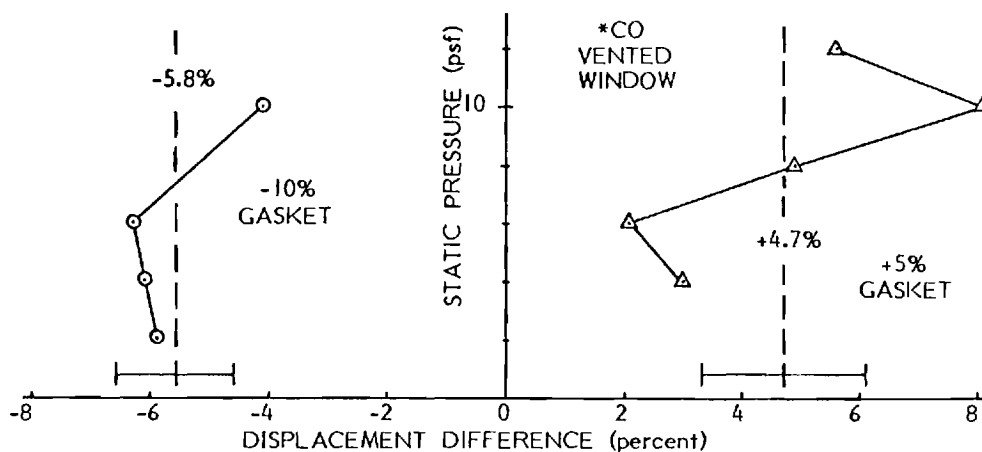
Figure 4.4-14 Static Deformation with Applied Pressure for the 10% Less Gasket Configuration



(a) Inner Pane Deflection Referenced to Standard Gasket - Sealed



(b) Inner Pane Deflection Referenced to Standard Gasket - Vented



(c) Outer Pane Deflection Referenced to Standard Gasket - Vented

Figure 4.4-15 Effect of Gasket Properties on the Static Deflection



variation is less than 0.5% at the maximum pressure). When a similar plot is made for the window with the interpane space vented as shown in Fig 4.4-15(b), a distinct difference in behavior is evident for both gasket configurations. With the reduced gasket the inner pane deflection is increased by an average of 3.8% over the standard configuration, while a 3.5% reduction is shown for the extra gasket case. Fig. 4.4-15(c) shows the same type of plot for the deflection at the midpoint of the outer pane, but now the behavior is just the opposite with the increased deflection associated with the extra gasket.

The most obvious conclusion from these data is that the effect of the glazing gasket on the static response is not significant for the nominal case of a sealed interpane space. Even when the seal is compromised, the effect on deflections and strains is not pronounced.

The particular reversal of the effect between the inner and the outer pane noted above can be explained by reference to the test configuration. The inner pane is supported directly by the neoprene foam gasket of higher "effective" stiffness while the 75 duro gasket with somewhat lower "effective" stiffness is used to make the outer glazing. Thus as more (or less) compression of the inner gasket with a consequent increase (or decrease) in the rotational restraint for the inner pane edge. Since the two panes are still bonded together, an increase in the inner gasket restraint force will reduce the corresponding inner glass deformation and shift more load to the outer pane which exhibits increased deformation. The reverse holds for a reduction in gasket restraint levels.

#### 4.5 Dynamic Testing

Following completion of the static tests, a series of dynamic response measurements were carried out using the same window unit and mounting arrangement. The simple static pressurization procedure used for the static testing was replaced with dynamic excitation. The same strain gage and LVDT instrumentation was employed, but the essentially static-level logging instrumentation was replaced by a minicomputer-based dynamic response analysis system. The test procedure and measurement methodology are described in the first subsection, and the results obtained from the tests are presented in the second.

##### 4.5(a) Measurement and Analysis Methodology

The strain gage and LVDT instrumentation used for the static tests and shown schematically in Fig. 4.4-10 was also used for the dynamic response measurements. The LVDT's were equipped with light weight cores and were preloaded with small springs to keep them in contact with the glass plate over the range of frequencies and accelerations encountered in the tests. Thus, the output from these transducers is proportional to the transverse displacement of the glass at that point. The strain gage outputs, on the other hand, are due to glass strain at the measurement point. For the range of deflections and loads used for the tests, the strain is due primarily to the plate curvature, and hence the gage output can be directly related to the plate deflection (this is not true when a substantial fraction of the strain is due to midplane stretching or to nonlinear effects). Strain gages are relatively inexpensive and simple to use and since in contrast to the LVDT's require no special reference mounting, they have been employed as the primary dynamic

response transducers. The LVDT's were used as a secondary source of information to resolve ambiguities in the strain signals. It should be noted that inertial accelerometers were not used because of their relatively high cost and the somewhat inaccessibility of the rear surface of the glass inside the test chamber.

The strain gage and LVDT signals were amplified and low-pass filtered and then were digitized using the HP5451B Fourier Analyzer described in Chapter 2. The analog-to-digital converter could accept only two signals simultaneously so that only pairs of transducers could be handled at a time. While substantially increasing the testing time, this restriction did not alter the basic measurement and analysis plan described below.

Modal Analysis: The primary objective of the dynamic response testing was to determine the basic modal characteristics of the window assembly when mounted in a manner similar to that in practice and when subjected to similar loading. In addition, the effect of systematic variations in the boundary restraint was evaluated. The classical approach to this type of problem has been to provide a means for dynamically exciting the structure and then to measure and identify the response at each of the natural modes of vibration within the frequency range of interest. Typically, a harmonic excitation is used and the frequency swept slowly over a range while the phase and amplitude measurements at selected points are used to identify normal mode response. In this connection various methods have been used to provide excitation of selected normal modes.

More recently, additional signal processing and analysis techniques

have become available that reduce the measurement time and allow a greater variety of excitation methods to be used. These are based on the "fast Fourier transform" or FFT algorithm and as such are implemented using digital time series realizations of the excitation and response signals. The methodology has been described in Chapter 2 in connection with the full scale building measurements and will not be repeated here. As before, the present work is based largely around the capabilities and features of the HP5451B Fourier Analyzer system which was used for the acquisition and processing of the dynamic response time series.

The measurement procedure was based on the linear response model for the system so that an influence or a transfer function could be measured which would relate the response at one point to a unit impulse applied at another. If sufficient points on the structure are chosen and a matrix of transfer functions are measured relating input at one point to output at another, it can be shown (84) that this is sufficient information to identify the modal response characteristics for most cases. Furthermore, it is possible to accomplish this using only a single row or column of the matrix, that is, using transfer functions obtained for a single input point and multiple response locations or for a single response measurement point and multiple input locations. The modal frequencies, and damping as well as the mode shapes can be obtained from the transfer functions by a variety of techniques based on certain features of the real and imaginary or magnitude and phase components as described in Chapter 2.

In the field studies reported in Chapter 2 and initially in the

present lab work, the modal identification process was carried out by examination of each of the transfer functions manually. Since these processes require relatively widely separated modal frequencies for accuracy, the results were limited in these studies to only the first two or three modes. After acquisition of the transfer functions and completion of the laboratory test program, access became available to a powerful and sophisticated analytically-based modal identification system that could be incorporated directly in the HP5451B Fourier Analyzer. This system (basically the HP Modal Analysis software package) is capable of mathematically fitting multiple degree of freedom analytical transfer functions to the measured data and adjusting the modal characteristics until the best match is made in a least squares sense. This ability has allowed more careful examination of the transfer function data than was possible manually, especially in situations where the modes are somewhat closely spaced. Unfortunately, since the system became available well after completion of the actual laboratory measurements, it has not been possible to go back and obtain improved quality data when the original data was found lacking for purposes of the modal analysis. It should be noted that this type of activity is now underway outside the scope of the present grant and will be reported on shortly.

Input Excitation: The full-scale test facility described in Section 4.2 was used to produce the dynamic excitation required for a response analysis of the window system. Two different procedures were employed depending which aspects of the response were being analyzed. The first procedure consisted of using the closed-loop servo-controlled hydraulic

system to produce time varying pneumatic pressures across the test window. As described in Section 4.3, a time varying electrical signal was applied to the control system which in turn caused the hydraulic actuators to displace the two large doors or pneumatic pistons located on the rear side of the test fixture. The door motion produced corresponding time varying pressure fluctuations in the test facility and thus provided the desired fluctuating differential pressure across the test window.

A variety of signal sources could be used including actual recorded cladding pressure waveforms. For purposed of the present modal analysis, however, it was desirable to provide simultaneous excitation over a band of frequencies encompassing the first several modes for the window assembly. Consequently, a band-limited stationary random signal with "white" spectral characteristics and a Gaussian amplitude distribution was used as the excitation source. When coupled with the digital time series analysis capability described above this allowed simultaneous measurement and analysis over a band of frequencies rather than at a single frequency as with the swept-sine approach. Appropriate ensemble averaging was used to provide input and output response measurements with a minimum acceptable level of variance. A single transducer (strain gage or LVDT) was defined as the input reference and the corresponding transfer function matrix row was determined by successively measuring the transfer functions between this and all other transducers (taken two at a time because of signal digitizing limitations noted earlier).

While this manner of excitation closely resembled the actual service conditions in character, the symmetric placement of the window in the

fixture coupled with the symmetric arrangement of the driving doors combined to strongly favor excitation of symmetric modes only. One of the modes of most interest, namely the 1'-2' mode in which the panes can move in opposite directions to each other, is not a symmetric mode and it was not possible to clearly excite this type of response using the uniform random pressure method. As a result, an alternate method was developed and used along with the random excitation. This method was basically an impulse excitation applied unsymmetrically at a point on the window. Again, a single transducer near the application point for the impulse was defined as the input reference and the corresponding transfer function matrix row was determined as above. Due to the relative inaccessibility of the window surface on the test chamber inside, a single input point was chosen and multiple response points were used. The impulse was applied by hand using a large mallet with a suitably cushioned head to provide the necessary impulse energy over the frequency range of interest. Ensemble averaging was also used to insure an acceptable level of variance in the data.

#### 4.5(b) Dynamic Response Results

Test Method: As described earlier, the same strain and displacement transducers (shown in Fig.4.4-10) were used for static and dynamic tests. In contrast to the static tests, dynamic analysis requires considerably more data be acquired from each transducer in order to adequately define the time-frequency behavior. For example, in order to obtain a frequency resolution of  $\Delta f$ , one must analyze a sample record roughly  $1/\Delta f$  long (see Chapter 2) which for the desired resolution of  $0.1 \text{ Hz}$  in the present tests resulted in collection of data typically over periods of 10 sec or

more. While this was readily done using the HP5451B analyzer, the off-line storage of this data for future analysis proved especially burdensome since the system was not originally equipped with any mass storage devices; the only medium available was 1" punched paper tape. As a result, a compact (from the standpoint of amount of data) test method was developed as follows. With reference to Fig. 4. 4-10, the 2 center and 4 lower quarter point strain gage locations (CHx, LLx, LRx) and the 5 LVDT locations on the midline (\*xx) were taken as the primary measurement sensors and only data from these were saved on paper tape. In this way, by comparing the phase between pairs of sensors it is possible to detect modes with up to 3 node points in the vertical and up to 2 in the transverse directions (up to the 3-2 mode):

Next, using either the random pressure or the point impulse excitation methods, a single sensor (LR0 strain or \* B0 LVDT) was selected as the input reference and the transfer function between that and the other locations (output points) were computed digitally in the HP5451B. For the strain gages this required 5 transfer functions which in turn defined a row in the transfer function matrix. Four transfer functions were computed in a similar manner for the LVDT's. For each test, ensemble averaging was performed to reduce the variance and the data were rejected if the computed coherence function was of poor quality (e.g. substantially below about 0.85).

Finally, the transfer functions were analyzed, first by hand and much later using the HP Modal Analysis software, to determine



the modal response characteristics. As described in 2.5(c) the modal frequencies and damping can be identified by certain features of the real and imaginary components of the transfer function, and for a given modal frequency, the phase between the input and each output point can be computed. This latter information is particularly critical for it effectively defines the mode shape. For example, in a pure (normal) mode, the frequency is identified by zero real and maximum imaginary components. Accordingly, the phase will be either  $0^\circ$  or  $180^\circ$  and the modal amplitude will be determined solely by the imaginary value. In reality, however, perfect  $0^\circ$  or  $180^\circ$  phase is seldom observed due to the effects of high or nonideal damping, close modal spacing, nonlinear effects, or analysis limitations. Such has been the case in the present study where the primary difficulty has been the tendency for each of the two plates to exhibit slight but noticeably different modal properties. The behavior was not constant and tended to be more pronounced with some modes than others.

Test-Configurations: The dynamic response tests were carried out for the same three different gasket configurations employed for the static tests (the same actual test set-ups were used) and for both a sealed and "vented" interpane void. The test glazings are described in section 4.4 and are referred to as: (i) nominal gasket installation, (ii) 10% less gasket (lineal basis), and (iii) 5% more gasket (lineal basis). Interpane venting involved drilling  $4 \frac{3}{8}$  inch diameter holes (one near each corner) through the interpane spacer. These were either sealed mechanically or

left clear and open as required. For clarity, the results are discussed below in two categories, first the general behavior for a given configuration and second the effects of variations in gasket properties and venting of the interpane void.

Nominal Gasket: Table 4.5-1 shows the modal frequencies and damping for the nominal gasket configuration as obtained using the transfer function curve-fitting algorithms incorporated in the Modal Analysis software on the HP5451B Analyzer. Each mode was identified by the phase relations at the measurement points. The results are shown for both the sealed and vented interpane void tests, and for comparison, the continuous and discrete analytical predictions are also listed. Several points are noteworthy.

(1) - The experimental results identify not only the first 4 "classical" plate modes but also the 1'-2' and possibly the 2'-1' modes.

In the classical modes, the two plates move together and the window behaves like an equivalent single plate, while for the primed modes, the two plates move in opposite directions. As noted in Chapter 3, the latter behavior may or may not involve a net interpane compression or breathing action, depending on the mode shape.

The primed modes above are of the even-odd type and hence do not involve a net void volume change (breathing), but they do depend on the interplate rotational coupling at the edge. This coupling stiffness was not quantitatively assessed (since it would involve destruction of a specimen) but from the continuous model results (section 3.4) the 1'-2' frequency could range from 18.09-22.44 Hz as the rotational stiffness varies from 0 to  $\infty$ . The observed

frequencies indicate that this type of coupling is not particularly strong for the present type of window. (2) - The measured damping is in excess of 5% for most modes which is somewhat large for metal or glass structures alone. This behavior is largely due to the viscoelastic properties of the boundary gasket, especially for the classical plate modes. The damping values for the primed modes are considerably smaller and reflect the reduced participation of the gasket in this type of response. From a design point of view, this suggests that glazing properties, through their viscoelastic behavior, might be used to control peak dynamic response in the classical plate modes. On the other hand, this will likely have little effect on the primed modes which could lead to the possibility of interplate contact, especially for very large panes with strong non-uniform pressure loadings. (3) - The effect of interpane venting is essentially negligible, for the present conditions. This is basically in agreement with the analysis since for the lower modes there should be no net interpane volume change. The effect would be a factor only for the breathing modes, the lowest of which is near 100 Hz, and at that frequency, the present holes would be of insignificant size. In contrast, the static response is highly affected. (4) - While not evident in Table 4.5-1, the presence of other modes at similar frequencies which involve substantial boundary participation was not initially anticipated and proved difficult to interpret. These modes are predicted by the discrete model in section 3.3 but due to test scheduling, appropriate boundary response instrumentation was not installed. (5) - A clear normal

mode response was not observed for some of the modes. This can be attributed in part to the high damping, but was largely due to the tendency of the two plates comprising the window to behave slightly differently for different modes. This effectively resulted in two very closely spaced modes which proved impossible to detect by hand and difficult at best to handle with the modal analysis software.

Boundary Gasket Effects: The effect of variation in the boundary gasket support stiffness is shown in Table 4.5-2. The data cover the first two modes for both vented and sealed interpane conditions. As before, these data were obtained at first in rough form by visual inspection of the transfer functions and later in the present form using the analytical curve-fitting software. Several features are noteworthy. (1) - There is little difference between the vented and sealed interpane void data. As was anticipated and observed in the previous results. For clarity, these two cases have been averaged together and listed in the lower lines as the "Average" modal values. (2) - There is a consistent trend for the frequencies to increase as the boundary support is stiffened. This is seen most clearly in the "average" values and is most pronounced for the higher mode. This type of behavior is fully consistent with the analytical predictions developed in Chapter 3. In the present tests the major effort has been to simulate the field conditions as closely as possible, using for example the actual cladding components (mullions, spandrells, etc.). Unfortunately, this approach while yielding results representative

TABLE 4.5-1

MEASURED MODAL FREQUENCIES AND DAMPING  
FOR NOMINAL GASKET AND SEALED INTERPANE VOID

Source	Parameter	Mode					
		1-1	1-2	1'-2'	1-3	2-1	2'-1'
Measured	Freq. (Hz)	10.2	19.4	19.6	28.6	32.8	42.3
	Damping (%)	6.0	3.3	1.8	10.9	7.6	1.1
Theory	Freq. (Hz)	12.2 <sup>1</sup>	18.1 <sup>2</sup>	18.1 <sup>3</sup>	34.7 <sup>1</sup>	34.2 <sup>2</sup>	36.9
				28.4			52.1

## Notes:

1. Discrete element model with measured properties.
2. Continuum model for simple support.
3. Continuum model for interplate restraint from 0 to infinity  
(these are out-of-phase modes involving opposing movement).

TABLE 4.5-2

## DEPENDENCE OF MODAL VALUES ON GASKET RESTRAINT

Configuration	1-1 MODE		1-2 MODE	
	Frequency (Hz)	Damping (%)	Frequency (Hz)	Damping (%)
Sealed				
-10%	10.18	5.7	17.54	3.5
0%	10.23	6.0	19.31	3.5
+5%	10.01	3.5	19.70	2.6
Vented				
-10%	9.95	5.7	17.49	3.1
0%	10.16	7.6	19.40	-
5%	10.48	7.8	19.67	2.9
Average				
-10%	10.06	5.7	17.52	3.3
0%	10.20	7.8	19.36	3.3
5%	10.25	5.8	19.68	2.8

of field conditions, does not readily allow for quantitative variations in these conditions. (For this purpose a special adjustable test frame would be required). Thus, it has not been possible to quantitatively assess the range of boundary restraint conditions (-10%, + 5% gasket material) used. Rather, the nominal configuration was carefully evaluated in a series of tests described in section 4.4 and two other configurations were then selected to represent the practical extremes of support for the present glazing system.

(3) - The damping is not appreciably affected by alteration of the glazing gasket as presently accomplished. However, as observed earlier, there is a substantial change in the damping from mode to mode.

Summary: The modal tests using either impulse or uniform random pressure have yielded frequencies, damping and mode shapes for the double glazed window when supported in a manner closely duplicating the actual service conditions. The measured frequencies agree reasonably well with analytical predictions, Although sufficient experimental measurements were not made to adequately define those modes involving significant boundary movement. Damping values tended to be rather large for glass or metal structures as a result of the viscoelastic boundary restraint provided by the glazing gaskets. However, variations in the amount of gasket material used in the glazing produced insignificant changes in damping. The most significant result of the tests was the detection of the primed modes in which the two panes move in opposite directions. Since the damping was found to be relatively low in the lab tests,

the possibility exists that these modes could be excited at relatively high amplitudes, especially in large windows where non-uniform pressure loading may occur. Several problems were encountered, however. For certain modes, there was a measurable difference in the dynamics of each plate which appeared in the measured transfer functions as two extremely closely spaced modes and greatly complicated the analysis. Finally, the major effort was made to closely simulate actual conditions and while this yielded results in close agreement with field and analytical studies, ( and provided confidence in the analysis validity), it did not allow enough time to carry out a parametric analysis for design purposes.

## 5.0 Conclusions

The study reported here has attempted to more precisely characterize the nature of the local cladding pressure loads and to identify the degree of causality for the dynamic response of the cladding--in particular, large insulating glass windows. This has involved a combination of field measurements, analytical modelling of the cladding, and laboratory simulation and testing.

The field measurement results are generally in agreement with the limited amount of other reported work. The chosen site exhibited a range of building geometries because of its unique shape, and consequently, it has been possible to measure the behavior of similar cladding components for a variety of flow conditions. Generally it was found that the differential cladding pressure contained relatively little power over frequency ranges encompassing the first few cladding natural frequencies and the causality (or coherence function value) was very low as well. Taken together this implies that substantial dynamic response is not likely due to pressure loading. On the other hand, for certain locations and under specific conditions, relative large pressure power levels were observed over frequency ranges including the cladding fundamental frequency. For these critical locations, a significant causality existed between the input pressure and the dynamic response.

The analytical studies centered around development of a general cladding model and computer program which would permit a wide range of cladding configurations and loading types to be considered. Panel response to static and dynamic pressure loading and racking



distortions introduced by differential motions of adjacent floors, both treated as deterministic functions of time, were the principal loadings considered. Sensitivity studies were conducted to determine the influence of system properties on the response of the panel-frame assembly. Response was shown to be sensitive to gasket rotational stiffness and frame flexural stiffness for selected panel aspect ratios. In general, the finite element model produced frequency and displacement results which were in good agreement with laboratory and field test values. Out-of-phase double panel modes were predicted both at relatively low frequencies and at frequencies beyond the range of practical interest.

Subsequent analytical studies of glass cladding should consider large displacement response of the lights. Viscoelastic properties of the gasket, damping mechanisms, and treatment of wind excitation as a random process are also worthy extensions to this study. The present discrete element model can be modified to include these effects. Finally, the interaction of the entire glass curtain wall with the structural frame should be investigated in future studies. Claddings of various types are likely to provide measurable in-plane stiffness which could be used to advantage by the structural designer to help control interstory drift and building motion affecting occupant comfort.

The laboratory phase of the study centered around development of a full-scale pressure simulation facility and subsequent use of this equipment with cladding components from the field site. The facility has generally met or exceeded its design objectives

and has been found of great value in simulating both steady state and transient test conditions. Extensive laboratory measurements were carried out to precisely define the local structural characteristics of the cladding system under study (which is representative of a number of common designs). Subsequent static and dynamic response measurements produced results that were in good agreement with the predictions of the analytical model and thus confirmed its validity and appropriateness. The dynamic response measurements using a point impulse forced excitation method revealed the presence as predicted, of double plate modes in which the plates move out of phase with respect to each other. One of these, modes involving no net interpane volume change, was found to occur at a low enough frequency (20 Hz) to pose potential problems in large windows where significant pressure non-uniformity (which is required to excite this mode) might occur. This type of motion, if severe enough, could obviously lead to plate-to-plate impact with the risk of consequent failure.

The lab studies were unfortunately somewhat limited in breadth since the primary requirement was to characterize a typical cladding system and validate the analysis for it. Future work with this facility should attempt to apply the model to and determine the characteristics of a variety of cladding systems. This approach will directly furnish design-oriented information so critical for efficient, sound and economical design.

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## APPENDIX

### Listing of FORTRAN Computer Program WINDOW

```
*DECK MAIN
  PROGRAM WINDOW(INPUT,OUTPUT,TEMPS,SFFMFF,REACS,TAPE5=INPUT,
#      TAPE6=OUTPUT,TAPE10=TEMPS,TAPE13=SFFMFF,
#      TAPE11=REACS)
```

C- -- -- PROGRAM WINDOW

C-----PROGRAMMER: DR. BARRY J. GOODNO 1975-6  
C SCHOOL OF CIVIL ENGINEERING  
C GEORGIA INSTITUTE OF TECHNOLOGY  
C ATLANTA, GEORGIA 30332  
C (404) 896-2227

## C-----DESCRIPTION:

C WINDOW - PROGRAM FOR THE LINEAR STATIC AND DYNAMIC ANALYSIS  
OF FRAMED PANELS WITH ARBITRARY SUPPORT CONDITIONS AND  
ARBITRARY SELECTION OF STRUCTURE DEGREES OF FREEDOM.

C IN THIS PROGRAM, PANELS ARE REPRESENTED BY A COMBINATION OF THE  
C BARBER-DAKBERG-WEAVER (BOW) REFINED PLANE STRESS RECTANGULAR  
C FINITE ELEMENT AND THE BOGNER-FOX-SCHMITT (BFS) RECTANGULAR PLATE  
C BENDING ELEMENT. THE EXTERIOR FRAMING (IF PRESENT) CONSISTS OF  
C BEAM AND COLUMN ELEMENTS AND THE BEAM ELEMENTS MAY HAVE NONRIGID  
C CONNECTIONS. BOUNDARY NODES ON THE PANEL AND ADJACENT FRAME NODES  
C MAY BE JOINED BY THREE TRANSLATIONAL AND ONE ROTATIONAL SPRING  
C (PARALLEL TO THE PANEL EDGE); THUS THE MODEL CAN BE USED TO  
C REPRESENT A WINDOW-GASKET-MULLION ASSEMBLAGE. IN ADDITION, LINEAR  
C ELASTIC SPRINGS MAY BE ATTACHED TO ANY PANEL NODE IF DESIRED.

C THE BASIC CHARACTER OF THE MODEL DEPENDS UPON THE VALUE OF THE V  
C E  
C "NFM" (=NUMBER OF BOUNDARY FRAME MEMBERS). IF NFM>0, A PANEL-  
C FRAME ASSEMBLAGE IS IMPLIED; IF NFM=0, AN ARBITRARILY-SUPPORTED  
C PANEL WITH NO FRAME IS IMPLIED. BOTH PANEL AND FRAME DOFS MAY  
C BE RESTRAINED AS DESIRED AND SUPPORT DISPLACEMENTS SPECIFIED AS  
C A LOADING CONDITION. PANEL AND FRAME ARE NUMBERED INDEPENDENTLY

C AN ARBITRARY NUMBER AND DISTRIBUTION OF PANEL NODE DEGREES OF FREEDOM  
C MAY BE DESIGNATED AS THE STRUCTURE DEGREE OF FREEDOM (CALLED F  
C DISPLACEMENT TYPES) AND EXTERNAL PANEL AND MULLION DEGREES OF  
C FREEDOM ARE ELIMINATED THROUGH MATRIX CONDENSATION (GUYAN REDUCTION).  
C AT THE END OF THE ANALYSIS, STRUCTURE STIFFNESS AND MASS MATRICES  
C COUPLING THE DEGREES OF FREEDOM WITH EACH OTHER AND WITH THE SUPPORT  
C RESTRAINTS (R DISPLACEMENT TYPES) ARE LEFT, AND STRUCTURE FREQUENCIES  
C AND MODE SHAPES MAY BE DETERMINED. THE STRUCTURE MASS MATRIX MAY BE  
C FORMED USING CONSISTENT MASS (CM), ASSEMBLED-LUMPED MASS (ALM), OR  
C INPUT-LUMPED MASS (ILM).

C TIME DEPENDENT NODAL LOADS, ACCELERATIONS OF THE RIGID SUPPORT SYSTEM  
C OR ARBITRARY INDEPENDENT DISPLACEMENTS OF SUPPORTS,  
C ALL TREATED AS PIECEWISE-LINEAR FUNCTIONS OF TIME IN THE STRUCTURE X,  
C Y, AND Z DIRECTIONS, ARE ACCEPTABLE LOADING CONDITIONS. THE NODAL  
C MODE METHOD IS USED TO DETERMINE THE DISPLACEMENT RESPONSE OF THE  
C SYSTEM. DISPLACEMENT-TIME HISTORIES ARE PRINTED AND/OR PLOTTED FOR  
C DEGREES OF FREEDOM (F TYPES) SELECTED BY THE ANALYST.  
C IN ADDITION, STATIC DISPLACEMENTS, PANEL STRESSES, AND SUPPORT  
C REACTIONS MAY BE DETERMINED IF STATIC BEHAVIOR IS OF INTEREST  
C (I.E., ISDA=0).  
C

MAI10000  
MAI00010  
MAI00229  
MAI00633  
MAI00743  
MAI00855  
MAI00866  
MAI00873  
MAI00966  
MAI00996  
MAI01010  
MAI01011  
MAI01030  
MAI01039  
MAI01050  
MAI01066  
MAI01070  
MAI01090  
MAI01099  
MAI01200  
MAI01220  
MAI01221  
MAI01230  
MAI01246  
MAI01259  
MAI01266  
MAI01270  
MAI01290  
MAI01296  
LMAI03030  
MAI03032  
MAI03110  
MAI03330  
MAI03346  
MAI03354  
MAI03366  
MAI03376  
MMAI03380  
MAI03396  
MAI03400  
MAI03416  
MAI03426  
MAI03430  
MAI03440  
MAI03459  
MAI03466  
MAI03476  
MAI03490  
MAI03499  
MAI03510  
MAI03522  
MAI03533  
MAI03550  
MAI03566  
MAI03570  
MAI03580  
MAI03590

## C-----SUBPROGRAMS IN PROGRAM WINDWAT

## C 1) PRIMARY SUBPROGRAMS

```

C 1) PRIMARY 308-400R
C DATA - READS AND PRINTS STRUCTURE DATA AND CONTROL PARAMETERS
C PANL5N - FORMS STIFFNESS AND MASS SUBMATRICES FOR PANEL
C FRAM5N - FORMS STIFFNESS AND MASS SUBMATRICES FOR FRAME
C SPRNG - FORMS STIFFNESS SUBMATRICES FOR PANEL AND FRAME DEGREES
C OF FREEDOM COUPLED BY ELASTIC SPRINGS
C CDENSE - CONDENSATION OF STIFFNESS AND MASS MATRICES
C EIGENV - DETERMINES FREQUENCIES AND MODE SHAPES FOR STRUCTURE DOFS
C STATC - COMPUTES STATIC DISPLACEMENTS, STRESSES, REACTIONS
C DYNAMC - FREQUENCIES, MODE SHAPES, AND DYNAMIC RESPONSE
C DYDATA - READS AND PRINTS DATA FOR DYNAMIC ANALYSIS
C RESPDN - COMPUTES AND PRINTS DISPLACEMENT RESPONSE OF STRUCTURE

```

## C 21 SECONDARY SUBPROGRAMS

```

C AXLBX, TREO2, TQL2 - EIGENVALUE PROBLEM SOLUTION ROUTINES (CALLED
C BY EIGENV)
C MATPR - PRINTS MATRICES
C COMPM - CHOLESKY DECOMPOSITION OF A MATRIX (CALLED BY COENSE)
C SOLVE - SOLVES SIMULTANEOUS EQUATIONS (CALLED BY COENSE)
C GRAPM - PRINTER PLOT OF SYSTEM RESPONSE (CALLED BY RESPON)
C TSET - FORMS PERMUTATION VECTOR FOR FORMATION OF ELEMENT
C STIFFNESS AND MASS MATRICES (CALLED BY PANLSM)
C ROALC - DYNAMIC RESPONSE CALCULATIONS (CALLED BY RESPON)
C SBOG - STIFFNESS MATRIX FOR BGS ELEMENT (CALLED BY PANLSM)
C SBOU - ELEMENT STIFFNESS MATRIX FOR REFINED PLANE STRESS ELEMENT
C (CALLED BY PANLSM)
C BOGCM - CONSISTENT MASS MATRIX FOR BGS ELEMENT (CALLED BY PANLSM)
C BOUCCM - CONSISTENT MASS MATRIX FOR BOU ELEMENT (CALLED BY PANLSM)
C BFSSTR - STRESS CALCULATIONS FOR BGS ELEMENT (CALLED BY STATIC)
C BOUSTR - STRESS CALCULATIONS FOR BOU ELEMENT (CALLED BY STATIC)

```

```
C-----MAIN PROGRAM OF WINDOW
```

```
REAL MFF,MFR,L,MP,MF,IX,IY,IZ,KS1,KS2
INTEGER SN,JKJ(3),PPV,FPV,PLIST
```

## COMMON

```

1/CBPNPXP(25),YPI(25),MP(25),KS1(25,5),JFN(25,2),JPN(25),INTYP
# IREP(16),ICNTY(16,A),KS2(25,5)
2/CBNFJFX(I(16),YF(16),MF(16),IX(16),IY(16),IZ(16),A(16),L(16),
# EZJ(16),EZK(16),EYJ(16),EYK(16),IFX(16),JJ(16),JK(16),
3/CBNPFDP/PPV(200),FPV(96),NFL(200),NRL(200),JRL(96)
4/CBSIF/SFF( 40, 40),SFF( 40, 120),HFF( 40, 45),SFR( 40, 56),
# SPP(120, 120),SPM(120,45),HFR(120, 56),
# SHM( 45, 45),SMR( 45, 56),
# SRP( 56, 56)
5/CBSTOR/TS1(120,45),TS2(120,56),TS3( 45, 120)
6/CBMISC/SW(32,32),SMD(120,21),DUM(20)
7/CBOATA/E,P,T,PKP,OENSP,E,F,G,T,O,DENSF,NS,SN,MCOOE,ISM,
# NRPM,NFT,NPT,NMT,NPT,NF,NM,NR,NP,NM,NPE,NPD,NPDOF,
# NFJ,NFM,NJG,NFNSPR,NSTPAT,ISDA,NKPR,NFR,NRF,J,
# NFR,I,REAC,ISTRES,IPSTAT,IL,IKKK,IME
8/CBOYNQ/PLIST(30),INDO,INQ,IAF,IGA,ISD,NOLC,LN,NPTS,NMOOSE,
# IOMAX,IOTHAND,DI_02,FX,FY,FZ,AYO,AZO,AYS,SYS,SZ,

```

MAI106000  
MAI106010  
MAI106200  
MAI106300  
MAI106400  
MAI106500  
MAI106600  
MAI106700  
MAI106800  
MAI106900  
MAI107000  
MAI107100  
MAI107200  
MAI107300  
MAI107500  
MAI107600  
MAI107700  
MAI107800  
MAI107900  
MAI108000  
MAI108100  
MAI108200  
MAI108300  
MAI108400  
MAI108500  
MAI108600  
MAI108700  
MAI108800  
MAI108900  
MAI109000  
MAI109100  
MAI109200  
MAI109300  
MAI109400  
MAI109500  
MAI109600  
MAI109700  
MAI109800  
MAI109900  
MAI110000  
MAI110100  
MAI110200  
MAI110300  
MAI110400  
MAI110500  
MAI110600  
MAI110700  
MAI110800  
MAI110900  
MAI111000  
MAI111100  
MAI111200  
MAI111300  
MAI111400  
MAI111500  
MAI111600  
MAI111700  
MAI111800  
MAI111900

```

      #          SXA,SYA,SZA,NTX,NTY,NTZ
C
C-----ARRAY DIMENSIONS IN COMMON (TS1: NPXNM: TS2: NPXNR: TS3: NMXNP)
C
      NF=40
      NP=120
      NM=45
      NR=56
C
C-----READ IN NUMBER OF STRUCTURES TO BE PROCESSED (NS):
C      ALSO, READ IN PARAMETER 'ISDA' (ISDA=0, STATIC ANALYSIS; ISDA=1,
C      DYNAMIC ANALYSIS)
C
      READ(5,*) NS
C
      SN=0
      1 SN=SN+1
C
C-----READ IN STRUCTURE DATA FOR STRUCTURE NUMBER SN
C
      CALL SDATA
C
C-----GENERATE STIFFNESS (ISM=1) AND MASS (ISM=2) MATRICES FOR STRUCTURE
C
      ISM=0
      2 ISM=ISM+1
C
      IF((ISDA.EQ.0).AND.(ISM.EQ.2))GO TO 10
      IF((ISM.EQ.2).AND.(MCODE.EQ.3))GO TO 3
C
C-----FORM PANEL STIFFNESS (ISM=1) AND MASS (ISM=2) SUBMATRICES
C
      CALL PANLSM
C
C-----FORM FRAME STIFFNESS (ISM=1) AND MASS (ISM=2) SUBMATRICES (UNLESS
C      NFM=0)
C
      IF(NFM.EQ.0)GO TO 4
C
      CALL FRAMSM
C
C-----ADD BOUNDARY SPRING CONSTANTS TO STIFFNESS SUBMATRICES
C
      4 IF((ISM.EQ.2).OR.((NFMSPR+MPNSPR).EQ.0))GO TO 5
C
      CALL SPRING
C
      5 CONTINUE
C
C-----CONDENSE STIFFNESS (ISM=1) AND MASS (ISM=2) MATRICES BY ELIMINATION
C      G
C
      P (NFM=0), OR P AND M (NFM>0), DISPLACEMENT TYPES
C
      CALL CDENSE
C
      GO TO 7
C
C-----FORM MASS SUBMATRICES FOR MCODE=3 CASE
C
      3 DO 102 I=1,NFT

```

```

      MAI01200
      MAI01210
      MAI01220
      MAI01230
      MAI01240
      MAI01250
      MAI01260
      MAI01270
      MAI01280
      MAI01290
      MAI01300
      MAI01310
      MAI01320
      MAI01330
      MAI01340
      MAI01350
      MAI01360
      MAI01370
      MAI01380
      MAI01390
      MAI01400
      MAI01410
      MAI01420
      MAI01430
      MAI01440
      MAI01450
      MAI01460
      MAI01470
      MAI01480
      MAI01490
      MAI01500
      MAI01510
      MAI01520
      MAI01530
      MAI01540
      MAI01550
      MAI01560
      MAI01570
      MAI01580
      MAI01590
      MAI01600
      MAI01610
      MAI01620
      MAI01630
      MAI01640
      MAI01650
      MAI01660
      MAI01670
      MAI01680
      MAI01690
      MAI01700
      MAI01710
      MAI01720
      MAI01730
      MAI01740
      MAI01750
      MAI01760
      MAI01770
      MAI01780
      MAI01790
      DO 103 J=1,NFT
      103 MFF(I,J)=J.
      DO 102 K=1,NRT
      102 MFR(I,K)=0.
C
      DO 104 I=1,NPN
      JKJ(1)=0*I-7
      JKJ(2)=JKJ(1)+1
      JKJ(3)=JKJ(1)+2
      DO 105 K=1,3
      J=PPV(JKJ(K))
      IF(J,GT,NFT)GO TO 105
      MFF(J,J)=MFF(J,J)+P(I)
      105 CONTINUE
      104 CONTINUE
      WRITE(33) ((MFF(I,J),J=1,NFT),I=1,NFT)
      WRITE(33) ((MFR(I,J),J=1,NRT),I=1,NFT)
      GO TO 9
C
C-----SAVE CONDENSED MATRICES SFF AND SFR (ISM=1) AND MFF AND MFR (ISM=2)
C      AND MCODE,NE.3) ON MASS STORAGE
C
      7 IF(ISM.EQ.1)REWIND 33
      WRITE(33) ((SFF(I,J),J=1,NFT),I=1,NFT)
      WRITE(33) ((SFR(I,J),J=1,NRT),I=1,NFT)
C
      IF(ISM.LT.2)GO TO 2
C
      DO 100 I=1,NFT
      DO 101 J=1,NFT
      101 MFF(I,J)=SFF(I,J)
      DO 100 K=1,NRT
      100 MFR(I,J)=SFR(I,J)
C
      REWIND 33
      READ(33) ((SFF(I,J),J=1,NFT),I=1,NFT)
      READ(33) ((SFR(I,J),J=1,NRT),I=1,NFT)
C
C-----DETERMINE FREQUENCIES, MODE SHAPES, AND DYNAMIC RESPONSE
C
      9 CALL DYNAMIC
      GO TO 11
C
C-----STATIC ANALYSIS
C
      10 CALL STATIC
C
C-----TEST FOR FINAL STRUCTURE
C
      11 IF(SN.LT.NS)GO TO 1
C
      STOP
      END
      MAI01800
      MAI01810
      MAI01820
      MAI01830
      MAI01840
      MAI01850
      MAI01860
      MAI01870
      MAI01880
      MAI01890
      MAI01900
      MAI01910
      MAI01920
      MAI01930
      MAI01940
      MAI01950
      MAI01960
      MAI01970
      MAI01980
      MAI01990
      MAI02000
      MAI02010
      MAI02020
      MAI02030
      MAI02040
      MAI02050
      MAI02060
      MAI02070
      MAI02080
      MAI02090
      MAI02100
      MAI02110
      MAI02120
      MAI02130
      MAI02140
      MAI02150
      MAI02160
      MAI02170
      MAI02180
      MAI02190
      MAI02200
      MAI02210
      MAI02220
      MAI02230
      MAI02240
      MAI02250
      MAI02260
      MAI02270
      MAI02280
      MAI02290
      MAI02300
      MAI02310
      MAI02320

```

\*DECK SDATA  
SUBROUTINE SDATA

C-----SUBPROGRAM TO READ IN STRUCTURE DATA FOR PANEL, FRAME, AND  
C BOUNDARY SPRINGS; CONTROL PARAMETERS, MATERIAL PROPERTIES,  
C NODAL COORDINATES, MASS DATA, AND CONNECTIVITY AND RESTRAINT  
C INFORMATION.

C REAL IX,IY,IZ,L,M,FP,MP,KS1,KS2,TITLE(8),UNITS(8)  
C INTEGER SM,PPV,PPV

C COMMON  
C 1/CBNPN/XP(25),YP(25),MP(25),KS1(25,5),JFN(25,2),JPN(25),INTYP(25),  
C IREP(16),ICNTY(16,4),KS2(25,5)  
C 2/CBNFJ/XP(16),YP(16),MF(16),IX(16),IY(16),IZ(16),A(16),L(16),  
C EZJ(16),EZK(16),EYJ(16),EYK(16),IFM(16),JJ(16),JK(16)  
C 3/CBNPFD/PPV(200),FPV(96),NFI(200),NRL(200),JRL(96)  
C 7/CBQDATA/EP,TP,PRP,DENSP,EP,GF,TF,DENSP,NS,SN,MCODE,ISM,  
C NRPN,NFT,NPT,NMT,NRT,NF,NP,NM,NR,NPN,NPE,NPD,NPDF,  
C NFJ,NFM,NJ6,NFNSPR,NPNSPR,ISDA,NPR,NFR,NFJ,  
C NFRT,IREAC,ISTRES,ISTAT,IL,IKKK,TIME

C 99 FORMAT(\*1 STATIC/DYNAMIC ANALYSIS OF AN ARBITRARILY-SUPPORTED\*,  
C \* PLATE STRUCTURE\*,//)

C 100 FORMAT(\*1 STATIC/DYNAMIC ANALYSIS OF FRAMED PANELS\*,//)  
C READ(5,101) TITLE  
C READ(5,101) UNITS

C 101 FORMAT(8A10)  
C READ(5,\*) NPE,NFM,ISDA  
C NFR=NPR=NRT=0

C-----I. PANEL DATA: CONTROL PARAMETERS, MATERIAL PROPERTIES, NODAL  
C COORDINATES, LUMPED MASSES, CONNECTIVITY  
C DATA, AND RESTRAINT INFORMATION (IF NRPN > 0)  
C NOTE: PANEL NODES ARE NUMBERED 1,...,NPN

C-----A. CONTROL PARAMETERS, MATERIAL PROPERTIES

C 103 FORMAT(\*0PANEL DATA\*,//, \* NPE NPN NRPN\*,  
C \* NMFT EP PRP DENSP TP\*)

C-----MCODE: 1=CHI+ILM; 2=ALM+ILM; 3=ILM ONLY

C READ(5,\*) NPN,NRPN,NPR,NMFT,NFT,EP,PRP,DENSP,TP,MCODE  
C IF(NFM.EQ.0)PRINT 99  
C IF(NFM.GT.0)PRINT 100  
C PRINT 102,TITLE,UNITS

C 102 FORMAT(\*0TITLE: \*,8A10,/, \*0UNITS: \*,8A10)  
C PRINT 500,SN,NS,ISDA

C 500 FORMAT(/, \*0STRUCTURE NUMBER\*, I3, \* DF\*, I3, /,  
C \* ISDA =\*, I2, /)

C PRINT 103  
C PRINT 104,NPE,NPN,NRPN,NMFT,EP,PRP,DENSP,TP

C 104 FORMAT(I4,2I5,I7,1P4E13.5)  
C NPD=0\*NPN

C NPDF=NPD-NPR  
C NPT=NPDF-NFT  
C NRT=NPR

C PRINT 105  
C 105 FORMAT(\*0 NPD NPDF NFT NPT NPR MCODE\*)

C PRINT 106,NPD,NPDF,NFT,NPT,NPR,MCODE  
C 106 FORMAT(I5,2I6,3I5)

SDAL0000  
SDAL0010  
SDAG0020

SDAL0030  
SDAL0040  
SDAL0050

SDAL0060  
SDAL0070  
SDAL0080

SDAL0090  
SDAL0100  
SDAL0110

SDAL0120  
SDAL0130  
SDAL0140

SDAL0180  
SDAL0190  
SDAL0200

SDAL0210  
SDAL0220  
SDAL0230

SDAL0240  
SDAL0250  
SDAL0260

SDAL0270  
SDAL0280  
SDAL0290

SDAL0300  
SDAL0310  
SDAL0320

SDAL0330  
SDAL0340  
SDAL0350

SDAL0360  
SDAL0370  
SDAL0380

SDAL0390  
SDAL0400  
SDAL0410

SDAL0420  
SDAL0430  
SDAL0440

SDAL0450  
SDAL0460  
SDAL0470

SDAL0480  
SDAL0490  
SDAL0500

SDAL0510  
SDAL0520  
SDAL0530

SDAL0540  
SDAL0550  
SDAL0560

SDAL0570  
SDAL0580  
SDAL0590

SDAL0590

DO 1061 I=1,NPD  
1061 NFL(I)=NRL(I)=0

C-----B. NODAL COORDINATES; LUMPED MASSES

C PRINT 107  
C 107 FORMAT(\*0NODAL COORDINATES\*,//, \* MODE X COORD. Y COORD. \*,  
C \* \* ADOEO MASS\*)  
C DO 108 I=1,NPN  
C READ(5,\*) N,XP(N),YP(N),MP(N)  
C 108 PRINT 109,N,XP(N),YP(N),MP(N)  
C 109 FORMAT(I4,1P3E13.5)

C-----C. CONNECTIVITY, DEGREE-OF-FREEDOM (F TYPES), AND RESTRAINT DATA

C PRINT 110  
C 110 FORMAT(\*0ELEMENT CONNECTIVITY DATA\*,//, \* ELEMENT NODE \*,  
C \* NUMBERS REPEITION\*,//, I4X, \* I J K L CODE\*)

C C-----ICNTY(K,J) = CGW NODE NOS. FOR ELEMENT #K; J=1 TO 4 FOR RECT.  
C ELEMENT; START NODE NOS. AT LOWER LEFT NODE

C C-----IREP(K) = REPEITION CODE FOR ELEMENT #K (0 = DIFFERENT;  
C 1 = SAME MATERIAL AND DIMENSIONS AS ELEMENT #K-1))

C DO 111 I=1,NPE  
C READ(5,\*) K,(ICNTY(K,J),J=1,4),IREP(K)

C 111 PRINT 112,K,(ICNTY(K,J),J=1,4),IREP(K)  
C 112 FORMAT(I6,I9,3I4,I12)

C C-----READ IN NODAL F-TYPE LIST (NFL) FOR NODE #J (1=F TYPE, 0=OTHER)

C PRINT 113  
C 113 FORMAT(\*0NODAL F DISPL. TYPES\*,//, \* NODE J1 J2 J3 J4 J5\*,  
C \* \* J6 J7 J8\*)

C DO 114 K=1,NMFT  
C READ(5,\*) J,NFL(8\*J-7),NFL(8\*J-6),NFL(8\*J-5),NFL(8\*J-4),  
C NFL(8\*J-3),NFL(8\*J-2),NFL(8\*J-1),NFL(8\*J)

C J1=8\*J-7  
C J8=8\*J

C 114 PRINT 115,J,(NFL(I),I=J1,J8)  
C 115 FORMAT(I4,I5,7I4)

C C-----READ IN RESTRAINT LIST (NRL) FOR PANEL (IF NRPN=0)  
C (NRL: 0 = NO RESTRAINT, 1 = RESTRAINT EXISTS)

C OTHERWISE (IF NRPN = 0) READ IN BOUNDARY FRAME DATA  
C IF(NRPN.EQ.0) GO TO 200

C PRINT 116  
C 116 FORMAT(\*0NODAL RESTRAINT LIST\*,//, \* NODE J1 J2 J3 J4 J5\*,  
C \* \* J6 J7 J8\*)

C DO 117 K=1,NRPN  
C READ(5,\*) J,NRL(8\*J-7),NRL(8\*J-6),NRL(8\*J-5),NRL(8\*J-4),  
C NRL(8\*J-3),NRL(8\*J-2),NRL(8\*J-1),NRL(8\*J)

C J1=8\*J-7  
C J8=8\*J

C 117 PRINT 115,J,(NRL(I),I=J1,J8)  
C IF(NFM.EQ.0) GO TO 300

C C---II. FRAME DATA: CONTROL PARAMETERS, JOINT COORDINATES, MEMBER  
C AND MATERIAL PROPERTIES, RESTRAINT INFORMATION, AND

C ELASTIC CONNECTION DATA  
C NOTE: FRAME JOINTS ARE NUMBERED 1,...,NFJ (1.E., NUMBER FRAME  
C JOINTS INDEPENDENTLY OF PANEL NODES)

C SDAL0600  
SDAL0610  
SDAL0620  
SDAL0630  
SDAL0640  
SDAL0650  
SDAL0660  
SDAL0670  
SDAL0680  
SDAL0690  
SDAL0700  
SDAL0710  
SDAL0720  
SDAL0730  
SDAL0740  
SDAL0750  
SDAL0760  
SDAL0770  
SDAL0780  
SDAL0790  
SDAL0800  
SDAL0810  
SDAL0820  
SDAL0830  
SDAL0840  
SDAL0850  
SDAL0860  
SDAL0870  
SDAL0880  
SDAL0890  
SDAL0900  
SDAL0910  
SDAL0920  
SDAL0930  
SDAL0940  
SDAL0950  
SDAL0960  
SDAL0970  
SDAL0980  
SDAL0990  
SDA10000  
SDA10100  
SDA10200  
SDA10300  
SDA10400  
SDA10500  
SDA10600  
SDA10700  
SDA10800  
SDA10900  
SDA11000  
SDA11100  
SDA11200  
SDA11300  
SDA11400  
SDA11500  
SDA11600  
SDA11700  
SDA11800  
SDA11900

```

C-----A. CONTROL PARAMETERS
C
200 PRINT 201
201 FORMAT('0FRAME DATA',/,*,* -----',/,*,* NFM NFF NFR NRFJ*,
  # * NMEC EF PRF DENSF*)
  READ(5,*) NFM,NFR,NRFJ,NMEC,EF,PRF,DENSF
  PRINT 202,NFM,NFJ,NFR,NRFJ,NMEC,EF,PRF,DENSF
202 FORMAT('14,3I5,16,1P3E13.5)
  GF=EF/(2*(1+PRF))
  NRT=NRT+NFR
  NJ6=6*NFM
  NMT=NJ6-NFR
C
C-----B. JOINT COORDINATES, MEMBER AND MATERIAL PROPERTIES, LUMPED MASS
C
  PRINT 206
206 FORMAT('0JOINT COORDINATES',/,*,* JOINT X COORD. Y COORD.*,
  # * ADDED MASS*)
  DO 207 I=1,NFM
  READ(5,*) J,XF(J),YF(J),MF(J)
207 PRINT 208,J,XF(J),YF(J),MF(J)
208 FORMAT('15,1P3E13.5)
  PRINT 203
203 FORMAT('0MEMBER PROPERTIES',/,*,* MEMBER JJ JK*,9X,*,*,7X,
  # * LENGTH*,11X,*,IX*,11X,*,IY*,11X,*,IZ TYPE*)
  DO 204 II=1,NFM
C-----IFM(II) = INDICATOR FOR FRAME MEMBER TYPE (0=COLUMN, 1=BEAM)
  READ(5,*) I,JJ(II),JK(II),A(II),IX(II),IY(II),IZ(II),IFM(II)
  IF(IFM(II).EQ.0) GO TO 220
  TYPE=6HBEAM
  L(II)=ABS(XF(JK(II))-XF(JJ(II)))
  GO TO 204
  220 TYPE=6HCOLUMN
  L(II)=ABS(YF(JK(II))-YF(JJ(II)))
204 PRINT 205,I,JJ(II),JK(II),A(II),L(II),IX(II),IY(II),IZ(II),TYPE
205 FORMAT('15,16,14,1P5E13.5,2X,A6)
C
C-----C. RESTRAINT INFORMATION (1 = RESTRAINT EXISTS, 0 = NO RESTRAINT)
C
  PRINT 209
209 FORMAT('0JOINT RESTRAINT LIST',/,*,* JOINT J1 J2 J3 J4*,
  # * J5 J6*)
  DO 210 I=1,NJ6
  JRL(II)=0
210 PRINT 211,I,NRFJ
  READ(5,*) J,JRL(6*J-5),JRL(6*J-4),JRL(6*J-3),JRL(6*J-2),
  # JRL(6*J-1),JRL(6*J)
  J1=6*J-5
  J6=6*J
211 PRINT 212,J,(JRL(K),K=J1,J6)
212 FORMAT('2I5,5I4)
C
C-----D. ELASTIC CONNECTION DATA (FOR BEAM MEMBERS ONLY)
C
  DO 216 I=1,NFM
216 EZK(II)=EZK(II)+EYJ(II)+EYK(II)=0.
  IF(NMC.EQ.0) GO TO 303
  PRINT 215
215 FORMAT('0ELASTIC CONNECTION DATA',/,*,* MEMBER*,8X,*,EZJ*,10X,
  SDA11200
  SDA11210
  SDA11220
  SDA11230
  SDA11240
  SDA11250
  SDA11260
  SDA11270
  SDA11280
  SDA11290
  SDA11300
  SDA11310
  SDA11320
  SDA11330
  SDA11340
  SDA11350
  SDA11360
  SDA11370
  SDA11380
  SDA11390
  SDA11400
  SDA11410
  SDA11420
  SDA11430
  SDA11440
  SDA11450
  SDA11460
  SDA11470
  SDA11480
  SDA11490
  SDA11500
  SDA11510
  SDA11520
  SDA11530
  SDA11540
  SDA11550
  SDA11560
  SDA11570
  SDA11580
  SDA11590
  SDA11600
  SDA11610
  SDA11620
  SDA11630
  SDA11640
  SDA11650
  SDA11660
  SDA11670
  SDA11680
  SDA11690
  SDA11700
  SDA11710
  SDA11720
  SDA11730
  SDA11740
  SDA11750
  SDA11760
  SDA11770
  SDA11780
  SDA11790
  # *EZK*,10X,*,EYJ*,10X,*,EYK*)
  DO 217 II=1,NMEC
  READ(5,*) I,EZJ(II),EZK(II),EYJ(II),EYK(II)
217 PRINT 218,I,EZJ(II),EZK(II),EYJ(II),EYK(II)
218 FORMAT('15,2X,1P4E13.5)
C
C---III. ELASTIC SPRING DATA (NFNSPR = NO. OF FRAME NODES WITH SPRINGS
  ATTACHED; NPNSPR = NO. OF PANEL NODES WITH SPRINGS ATTACHED)
C
  300 READ(5,*) NFNSPR,NPNSPR
  DO 320 I=1,NPN
  DO 320 J=1,5
  320 KS1(II,J)=KS2(II,J)=0.
  PRINT 301,NFNSPR,NPNSPR
  301 FORMAT('0SPRING DATA',/,*,* -----',/,*,* NFNSPR =',I3,
  # * NPNSPR =',I3,/)
  IF(NFNSPR.EQ.0) GO TO 310
C-----A. FRAME-PANEL BOUNDARY SPRINGS
C
C-----THE FOLLOWING SECTION PERTAINS TO THE 12 ELEMENT 1/4 MODEL ONLY
C-----AND INVOLVES THE DETERMINATION OF KS1,KS2,... BY FORMULAS
C
  READ(5,*) SKP0,SKL,SKN,SKR
  PRINT 400,SKP0,SKL,SKN,SKR
  400 FORMAT(/,*,* BASIC BOUNDARY SPRING CONSTANTS',/,*,* PULL-OUT =',
  # F0,2,*,* LONGITUDINAL =',F0,2,/,*,* ROTATIONAL =',F0,2,
  # *,* NORMAL =',F0,2)
  AAS=XF(5)
  BBS=YF(5)
C-----FIND KS1,KS2,KS3,KS4,KS5 FOR BOUNDARY NODES
  DO 410 I=1,8
C-----I=FRAME NODE NO.
  JFN(I,1)=I
  IF(I.LE.5)JFN(I,2)=4*I
  IF(I.GT.5)JFN(I,2)=25-I
  GO TO (401,402,402,402,405,406,406,408),I
  401 KS1(I,1)=SKP0*BBS/8.
  KS1(I,2)=SKL*BBS/8.
  KS1(I,3)=SKN*BBS/8.
  KS1(I,4)=2.5E7
  KS1(I,5)=SKR*BBS/8.
  GO TO 410
  402 KS1(I,1)=SKP0*BBS/4.
  KS1(I,2)=SKL*BBS/4.
  KS1(I,3)=SKN*BBS/4.
  KS1(I,4)=5.E7
  KS1(I,5)=SKR*BBS/4.
  GO TO 410
  405 KS1(I,1)=(SKP0*BBS/8.)+(SKL*AAS/6.)
  KS1(I,2)=(SKP0*AAS/6.)+(SKL*BBS/8.)
  KS1(I,3)=SKN*(BBS/8.)+(AAS/6.)
  KS1(I,4)=SKR*(AAS/6.)
  KS1(I,5)=SKR*BBS/8.
  GO TO 410
  406 KS1(I,1)=SKL*AAS/3.
  KS1(I,2)=SKP0*AAS/3.
  KS1(I,3)=SKN*AAS/3.
  KS1(I,4)=SKR*AAS/3.
  KS1(I,5)=5.E7
  GO TO 410
  SDA11800
  SDA11810
  SDA11820
  SDA11830
  SDA11840
  SDA11850
  SDA11860
  SDA11870
  SDA11880
  SDA11890
  SDA11900
  SDA11910
  SDA11920
  SDA11930
  SDA11940
  SDA11950
  SDA11960
  SDA11970
  SDA11980
  SDA11990
  SDA12000
  SDA12010
  SDA12020
  SDA12030
  SDA12040
  SDA12050
  SDA12060
  SDA12070
  SDA12080
  SDA12090
  SDA12100
  SDA12110
  SDA12120
  SDA12130
  SDA12140
  SDA12150
  SDA12160
  SDA12170
  SDA12180
  SDA12190
  SDA12200
  SDA12210
  SDA12220
  SDA12230
  SDA12240
  SDA12250
  SDA12260
  SDA12270
  SDA12280
  SDA12290
  SDA12300
  SDA12310
  SDA12320
  SDA12330
  SDA12340
  SDA12350
  SDA12360
  SDA12370
  SDA12380
  SDA12390

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408 KS1(I,1)=SKL*AAS/6.
      KS1(I,2)=SKPO*AAS/6.
      KS1(I,3)=SKN*AAS/6.
      KS1(I,4)=SKR*AAS/6.
      KS1(I,5)=2.5E7
410 CONTINUE
C
      PRINT 302
302 FORMAT(10NODS AND SPRING CONSTANTS*,/,* FRAME NODE PANEL NODE*,
      / 0X,*KS1*,10X,*KS2*,10X,*KS3*,10X,*KS4*,10X,*KS5*)
C-----KS1(I,J) = SPRING CONSTANT FOR ELASTIC SPRING CONNECTING NODES
C      JFN(I,1) AND JFN(I,2) (I.E., FRAME AND CORRESPONDING PANEL NODE)
C      AND CORRESPONDING TO NODAL DOF J
      DO 303 I=1,NFNSPR
303 PRINT 304,JFN(I,1),JFN(I,2),(KS1(I,J),J=1,5)
304 FORMAI(I6,I13,4X,1P5E13,5)
C-----0. PANEL NODE SPRINGS
310 IF(NPNSPR.EQ.0)RETURN
      PRINT 311
311 FORMAT(* PANEL NODE*,0X,*KS1*,10X,*KS2*,10X,*KS3*,10X,*KS4*,
      / 0X,*KS5*)
C-----KS2(I,J) = SPRING CONSTIANT FOR SPRING ATTACHED TO PANEL NODE
C      JPN(I) AND CORRESPONDING TO NODAL DOF J
      DO 312 I=1,NPNSPR
312 PRINT 313,JPN(I),(KS2(I,J),J=1,5)
313 FORMAT(I7,4X,1P5E13,5)
C
      RETURN
      END

```

```

SOAL2405
SOAL2410
SOAL2420
SOAL2430
SOAL2440
SOAL2450
SOAL2460
SOAL2470
SOAL2480
SOAL2490
SOAL2500
SOAL2510
SOAL2520
SOAL2530
SOAL2540
SOAL2550
SOAL2560
SOAL2570
SOAL2580
SOAL2590
SOAL2600
SOAL2610
SOAL2620
SOAL2630
SOAL2640
SOAL2650
SOAL2660
SOAL2670
SOAL2680
SOAL2690

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```

*DECK PANLSM
      SUBROUTINE PANLSM
C
C-----SUBPROGRAM TO FORM PANEL STIFFNESS (ISM=1) AND MASS (ISM=2 AND
C      MCODE DOES NOT EQUAL 3) SUPMATRICES SFF, SFP, SFR, SPP, SPO,
C      SRR (OR MASS ARRAYS MFF,...,MFR)
C
      REAL MP,KS1,KS2,SB(144)
      INTEGER SN,PPV,FPV,EPV(32),KL(16)
      COMMON
      1/CBNPN/XP(25),YP(25),MP(25),KS1(25,5),JFN(25,2),JPN(25),INTYP(25),
      0 IREP(16),ICNTY(16,4),KS2(25,5)
      3/CBNPFD/PPV(200),FPV(96),NFL(200),NRL(200),JFL(96)
      4/CBSTIF/SFF( 40, 40),SFP( 40, 120),SFM( 40, 45),SFR( 40, 56),
      0 SPP(120,120),SPM(120,45),SPR(120, 56),
      0 SMM( 45, 45),SMR( 45, 56),
      0 SRR( 56, 56)
      6/CBMISC/SH(32,32),SMD(12,12),DUM(20)
      7/CBDATA/EP,IP,PRP,DENSP,EF,GF,TF,DENSF,NS,SN,MCODE,ISM,
      0 NRP,NFT,NPT,NMT,NRT,NF,NP,NM,NR,NPN,NPE,NPD,NPOOF,
      0 NFJ,NFM,NJ6,NFNSFR,NPNSPR,ISOA,NPR,NFR,NRFJ,
      0 NFRI,IREAC,ISTRES,ISTAT,IL,IKKK,TIME
C
      EQUIVALENCE (SB,SPD)
C
C-----CLEAR ARRAYS
C
      DO 192 I=1,NFT
      DO 193 J1=1,NFT
193 SFF(I,J1)=0.
      DO 194 J2=1,NPT
194 SFP(I,J2)=0.
      DO 195 J3=1,NRT
195 SFR(I,J3)=0.
192 CONTINUE
      DO 196 I=1,NPT
      DO 197 J1=1,NPT
197 SPP(I,J1)=0.
      DO 198 J2=1,NRT
198 SPR(I,J2)=0.
196 CONTINUE
      DO 199 I=1,NRT
      DO 199 J=1,NRT
199 SRR(I,J)=0.
C
C-----CONSTRUCT THE PANEL PERMUTATION VECTOR (PPV(I))
C
      IF (ISM.GE.2)GO TO 103
      I1=0
      I2=NFT
      I3=NPOOF
      DO 100 I=1,NPD
      IF(NRL(I),EQ.1)GO TO 101
      IF(NFL(I),EQ.1)GO TO 102
C-----P TYPE DEGREE OF FREEDOM (DOF)
      I2=I2+1
      PPV(I)=I2
      GO TO 100
C-----F TYPE DOF
      102 I1=I1+1

```

```

PAN0000
PAN0001
PAN0002
PAN0003
PAN0004
PAN0005
PAN0006
PAN0007
PAN0008
PAN0009
PAN0010
PAN0011
PAN0012
PAN0013
PAN0014
PAN0015
PAN0016
PAN0017
PAN0018
PAN0019
PAN0020
PAN0021
PAN0022
PAN0023
PAN0024
PAN0025
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PAN0028
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PAN0032
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PAN0036
PAN0037
PAN0038
PAN0039
PAN0040
PAN0041
PAN0042
PAN0043
PAN0044
PAN0045
PAN0046
PAN0047
PAN0048
PAN0049
PAN0050
PAN0051
PAN0052
PAN0053
PAN0054
PAN0055
PAN0056
PAN0057
PAN0058
PAN0059

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```

      PPV(I)=I1
      GO TO 100
C-----R TYPE DOF
      101 I3=I3+1
      PPV(I)=I3
      100 CONTINUE
C
C-----NOW CYCLE THROUGH PANEL ELEMENTS TO FORM STIFFNESS (OR MASS FOR
      ISM=2) SUBMATRICES (SFF, ETC)
C
      103 IREP(I)=0
      DO 110 IE=1,NPE
C
C-----SKIP GENERATION OF SE AND SW IF CURRENT ELEMENT (IE) IS IDENTICAL
      (IN SIZE AND MATERIAL) TO # (IF-1)
C
      IF(IREP(IE).EQ.1) GO TO 111
      DO 191 I=1,32
      DO 191 J=1,32
      191 SW(I,J)=0.
C
C-----FORM PERMUTATION VECTOR FOR ORDERING OF BOGNER AND R-D-W ELEMENTS
      DOF INDICES
C
      CALL TSET(EPV)
C
      J1=ICNTY(IE,1)
      J2=ICNTY(IE,2)
      J3=ICNTY(IE,3)
      A1=XP(J2)-XP(J1)
      B1=YF(J3)-YF(J2)
C-----GENERATE ELEMENT STIFFNESS (OR MASS) ARRAY FOR BOGNER ELEMENT
      IF(ISM.EQ.1) GO TO 112
      IF(MCODE.EQ.2) GO TO 113
C
      CALL BOGCM(DENSP,A1,B1,TP,SB)
C
      GO TO 114
      113 XM=DENSP*A1*B1*TP/4.
      SW(3,3)=SW(11,11)=SW(19,19)=SW(27,27)=XM
      GO TO 150
C
      112 CALL SBOG(A1,B1,TP,EP,PRF,SB)
C
      114 KL(I)=1
      DO 115 I=2,16
      115 KL(I)=KL(I-1)+10-I
      DO 116 I=1,16
      KL=KL(I)
      DO 116 J=1,16
      116 SW(EPV(I),EPV(J))=SW(EPV(J),EPV(I))=SB(KLT-I+J)
C-----GENERATE ELEMENT STIFFNESS (OR MASS) ARRAY FOR REFINED PLANE STRESS
      ELEMENT
      150 IF(ISM.EQ.1) GO TO 149
      IF(MCODE.EQ.2) GO TO 151
C
      CALL BOMCM(DENSP,A1,B1,TP,SB)

```

```

PAM00600
PAN:0610
PAN:0620
PANC0630
PANC0640
PANC0650
PANC0660
PANC0670
PANC0680
PANC0690
PANC0700
PANC0710
PANC0720
PANC0730
PANC0740
PANC0750
PANC0760
PANC0770
PANC0780
PANC0790
PANC0800
PANC0810
PANC0820
PANC0830
PANC0840
PANC0850
PANC0860
PANC0870
PANC0880
PANC0890
PANC0900
PANC0910
PANC0920
PANC0930
PANC0940
PANC0950
PANC0960
PANC0970
PANC0980
PANC0990
PANC1000
PANC1010
PANC1020
PANC1030
PANC1040
PANC1050
PANC1060
PANC1070
PANC1080
PANC1090
PANC1100
PANC1110
PANC1120
PANC1130
PANC1140
PANC1150
PANC1160
PANC1170
PANC1180
PANC1190

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```

C
      GO TO 152
      151 SW(1,1)=SW(2,2)=SW(9,9)=SW(10,10)=XM
      SW(17,17)=SW(18,18)=SW(25,25)=SW(26,26)=XM
      GO TO 111
C
      149 CALL S90M(A1,B1,TP,EP,PRF,SB)
C
      152 DO 153 I=1,16
      KLT=KL(I)
      II=I+16
      DO 153 J=1,16
      JJ=J+16
      153 SW(EPV(II),EPV(JJ))=SW(EPV(JJ),EPV(II))=SB(KLT-I+J)
C
C-----FORM THE ELEMENT PERMUTATION VECTOR (EPV) FROM THE PANEL PERMUTATION
      VECTOR (PPV) FOR ELEMENT NUMBER IE
C
      111 JJ=0
      DO 160 I=1,4
      JP=ICNTY(IE,I)
      JP=8*JP-8
      DO 160 J=1,8
      JP=JP+1
      JJ=JJ+1
      160 EPV(JJ)=PPV(JP)
C
C-----USE THE EPV(I) TO ADD CONTRIBUTION OF CURRENT PANEL ELEMENT (IE)
      TO ARRAYS SFF, SFP, SFR, SPP, SPR, SRF (OR MFF,...MRF FOR ISM=2
      AND MCODE NOT EQUAL 3) NOTE THAT ARRAYS SFR, SPR, AND SRR DO NOT
      RECEIVE ANY PANEL CONTRIBUTIONS WHEN NRPN=0
C
      DO 170 I=1,32
      KR=EPV(I)
      IF(KR.LE.NFT) GO TO 171
      IF((KR.GT.NFT).AND.(KR.LE.NPOOF)) GO TO 172
C-----R DISPLACEMENT TYPES (NO R TYPES ON PANEL FOR NRPN=0)
      KR=KR-NPOOF
      DO 183 J=1,32
      KC=EPV(J)-NPOOF
      IF(KC.LE.0) GO TO 183
      SRR(KR,KC)=SRR(KR,KC)+SW(I,J)
      183 CONTINUE
      GO TO 170
C-----F DISPLACEMENT TYPES
      171 DO 181 J=1,32
      KC=EPV(J)
      K1=KC-NFT
      K2=KC-NPOOF
      IF(KC.LE.NFT) SFF(KR,KC)=SFF(KR,KC)+SW(I,J)
      IF((KC.GT.NFT).AND.(KC.LE.NPOOF)) SFP(KR,K1)=SFP(KR,K1)+SW(I,J)
      IF(KC.GT.NPOOF) SFR(KR,K2)=SFR(KR,K2)+SW(I,J)
      181 CONTINUE
      GO TO 170
C-----P DISPLACEMENT TYPES
      172 KR=KR-NFT
      DO 182 J=1,32
      KC=EPV(J)
      IF(KC.LE.NFT) GO TO 182
      K1=KC-NFT

```

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PAN:1200
PAN:1210
PAN:1220
PAN:1230
PAN:1240
PAN:1250
PAN:1260
PAN:1270
PAN:1280
PAN:1290
PAN:1300
PAN:1310
PAN:1320
PAN:1330
PAN:1340
PAN:1350
PAN:1360
PAN:1370
PAN:1380
PAN:1390
PAN:1400
PAN:1410
PAN:1420
PAN:1430
PAN:1440
PAN:1450
PAN:1460
PAN:1470
PAN:1480
PAN:1490
PAN:1500
PAN:1510
PAN:1520
PAN:1530
PAN:1540
PAN:1550
PAN:1560
PAN:1570
PAN:1580
PAN:1590
PAN:1600
PAN:1610
PAN:1620
PAN:1630
PAN:1640
PAN:1650
PAN:1660
PAN:1670
PAN:1680
PAN:1690
PAN:1700
PAN:1710
PAN:1720
PAN:1730
PAN:1740
PAN:1750
PAN:1760
PAN:1770
PAN:1780
PAN:1790

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PANL1800  
PANL1810  
PANL1820  
PANL1830  
PANL1846  
PANL1853  
PANL1863  
PANL1870  
PANL1880  
PANL1890  
PANL1900  
PANL1910  
PANL1920  
PANL1930  
PANL1940  
PANL1950  
PANL1960  
PANL1970  
PANL1980  
PANL1990  
PANL2000  
PANL2010  
PANL2620  
PANL2630  
PANL2040

```

SUBROUTINE FRAMSH
C
C-----SUBPROGRAM TO FORM THE FRAME STIFFNESS (ISM=1) AND MASS (ISM=2) AND PERMUTATION VECTOR (ISM=3)
C MCODE NOT EQUAL TO 3) SUBMATRICES SHM, SMP, SMR, AND MMH, MMH, MHRFAC(0,40)
C *** UNLESS *** NFM=0 (I.E., STRUCTURE = APPLICABLY-SUPPORTED)
C PANEL ONLY)
C
REAL MF,L,IX1,IX,IY,IZ
INTEGER SN,PPV,FPV,EPV(12)
C
COMMON
2/CBNFJ/XF(16),YF(16),MF(16),IX(16),IY(16),IZ(16),A(16),L(16),
0 E7J(16),EZK(16),EYJ(16),ZYK(16),IFM(16),JJ(16),JK(16)
3/CBNPFD/PPV(200),FPV(96),NFL(200),NRL(200),JRL(96)
4/CBSTIF/SFF( 40, 40),SFP( 40, 120),SFM( 40, 45),SFR( 40, 56),
0 SPP(120,120),SPH(120,45),SPR(120, 56),
0 SHM( 45, 45),SMR( 45, 56),
0 SRR( 56, 56)
6/CBMISC/SW(32,32),SMD(12,12),DUM(20)
7/CBDATA/EP,TP,PRP,DENS,EF,GF,TF,DENSF,NS,NM,MCODE,ISM,
0 NRPN,NFT,NPT,NMT,NRT,NF,NP,NH,NR,NPN,NPE,NPD,NPODF,
0 NFJ,NFM,NJS,NFNSPR,NPNSPR,ISDA,NPR,NF,NR,FJ
0 NFRT,IREAC,ISTRES,ISTAT,IL,IKKK,TIME
C
C-----CLEAR ARRAYS
C
DO 100 I=1,NMT
DO 101 J1=1,NMT
101 SHM(I,J1)=0.
DO 102 J2=1,NRT
102 SMR(I,J2)=0.
DO 104 J3=1,NPT
104 SPH(I,J3)=0.
DO 105 J4=1,NFT
105 SFH(J4,I)=0.
100 CONTINUE
IF(NRPN,GT,0)GO TO 1031
DO 103 I=1,NRT
DO 103 J=1,NRT
103 SRR(I,J)=0.
C
C-----FORM FRAME PERMUTATION VECTOR
C
1031 I1=0
I2=NMT+NPR
DO 104 I=1,NJ6
IF(JRL(I),EQ,0)GO TO 105
C-----R TYPE ODF
I2=I2+1
FPV(I)=I2
GO TO 104
C-----M TYPE ODF
105 I1=I1+1
FPV(I)=I1
104 CONTINUE
C
C-----FORM STIFFNESS SUBMATRICES SHM, SMP, AND SRF
C

```

```

00 106 I=1,NFM
00 107 I1=1,12
00 107 J1=1,12
107 SMD(I1,J1)=0.
IF(ISH.EQ.2)GO TO 150
IF(IFM(I).EQ.0)GO TO 108
C-----FORM BEAM ELEMENT STIFFNESS MATRIX (INCLUDES NONRIGID CONNECTIONS)
E22J=2*EZJ(I)+1
E22K=2*EZK(I)+1
E23J=3*EZJ(I)+1
E23K=3*EZK(I)+1
E21S=E2J(I)+EZK(I)+1
E2S=12*EZJ(I)*EZK(I)+4*(EZJ(I)+EZK(I))+1
EY2J=2*EYJ(I)+1
EY2K=2*EYK(I)+1
EY3J=3*EYJ(I)+1
EY3K=3*EYK(I)+1
EY1S=EYJ(I)+EYK(I)+1
EYS=12*EYJ(I)*EYK(I)+4*(EYJ(I)+EYK(I))+1
XL=(12*EF)/(L(I))*L(I)*L(I)
B1=EF*A(I)/L(I)
B2=(EZ1S/EZS)*XL*I2(I)
B3=(EY1S/EYS)*XL*IY(I)
B4=(XL*L(I)/2)*I2(I)*EZ2J/EZS
B5=(XL*L(I)/2)*IY(I)*EY2J/EYS
B6=GF*I2(I)/L(I)
B7=(2*EF*I2(I))/(L(I)*EZS)
B8=(2*EF*IY(I))/(L(I)*EYS)
SMD(1,1)=SMD(7,7)=B1
SMD(7,1)=SMD(1,7)=B1
SMD(2,2)=SMD(8,8)=B2
SMD(8,2)=SMD(2,8)=B2
SMD(3,3)=SMD(9,9)=B3
SMD(9,3)=SMD(3,9)=B3
SMD(12,2)=SMD(2,12)=B4
SMD(12,8)=SMD(8,12)=B4
SMD(6,2)=SMD(2,6)=B4*EZ2K/EZ2J
SMD(6,6)=SMD(6,6)=B4*EZ2K/EZ2J
SMD(11,9)=SMD(9,11)=B5
SMD(11,3)=SMD(3,11)=B5
SMD(5,3)=SMD(3,5)=B5*EY2K/EY2J
SMD(5,9)=SMD(9,5)=B5*EY2K/EY2J
SMD(4,4)=SMD(10,10)=B6
SMD(4,10)=SMD(10,4)=B6
SMD(5,5)=SMD(11,11)=2*B8*EY3K
SMD(11,5)=SMD(5,11)=B8
SMD(6,6)=SMD(12,12)=2*B7*EZ3K
SMD(12,6)=SMD(6,12)=B7
GO TO 109
C-----FORM COLUMN ELEMENT STIFFNESS MATRIX (WITHOUT NONRIGID CONNECTIONS)
108 C1=12*EF/(L(I))*L(I)*L(I)
C2=C1*L(I)/2
C3=EF*A(I)/L(I)
C4=GF*IY(I)/L(I)
C5=4*EF/L(I)
SMD(1,1)=SMD(7,7)=C1*I2(I)
SMD(1,7)=SMD(7,1)=C1*I2(I)
SMD(12,7)=SMD(7,12)=SMD(6,7)=SMD(7,6)=C2*I2(I)
SMD(6,1)=SMD(1,6)=SMD(1,12)=SMD(12,1)=C2*I2(I)

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```

FRA00600
FRA00610
FRA00620
FRA00630
FRA00640
FRA00650
FRA00660
FRA00670
FRA00680
FRA00690
FRA00700
FRA00710
FRA00720
FRA00730
FRA00740
FRA00750
FRA00760
FRA00770
FRA00780
FRA00790
FRA00800
FRA00810
FRA00820
FRA00830
FRA00840
FRA00850
FRA00860
FRA00870
FRA00880
FRA00890
FRA00900
FRA00910
FRA00920
FRA00930
FRA00940
FRA00950
FRA00960
FRA00970
FRA00980
FRA00990
FRA01000
FRA01010
FRA01020
FRA01030
FRA01040
FRA01050
FRA01060
FRA01070
FRA01080
FRA01090
FRA01100
FRA01110
FRA01120
FRA01130
FRA01140
FRA01150
FRA01160
FRA01170
FRA01180
FRA01190

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SMD(2,2)=SMD(8,8)=C3
SMD(2,8)=SMD(8,2)=C3
SMD(3,3)=SMD(9,9)=C1*I2(I)
SMD(3,9)=SMD(9,3)=C1*I2(I)
SMD(4,3)=SMD(3,4)=SMD(3,10)=SMD(10,3)=C2*I2(I)
SMD(10,9)=SMD(9,10)=SMD(4,9)=SMD(9,4)=SMD(4,3)
SMD(4,4)=SMD(10,10)=C5*I2(I)
SMD(4,10)=SMD(10,4)=SMD(4,4)/2
SMD(5,5)=SMD(11,11)=C4
SMD(5,11)=SMD(11,5)=C4
SMD(6,6)=SMD(12,12)=C5*I2(I)
SMD(6,12)=SMD(12,6)=SMD(6,6)/2
GO TO 109
C
C-----ASSEMBLE MASS SUBMATRICES MMH, MMR, MRR; REUSE STORAGE AND ARRAY
C NAMES OF STIFFNESS COMPUTATIONS
C
150 IF(MCODE.EQ.1) GO TO 151
C-----ASSEMBLED LUMPED MASS (ALM) FORMULATION
XM=DENSF*A(I)*L(I)
SMD(1,1)=SMD(2,2)=SMD(3,3)=SMD(7,7)=SMD(8,8)=SMD(9,9)=XM
GO TO 109
C-----CONSISTENT MASS (CM) FORMULATION
151 A1=DENSF*A(I)*L(I)/3
IX1=IY(I)+I2(I)
IF(IFM(I).EQ.0)IX1=IX(I)+I2(I)
A2=A1*IX1/A(I)
A3=A1*39./35.
A4=A1*27./70.
A5=(A1*L(I)*L(I))/35.
A6=(A1*L(I)*33.)/210.
A7=-A5*0.75
A8=(A6*13.)/22.
IF(IFM(I).EQ.1)GO TO 152
C-----COLUMNS--CONSISTENT MASS (CM)
SMD(2,2)=SMD(8,8)=A1
SMD(1,1)=SMD(3,3)=SMD(7,7)=SMD(9,9)=A3
SMD(5,5)=SMD(11,11)=A2
SMD(5,11)=SMD(11,5)=A2/2.
SMD(1,8)=SMD(8,2)=A1/2.
SMD(4,4)=SMD(10,10)=SMD(6,6)=SMD(12,12)=A5
SMD(4,10)=SMD(10,4)=SMD(6,12)=SMD(12,6)=A7
SMD(4,12)=SMD(12,7)=SMD(3,4)=SMD(4,3)=A6
SMD(9,10)=SMD(10,9)=SMD(1,6)=SMD(6,11)=A6
SMD(7,1)=SMD(1,7)=SMD(9,3)=SMD(3,9)=A4
SMD(12,1)=SMD(1,12)=SMD(4,9)=SMD(9,4)=A8
SMD(1,3)=SMD(3,10)=SMD(6,7)=SMD(7,6)=A8
GO TO 109
C-----BEAMS--CONSISTENT MASS (CM)
152 SMD(4,4)=SMD(10,10)=A2
SMD(4,10)=SMD(10,4)=A2/2.
SMD(1,1)=SMD(7,7)=A1
SMD(12,2)=SMD(8,8)=SMD(3,3)=SMD(9,9)=A3
SMD(1,7)=SMD(7,1)=A1/2.
SMD(5,5)=SMD(11,11)=SMD(6,6)=SMD(12,12)=A5
SMD(5,11)=SMD(11,5)=SMD(6,12)=SMD(12,6)=A7
SMD(6,2)=SMD(2,6)=SMD(9,11)=SMD(11,9)=A6
SMD(5,3)=SMD(3,5)=SMD(8,12)=SMD(12,8)=A6
SMD(2,8)=SMD(8,2)=SMD(3,9)=SMD(9,3)=A4
SMD(6,8)=SMD(8,6)=SMD(11,3)=SMD(3,11)=A8

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```

FRA 1200
FRA01210
FRA01220
FRA01230
FRA01240
FRA01250
FRA01260
FRA01270
FRA01280
FRA01290
FRA01300
FRA01310
FRA01320
FRA01330
FRA01340
FRA01350
FRA01360
FRA01370
FRA01380
FRA01390
FRA01400
FRA01410
FRA01420
FRA01430
FRA01440
FRA01450
FRA01460
FRA01470
FRA01480
FRA01490
FRA01500
FRA01510
FRA01520
FRA01530
FRA01540
FRA01550
FRA01560
FRA01570
FRA01580
FRA01590
FRA01600
FRA01610
FRA01620
FRA01630
FRA01640
FRA01650
FRA01660
FRA01670
FRA01680
FRA01690
FRA01700
FRA01710
FRA01720
FRA01730
FRA01740
FRA01750
FRA01760
FRA01770
FRA01780
FRA01790

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      SMD(5,9)=SMD(9,5)=SMD(12,2)=SMD(2,12)=-A8
C
C-----ASSEMBLE ELEMENT PERMUTATION VECTOR (EPV) FROM FRAME PERMUTATION
C      VECTOR (FPV)
C
      109 JJ1=0
          JA=6*JJ(I)-6
          DO 110 J=1,6
              JA=JA+1
              JJ1=JJ1+1
      110 EPV(JJ1)=FPV(JA)
          JB=6*JK(I)-6
          DO 111 J=1,6
              JB=JB+1
              JJ1=JJ1+1
      111 EPV(JJ1)=FPV(JB)
C
C-----USE EPV(I) TO ADD CONTRIBUTION OF CURRENT FRAME MEMBER (I) TO
C      STIFFNESS ARRAYS SHM, SMR, SRR (OR CORRESPONDING MASS ARRAYS FOR
C      ISM=2)
C
          DO 112 K=1,12
              KR=EPV(K)
              IF(KR.LE.NMT)GO TO 113
C-----R DISPLACEMENT TYPES
              KR=KR-NMT
              DO 114 J=1,12
                  KC=EPV(J)-NMT
                  IF(KC.LE.0)GO TO 114
                  SRR(KR,KC)=SRR(KR,KC)+SMD(K,J)
              114 CONTINUE
              GO TO 112
C-----M DISPLACEMENT TYPES
      113 DO 115 J=1,12
          KC=EPV(J)
          K1=KC-NMT
          IF(KC.LE.NMT)GO TO 116
          SMR(KR,K1)=SMR(KR,K1)+SMD(K,J)
          GO TO 115
      116 SHM(KR,KC)=SHM(KR,KC)+SMD(K,J)
      115 CONTINUE
      112 CONTINUE
C
      106 CONTINUE
C
      IF(ISM.EQ.1)RETURN
C
C-----ADD LUMPED MASSES TO FRAME FOR MCODE=1,2 CASES
C
          DO 200 I=1,NFJ
              EPV(1)=6*I-5
              EPV(2)=EPV(1)+1
              EPV(3)=EPV(1)+2
              DO 201 K=1,3
                  J=FPV(EPV(K))
                  IF(J.LE.NMT)ISM(J,J)=SHM(J,J)+MF(I)
                  IF(J.GT.NMT)SRR(J,J)=SRR(J,J)+MF(I)
              201 CONTINUE
          200 CONTINUE
C

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      FRA11800
      FRA01810
      FRA01820
      FRA01830
      FRA01840
      FRA01850
      FRA01860
      FRA01870
      FRA01880
      FRA01890
      FRA01900
      FRA01910
      FRA01920
      FRA01930
      FRA01940
      FRA01950
      FRA01960
      FRA01970
      FRA01980
      FRA01990
      FRA02000
      FRA02010
      FRA02020
      FRA02030
      FRA02040
      FRA02050
      FRA02060
      FRA02070
      FRA02080
      FRA02090
      FRA02100
      FRA02110
      FRA02120
      FRA02130
      FRA02140
      FRA02150
      FRA02160
      FRA02170
      FRA02180
      FRA02190
      FRA02200
      FRA02210
      FRA02220
      FRA02230
      FRA02240
      FRA02250
      FRA02260
      FRA02270
      FRA02280
      FRA02290
      FRA02300
      FRA02310
      FRA02320
      FRA02330
      FRA02340
      FRA02350
      FRA02360
      FRA02370
      FRA02380
      FRA02390

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      RETURN
      END

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      FRA02400
      FRA02410

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*DECK SPRING
SUBROUTINE SPRING
C
C-----SUBPROGRAM TO ADD BOUNDARY STIFFNESS TERMS (MODELLED AS LINEARLY-
ELASTIC SPRINGS) TO STIFFNESS SUMMATRICES FOR FRAME-PANEL
ASSEMBLAGE (NFM>0) OR FOR PANEL-ONLY CASE (NFM=0).
C
C BOUNDARY SPRINGS: FORM SFM, SPH AND AUGMENT SFF, SPP, SMH, SRR,
SFR, SPR, AND SMR
C
C PANEL SPRINGS: AUGMENT SFF, SPP, AND SRR FOR ARBITRARY NUMBER
AND LOCATION OF NODAL SPRINGS
C
C INTEGER F1,P1,F2,P2,P3,PPV,FPV,SN,PNPV(8),FNPV(6)
REAL KS1,KS2,MP
C
COMMON
1/CBNPN/XP(25),YP(25),MP(25),KS1(25,5),JFN(25,2),JPN(25),INTYP(25),
0 IREP(16),ICNTY(16,4),KS2(25,5)
3/CBNPFD/PPV(200),FPV(96),NFL(200),NRL(200),JRL(96)
4/CBSTIF/SFF( 40, 40),SFF( 40, 120),SFM( 40, 45),SFR( 40, 56),
0 SPP(120,120),SPH(120,45),SPR(120, 56),
0 SMH( 45, 45),SMR( 45, 56),
0 SRR( 56, 56)
7/CBDDATA/EP,TP,PRP,DENSP,EF,GF,TF,DENSF,NS,SN,MCODE,ISM,
0 NRPN,NFT,NPT,NMT,NRT,NF,NP,NM,NR,NPN,NPE,NPO,NPDOF,
0 NFJ,NFM,NJ6,NFNSPR,NPNSPR,ISQA,NPR,NFR,NFJ,
0 NFRT,IREAC,ISTRES,ISTAT,IL,IKKK,TIME
C
C-----CLEAR ARRAYS
C
DO 90 I=1,NMT
DO 91 J1=1,NFT
91 SFM(J1,I)=0.
DO 92 J2=1,NPT
92 SPH(J2,I)=0.
90 CONTINUE
C
IF(NRPN.GT.0)GO TO 151
C
DO 93 I=1,NRT
DO 94 J1=1,NFT
94 SFR(J1,I)=0.
DO 95 J2=1,NPT
95 SPR(J2,I)=0.
93 CONTINUE
C
C-----FRAME-PANEL ASSEMBLY BOUNDARY SPRINGS (NFM>0)
C
151 IF(NFNSPR.EQ.0)GO TO 150
DO 50 I=1,NFNSPR
C
JFN(I,1) AND JFN(I,2) = FRAME AND CORRESPONDING PANEL NODES
(RESPECTIVELY) WHERE BOUNDARY SPRINGS ARE ATTACHED
C
JF1=6*JFN(I,1)-6
JP1=8*JFN(I,2)-8
C
SPR(000)
SPR(0010)
SPR(0020)
SPR(0030)
SPR(0040)
SPR(0050)
SPR(0060)
SPR(0070)
SPR(0080)
SPR(0090)
SPR(0100)
SPR(0110)
SPR(0120)
SPR(0130)
SPR(0140)
SPR(0150)
SPR(0160)
SPR(0170)
SPR(0180)
SPR(0190)
SPR(0200)
SPR(0210)
SPR(0220)
SPR(0230)
SPR(0240)
SPR(0250)
SPR(0260)
SPR(0270)
SPR(0280)
SPR(0290)
SPR(0300)
SPR(0310)
SPR(0320)
SPR(0330)
SPR(0340)
SPR(0350)
SPR(0360)
SPR(0370)
SPR(0380)
SPR(0390)
SPR(0400)
SPR(0410)
SPR(0420)
SPR(0430)
SPR(0440)
SPR(0450)
SPR(0460)
SPR(0470)
SPR(0480)
SPR(0490)
SPR(0500)
SPR(0510)
SPR(0520)
SPR(0530)
SPR(0540)
SPR(0550)
SPR(0560)
SPR(0570)
SPR(0580)
SPR(0590)
C
FORM FRAME NODE PERMUTATION VECTOR (FNPV) FOR NODE JFN(I,1) FROM
FPV(I)
C
DO 51 J=1,6
JF1=JF1+1
51 FNPV(J)=FPV(JF1)
C
FORM PANEL NODE PERMUTATION VECTOR (PNPV) FOR NODE JFN(I,2) FROM
PPV(I)
C
DO 52 J=1,8
JP1=JP1+1
52 PNPV(J)=PPV(JP1)
C
FORM ARRAYS SFM,SPH; AUGMENT ARRAYS SFF,SPP,SMH,SRR,SFR,SPR,SMR
WITH SPRING STIFFNESSES KS1(I,J) (FOR DOFS 1,...,5 ONLY)
C
DO 53 J=1,5
C
F1 1...NFT; P1 1...NPT; M1 1...NMT; R1 1...NRT
C
P1=PNPV(J)
F1=FNPV(J)
P2=P1-NFT
F2=F1-NMT
P3=P2-NPT
IF(P2)55,55,561
561 IF(P3)56,56,562
C-----F, M, R TYPE DOF'S; FORM SFM, AUGMENT SFF, SMH, SRR, SFR
55 SFF(P1,P1)=SFF(P1,P1)+KS1(I,J)
IF (F2)57,57,58
57 SFH(P1,F1)=-KS1(I,J)
SMH(F1,F1)=SMH(F1,F1)+KS1(I,J)
GO TO 53
58 SFR(P1,F2)=-KS1(I,J)
SRR(F2,F2)=SRR(F2,F2)+KS1(I,J)
GO TO 53
C
P, M, R TYPE DOF'S; FORM SPH, AUGMENT SPP, SMH, SRR, SPR
C
56 SPP(P2,P2)=SPP(P2,P2)+KS1(I,J)
IF(F2)59,59,60
59 SPH(P2,F1)=-KS1(I,J)
SMH(F1,F1)=SMH(F1,F1)+KS1(I,J)
GO TO 53
60 SPR(P2,F2)=-KS1(I,J)
SRR(F2,F2)=SRR(F2,F2)+KS1(I,J)
C
GO TO 53
562 SRR(P3,P3)=SRR(P3,P3)+KS1(I,J)
IF(F2)563,563,564
563 SMR(F2,P3)=-KS1(I,J)
SMH(F2,F2)=SMH(F2,F2)+KS1(I,J)
GO TO 53
564 SRR(P3,F2)=SRR(P3,F2)-KS1(I,J)
SRR(F2,P3)=SRR(P3,F2)
53 CONTINUE
50 CONTINUE
C-----SPRINGS AT PANEL NODES (NOT CONNECTED TO FRAME)
SPR(0600)
SPR(0610)
SPR(0620)
SPR(0630)
SPR(0640)
SPR(0650)
SPR(0660)
SPR(0670)
SPR(0680)
SPR(0690)
SPR(0700)
SPR(0710)
SPR(0720)
SPR(0730)
SPR(0740)
SPR(0750)
SPR(0760)
SPR(0770)
SPR(0780)
SPR(0790)
SPR(0800)
SPR(0810)
SPR(0820)
SPR(0830)
SPR(0840)
SPR(0850)
SPR(0860)
SPR(0870)
SPR(0880)
SPR(0890)
SPR(0900)
SPR(0910)
SPR(0920)
SPR(0930)
SPR(0940)
SPR(0950)
SPR(0960)
SPR(0970)
SPR(0980)
SPR(0990)
SPR(1000)
SPR(1010)
SPR(1020)
SPR(1030)
SPR(1040)
SPR(1050)
SPR(1060)
SPR(1070)
SPR(1080)
SPR(1090)
SPR(1100)
SPR(1110)
SPR(1120)
SPR(1130)
SPR(1140)
SPR(1150)
SPR(1160)
SPR(1170)
SPR(1180)
SPR(1190)

```

```

C
150 IF (NPNSPR.EQ.0) RETURN
   DO 100 I=1,NPNSPR
C
   JPN(I) = NODE NUMBER OF PANEL WITH SPRINGS ATTACHED
C
   JP1=8*JPN(I)-8
C
   FORM PANEL NODE PERMUTATION VECTOR (PNPV) FOR PANEL NODE JP
   FROM PPV(I)
C
   DO 101 J=1,8
   JP1=JP1+1
101 PNPV(J)=PPV(JP1)
C
C-----AUGMENT ARRAYS SFF,SPP,SRR WITH NODE JPN(I) SPRING STIFFNESSES
C
   DO 102 J=1,5
   K1=PNPV(J)
   K2=K1-NFT
   K3=K2-NPT
   IF (K1.GT.NFT) GO TO 1031
   SFF(K1,K1)=SFF(K1,K1)+KS2(I,J)
   GO TO 102
1031 IF (K3.I103,103,1032
103 SPP(K2,K2)=SPP(K2,K2)+KS2(I,J)
   GO TO 102
1032 SRR(K3,K3)=SRR(K3,K3)+KS2(I,J)
102 CONTINUE
100 CONTINUE
C
   RETURN
   END

```

```

SPR,1200
SPR,1210
SPR,1220
SPR,1230
SPR,1240
SPR,1250
SPR,1260
SPR,1270
SPR,1280
SPR,1290
SPR,1300
SPR,1310
SPR,1320
SPR,1330
SPR,1340
SPR,1350
SPR,1360
SPR,1370
SPR,1380
SPR,1390
SPR,1400
SPR,1410
SPR,1420
SPR,1430
SPR,1440
SPR,1450
SPR,1460
SPR,1470
SPR,1480
SPR,1490
SPR,1500
SPR,1510
SPR,1520

*DECK CDENSE
SUBROUTINE CDENSE
C
C-----SUBPROGRAM TO ELIMINATE P AND H TYPES (P TYPES ONLY FOR
C NFM=0 CASE) FROM STIFFNESS (ISM=1) AND MASS (ISM=2) MATRICES
C USING GUYAN REDUCTION
C
   REAL TPF(120,45),TPM(120,45),TMFS(120,45),TPR(120,56),THRS(120,56)
   INTEGER SN
C
   COMMON
   4/CBSTIF/SFF( 40, 40),SFP( 40, 120),SFM( 40, 45),SFR( 40, 56),
   # SPP(120,120),SPH(120,45),SPR(120, 56),
   # SHM( 45, 45),SMR( 45, 56),
   # SRR( 56, 56)
   5/CBSTOR/TS1(120,45),TS2(120, 56),TS3( 45, 120)
   7/CBDATA/EP,TP,PRP,DENSP,EF,GF,TF,DENSEF,NS,SN,MDOE,ISM,
   # NRP,NFT,NPT,NMT,NRT,NF,NP,NM,NR,NPN,NPE,NPO,NPOOF,
   # NFJ,NFM,NJB,NFNSPR,NPNSPR,ISDA,NPR,NFR,NRJ,
   # NFRT,IREAC,ISTRES,ISTAT,IL,IKKK,TIME
C
   EQUIVALENCE (TPF,TPM,TMFS,TS1),ITPR,THRS,TS2)
C
   IF (ISM.EQ.2) GO TO 150
C
C-----SAVE SFR AND SPR (ALSO SMR FOR NFM=0 CASE) FOR REACTION CALCULATION
C
   NS
C
   REWIND 11
   WRITE(11) ((SFR(I,J),J=1,NRT),I=1,NFT)
   WRITE(11) ((SPR(I,J),J=1,NRT),I=1,NPT)
   WRITE(11) ((SRR(I,J),J=1,NRT),I=1,NRT)
   IF (NFM.EQ.0) GO TO 200
   WRITE(11) ((SMR(I,J),J=1,NRT),I=1,NMT)
200 CONTINUE
C
C-----STIFFNESS CONDENSATION: COMPUTE SFF* AND SFR* (ALSO SFM*,
C SHM*, AND SMR* IF NFM=0)
C
   REWIND 10
   CALL DCOMP(NPT,SPP,NP,NPE,RETURNS(100))
   GO TO 101
100 PRINT 102
102 FORMAT(///,* DECOMPOSITION OF SPP FAILS IN CDENSE*,///)
   RETURN
101 DO 103 I=1,NFT
103 CALL SOLVE(NPT,SPP,SPP,TPF,I,2,NFT,NP,NP,NF,NP,NP,NM)
C
   IF (MDOE.NE.3) WRITE(10) ((TPF(I,J),J=1,NFT),I=1,NPT)
C
   DO 106 I=1,NFT
   DO 106 J=1,NFT
   TEMP=SFF(I,J)
   DO 107 JJ=1,NPT
107 TEMP=TEMP-SFP(I,JJ)*TPF(JJ,J)
106 SFF(I,J)=TEMP
C
   DO 108 I=1,NFT
   DO 108 J=1,NRT
   TEMP=SFR(I,J)

```

```

COE/1200
COE/1210
COE/1220
COE/1230
COE/1240
COE/1250
COE/1260
COE/1270
COE/1280
COE/1290
COE/1300
COE/1310
COE/1320
COE/1330
COE/1340
COE/1350
COE/1360
COE/1370
COE/1380
COE/1390
COE/1400
COE/1410
COE/1420
COE/1430
COE/1440
COE/1450
COE/1460
COE/1470
COE/1480
COE/1490
COE/1500
COE/1510
COE/1520
COE/1530
COE/1540
COE/1550
COE/1560
COE/1570
COE/1580
COE/1590

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```

C      DO 109 JJ=1,NPT
109  TEMP=TEMP-TPF(JJ,I)*SPR(JJ,J)
108  SFR(I,J)=TEMP
C
      IF(MCODE.EQ.3)GO TO 149
C
      DO 130 I=1,NRT
130  CALL SOLVE(NPT,SPP,SPR,TPR,I,1,NRT,NP,NP,NP,NK,NP,NR)
      WRITE(10) ((TPR(I,J),J=1,NRT),I=1,NPT)
C
149  IF(NFM.EQ.0)RETURN
C
C-----NOW COMPUTE SFM*, SMH*, AND SMR* FOR NFM>0 CASE
C
      DO 131 I=1,NMT
131  CALL SOLVE(NPT,SPP,SPM,TPM,I,1,NMT,NP,NP,NP,NM,NP,NM)
      IF(MCODE.NE.3) WRITE(10) ((TPM(I,J),J=1,NMT),I=1,NPT)
C
      DO 111 I=1,NFT
      DO 111 J=1,NMT
      TEMP=SFM(I,J)
      DO 112 JJ=1,NPT
112  TEMP=TEMP-SFM(I,JJ)*TPM(JJ,J)
111  SFM(I,J)=TEMP
C
      DO 113 I=1,NMT
      DO 113 J=1,NMT
      TEMP=SMH(I,J)
      DO 114 JJ=1,NPT
114  TEMP=TEMP-SPM(JJ,I)*TPM(JJ,J)
113  SMH(I,J)=TEMP
C
      DO 123 I=1,NMT
      DO 123 J=1,NRT
      TEMP=SMR(I,J)
      DO 124 JJ=1,NPT
124  TEMP=TEMP-TPM(JJ,I)*SPR(JJ,J)
123  SMR(I,J)=TEMP
C
C-----NOW COMPUTE SFF** AND SFR**
C
      CALL OCOMPL(NMT,SMH,NM,NM),RETURNS(115)
      GO TO 116
115  PRINT 117
117  FORMAT(///,30H DECOMPOSITION OF SMH* FAILS IN CGENSE,///)
      RETURN
116  DO 118 I=1,NFT
118  CALL SOLVE(NMT,SMH,SFM,TPFS,I,2,NFT,NM,NM,NF,NM,NP,NM)
      IF(MCODE.NE.3) WRITE(10) ((TMFS(I,J),J=1,NFT),I=1,NMT)
C
      DO 119 I=1,NFT
      DO 119 J=1,NFT
      TEMP=SFF(I,J)
      DO 120 JJ=1,NMT
120  TEMP=TEMP-SFM(I,JJ)*TMFS(JJ,J)
119  SFF(I,J)=TEMP
C
      DO 121 I=1,NFT
      DO 121 J=1,NFT
      TEMP=SFR(I,J)

```

```

CODE0000      DO 122 JJ=1,NMT                                COE01200
CODE0010      122 TEMP=TEMP-TMFS(JJ,I)*SMR(JJ,J)             COE01210
CODE0020      121 SFR(I,J)=TEMP                               COE01220
CODE0030      C                                              COE01230
CODE0040      IF (MCOE.EQ.3) RETURN                            COE01240
CODE0050      C                                              COE01250
CODE0060      DO 125 I=1,NRT                                  COE01260
CODE0070      125 CALL SOLVE(NMT,SMH,SMR,TMRS,I,1,NRT,NM,NM,NM,NR,NP,NRI COE01270
CODE0080      WRITE(10) ((TMRS(I,J),J=1,NRT),I=1,NMT)        COE01280
CODE0090      C                                              COE01290
CODE0100      RETURN                                           COE01300
CODE0110      C                                              COE01310
CODE0120      C-----MASS CONDENSATION                       COE01320
CODE0130      C                                              COE01330
CODE0140      150 IF (MCOE.EQ.3) RETURN                        COE01340
CODE0150      REMIND 10                                         COE01350
CODE0160      C                                              COE01360
CODE0170      C-----COMPUTE MFF* = MFF - MFPXTPF - (MFPXTPF)' + TPF*XMPXTPF COE01370
CODE0180      C AND MFR* = MFR - MFPXTPR - TPF*XMR + TPF*XMPXTPR COE01380
CODE0190      C LET TS3 = TPF*XMPX                             COE01390
CODE0200      C                                              COE01400
CODE0210      READ(10) ((TPF(I,J),J=1,NFT),I=1,NPT)          COE01410
CODE0220      READ(10) ((TPR(I,J),J=1,NRT),I=1,NPT)          COE01420
CODE0230      C                                              COE01430
CODE0240      DO 151 I=1,NFT                                   COE01440
CODE0250      DO 151 J=1,NPT                                   COE01450
CODE0260      TEMP=0.                                           COE01460
CODE0270      DO 152 JJ=1,NPT                                   COE01470
CODE0280      152 TEMP=TEMP+TPF(JJ,I)*SPP(JJ,J)              COE01480
CODE0290      151 TS3(I,J)=TEMP                                COE01490
CODE0300      C                                              COE01500
CODE0310      DO 153 I=1,NFT                                   COE01510
CODE0320      DO 153 J=1,NFT                                   COE01520
CODE0330      QP=SFF(I,J)                                       COE01530
CODE0340      DO 154 JJ=1,NPT                                   COE01540
CODE0350      154 QP=QP-SFP(I,JJ)*TPF(JJ,J)-TPF(JJ,I)*SFP(J,JJ)+TS3(I,JJ)*TPF(JJ,J) COE01550
CODE0360      153 SFF(I,J)=QP                                   COE01560
CODE0370      C                                              COE01570
CODE0380      DO 155 I=1,NFT                                   COE01580
CODE0390      DO 155 J=1,NRT                                   COE01590
CODE0400      QP=SFR(I,J)                                       COE01600
CODE0410      DO 156 JJ=1,NPT                                   COE01610
CODE0420      156 QP=QP-SFP(I,JJ)*TPR(JJ,J)-TPF(JJ,I)*SPR(JJ,J)+TS3(I,JJ)*TPR(JJ,J) COE01620
CODE0430      155 SFR(I,J)=QP                                   COE01630
CODE0440      C                                              COE01640
CODE0450      C                                              COE01650
CODE0460      IF (NFM.EQ.0) RETURN                             COE01660
CODE0470      C                                              COE01670
CODE0480      C-----NOW COMPUTE NFM*, MMH*, AND MMR* FOR NFM>0 CASE (PANEL-FRAME ASSEMBLY) COE01680
CODE0490      C                                              COE01690
CODE0500      C                                              COE01700
CODE0510      C MFM*=NFM-MFPXTPM-TPF*XMPX+TS3XTPM WHERE TS3=TPF*XMPX COE01710
CODE0520      C MMH*=MMH-MPH*XTPM-TPM*XNFM+TPM*XMPXTPM (NOW LET TS3=TPM*XMPX) COE01720
CODE0530      C MMR*=MMR-MPR*XTPR-TPM*XMPX+TPM*XMPXTPR (SAME TS3 AS FOR MMH*) COE01730
CODE0540      C                                              COE01740
CODE0550      DO 157 I=1,NFT                                   COE01750
CODE0560      DO 157 J=1,NMT                                   COE01760
CODE0570      QP=SFM(I,J)                                       COE01770
CODE0580      DO 158 JJ=1,NPT                                   COE01780
CODE0590      158 QP=QP-TPF(JJ,I)*SPM(JJ,J)                  COE01790

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```

157 SFM(I,J)=QP                                CODE 1840
C                                                CODE 1840
C-----REPLACE TPF WITH TPM (SHARE SAME STORAGE)  CODE 1824
C                                                CODE 1836
      READ(10) ((TPM(I,J),J=1,NMT),I=1,NPT)      CODE 1840
      DO 159 I=1,NPT                             CODE 1850
      DO 159 J=1,NMT                             CODE 1860
      QP=SFH(I,J)                                CODE 1870
      DO 160 JJ=1,NPT                             CODE 1880
160 QP=QP-SFP(I,JJ)*TPM(JJ,J)+TS3(I,JJ)*TPM(JJ,J)  CODE 1890
159 SFH(I,J)=QP                                CODE 1900
C                                                CODE 1910
C-----REDEFINE TS3 FOR DETERMINATION OF MMH* AND MMR*  CODE 1920
C                                                CODE 1930
      DO 161 I=1,NMT                             CODE 1940
      DO 161 J=1,NPT                             CODE 1950
      TEMP=J,                                     CODE 1960
      DO 162 JJ=1,NPT                             CODE 1970
162 TEMP=TEMP+TPM(JJ,I)*SPP(JJ,J)               CODE 1980
161 TS3(I,J)=TEMP                               CODE 1990
C                                                CODE 2000
      DO 163 I=1,NMT                             CODE 2010
      DO 163 J=1,NMT                             CODE 2020
      QP=SMH(I,J)                                CODE 2030
      DO 164 JJ=1,NPT                             CODE 2040
164 QP=QP-SPH(JJ,I)*TPM(JJ,J)-TPM(JJ,I)*SPH(JJ,J)+TS3(I,JJ)*TPM(JJ,J)  CODE 2050
163 SMH(I,J)=QP                                CODE 2060
C                                                CODE 2070
      DO 165 I=1,NMT                             CODE 2080
      DO 165 J=1,NPT                             CODE 2090
      QP=SMR(I,J)                                CODE 2100
      DO 166 JJ=1,NPT                             CODE 2110
166 QP=QP-SPH(JJ,I)*TPR(JJ,J)-TPM(JJ,I)*SPR(JJ,J)+TS3(I,JJ)*TPR(JJ,J)  CODE 2120
165 SMR(I,J)=QP                                CODE 2130
C                                                CODE 2140
C-----FINALLY COMPUTE MFF* AND MFP*              CODE 2150
C      WHERE                                         CODE 2160
C      MFF*=-MFF*-MFM*XTMF*-TMF*XMFM*+TS3XTMF*, TS3=TMF*XMFM*  CODE 2170
C      MFR*=-MFR*-MFM*XTMR*-TMF*XMFR*+TS3XTMR*  CODE 2180
C                                                CODE 2190
C-----REPLACE TPM AND TPR WITH TMFS AND TMRS, RESPECTIVELY  CODE 2200
C                                                CODE 2210
      READ(10) ((TMFS(I,J),J=1,NPT),I=1,NMT)      CODE 2220
      READ(10) ((TMRS(I,J),J=1,NPT),I=1,NMT)      CODE 2230
C                                                CODE 2240
      DO 167 I=1,NPT                             CODE 2250
      DO 167 J=1,NMT                             CODE 2260
      TEMP=J,                                     CODE 2270
      DO 168 JJ=1,NMT                             CODE 2280
168 TEMP=TEMP+TMFS(JJ,I)*SMH(JJ,J)               CODE 2290
167 TS3(I,J)=TEMP                               CODE 2300
C                                                CODE 2310
      DO 169 I=1,NPT                             CODE 2320
      DO 169 J=1,NPT                             CODE 2330
      QP=SFF(I,J)                                CODE 2340
      DO 170 JJ=1,NMT                             CODE 2350
170 QP=QP-SFM(I,JJ)*TMFS(JJ,J)-TMFS(JJ,I)*SFP(I,JJ)+TS3(I,JJ)*TMFS(JJ,J)  CODE 2360
      # J)                                         CODE 2370
169 SFF(I,J)=QP                                CODE 2380
C                                                CODE 2390

      DO 171 I=1,NPT                             CODE 2400
      DO 171 J=1,NPT                             CODE 2410
      QP=SFR(I,J)                                CODE 2420
      DO 172 JJ=1,NMT                             CODE 2430
172 QP=QP-SFM(I,JJ)*TMRS(JJ,J)-TMFS(JJ,I)*SFP(I,JJ)+TS3(I,JJ)*TMRS(JJ,J)  CODE 2440
      # J)                                         CODE 2450
171 SFP(I,J)=QP                                CODE 2460
C                                                CODE 2470
      RETURN                                       CODE 2480
      END                                          CODE 2490

```

```

*DECK EIGENV
SUBROUTINE EIGENV
C
C-----SUBPROGRAM TO SOLVE THE NON-STANDARD FORM OF THE EIGENVALUE
C      PROBLEM:
C      M*X = (1/P**2)*S*X
C      OR      A*X = LAMBOA*B*X
C
C      PARAMETERS:
C      ON INPUT -  N=NO. DOF
C                  A=MASS MATRIX
C                  B=STIFFNESS MATRIX
C      ON OUTPUT - SA=PSQ (SQUARES OF NAT. CIRC. FREQUENCIES)
C                  A=MASS MATRIX
C                  S=STIFFNESS MATRIX
C                  X=XM1 (INVERSE OF XM)
C                  B=XM (MODAL MATRIX NORMALIZED WRT MASS MATRIX)
C
C      REAL LAMBOA,MFR
C      INTEGER SN
C
C      COMMON
C      4/CBSTIF/ B( 40, 40),S( 40, 120),A( 40, 45),SFR( 40, 56),
C              SPP(120,120),SPM(120,45),MFR(120, 56),
C              SMH( 45, 45),SHR( 45, 56),
C              SRR( 56, 56)
C      5/CBSTOR/X(80,80),FREQ(100,3),SA(4060),LAMBOA(4060)
C      7/CODATA/EP,TP,PRP,DENSP,EF,GF,TF,DENSF,NS,SN,MCODE,ISM,
C              NRP,N,NFT,NPT,NMT,NRT,NF,NP,NM,NR,NPN,NPE,NPD,NPOOF,
C              NFJ,NFN,NJ6,NFNSPR,NPNSPR,ISDA,NPR,NFR,NRFJ,
C              NFR1,IREAC,ISTRES,ISTAT,IL,IKKK,TIME
C
C      N=NFT
C      NA=MF
C      NB=NF
C      NL=4060
C      NX=80
C
C      101 FORMAT(*1FREQUENCIES AND MODE SHAPES*,/,
C      *-----*)
C      102 FORMAT(*0*,* SYSTEM MASS MATPIX*)
C      103 FORMAT(*0*,* SYSTEM STIFFNESS MATRIX*)
C      104 FORMAT(*0*,* CIRCULAR FREQUENCIES (RAD/SEC,)* )
C      110 FORMAT(*0*,* NATURAL FREQUENCIES (HERTZ)* )
C      111 FORMAT(*0*,* NATURAL FREQUENCIES (SECONDS)* )
C      112 FORMAT(*0*,* MODAL MATRIX XM (SCALED TO MAKE LARGEST ENTRY 1.0)* )
C      113 FORMAT(*0*,* )
C      104 FORMAT(1P9E13,5)
C      105 FORMAT(* ERROR IN AXLEX *)
C
C      PRINT 101
C      DO 114 I=1,N
C      SA(I)=A(I,I)
C      DO 114 J=1,N
C      114 S(I,J)=B(I,J)
C      PRINT 102
C      CALL MATPRT(N,N,A,NA,NM)
C      PRINT 103
C      CALL MATPRT(N,N,B,NB,NB)
C      CALL AXLBX(N,NA,NB,NL,NX,A,B,LAMBOA,X),RETURNS(106)

```

```

EIG00000
EIG00010
EIG00020
EIG00030
EIG00040
EIG00050
EIG00060
EIG00070
EIG00080
EIG00090
EIG00100
EIG00110
EIG00120
EIG00130
EIG00140
EIG00150
EIG00160
EIG00170
EIG00180
EIG00190
EIG00200
EIG00210
EIG00220
EIG00230
EIG00240
EIG00250
EIG00260
EIG00270
EIG00280
EIG00290
EIG00300
EIG00310
EIG00320
EIG00330
EIG00340
EIG00350
EIG00360
EIG00370
EIG00380
EIG00390
EIG00400
EIG00410
EIG00420
EIG00430
EIG00440
EIG00450
EIG00460
EIG00470
EIG00480
EIG00490
EIG00500
EIG00510
EIG00520
EIG00530
EIG00540
EIG00550
EIG00560
EIG00570
EIG00580
EIG00590

C-----NOW RESTORE MATRIX A, THE MASS MATRIX
N1=N-1
A(N,N)=SA(N)
DO 115 I=1,M1
A(I,I)=SA(I)
J1=I+1
DO 115 J=J1,N
115 A(J,I)=A(I,J)
GO TO 107
106 PRINT 105
107 CONTINUE
DO 103 I=1,N
FREQ(I,1)=SQRT(1/LAMBOA(I))
FREQ(I,2)=FREQ(I,1)/6.283185308
103 FREQ(I,3)=1./FREQ(I,2)
PRINT 110
PRINT 104,(FREQ(I,1),I=1,N)
PRINT 111
PRINT 104,(FREQ(I,2),I=1,N)
PRINT 112
PRINT 104,(FREQ(I,3),I=1,N)
PRINT 113
CALL MATPRT(N,N,X,NX,NX)
C-----NORMALIZE THE MODAL MATRIX (XM) WRT THE MASS MATRIX TO OBTAIN XM
DO 118 I=1,N
DO 119 J=1,N
SA(J)=0.
DO 119 K=1,N
119 SA(J)=SA(J)+A(J,K)*X(K,I)
SUM=0.
DO 120 JJ=1,N
120 SUM=SUM+X(JJ,I)*SA(JJ)
TEMP=SQRT(SUM)
DO 121 KK=1,N
121 X(KK,I)=X(KK,I)/TEMP
118 CONTINUE
C-----COMPUTE (XM)**-1 = XM* X M
DO 122 I=1,N
DO 122 J=1,N
SUM=0.
DO 123 K=1,N
123 SUM=SUM+X(K,I)*A(K,J)
122 X(I,J)=SUM
C-----FINALLY, STORE SQUARES OF NATURAL CIRCULAR FREQUENCIES IN SA(I)
DO 124 I=1,N
124 SA(I)=FREQ(I,1)*FREQ(I,1)
C
RETURN
END

```

```

FIG00600
FIG00610
FIG00620
FIG00630
FIG00640
FIG00650
FIG00660
FIG00670
FIG00680
FIG00690
FIG00700
FIG00710
FIG00720
FIG00730
FIG00740
FIG00750
FIG00760
FIG00770
FIG00780
FIG00790
FIG00800
FIG00810
FIG00820
FIG00830
FIG00840
FIG00850
FIG00860
FIG00870
FIG00880
FIG00890
FIG00900
FIG00910
FIG00920
FIG00930
FIG00940
FIG00950
FIG00960
FIG00970
FIG00980
FIG00990
FIG01000
FIG01010
FIG01020
FIG01030
FIG01040
FIG01050
FIG01060
FIG01070
FIG01080

```



```
*DECK AXLBX
SUBROUTINE AXLBX(N,NA,NB,NL,NX,A,B,LAMBDA,X),RETURNS(JJ)
```

```
C
C-----SOLVES EIGENPROBLEM A*X = LAMBDA*B*X
C A SYMMETRIC, B SYMMETRIC POSITIVE DEFINITE
C INPUT: A AND B (ONLY DIAGONAL AND UPPER TRIANGLE NEEDED)
C OUTPUT: LAMBDA AND X
C ALTERED: DIAGONAL AND LOWER TRIANGLE OF A AND B
C ERROR RETURN: B IS NOT POSITIVE DEFINITE OR TQL2 DID NOT CONVERGE
C REFERENCES: NUMERISCHE MATHEMATIK, VOL. 11, PP. 99, 181 AND 293.
C
C*****NOTE: NA,...,NX ARE THE DIMENSIONS OF ARRAYS A, B, LAMBDA, AND X
C IN THE CALLING PROGRAM.
C
C-----CALLED BY MAIN PROGRAM
C
C-----REQUIRES SUBPROGRAMS TRE02 AND TQL2
C
C      REAL LAMBDA(NL),A(NA,NA),B(NB,NB),X(NX,NX),E(256)
C
C      DO 4 I=1,N
C      DO 4 J=1,N
C      S=B(I,J)
C      IF(I.EQ.1) GO TO 2
C      I1=I-1
C      DO 1 K=1,I1
C      S=S-B(I,K)*B(J,K)
C 1 IF(J.NE.I) GO TO 3
C 2 IF(S.LE.0.) GO TO 13
C T=SQRT(S)
C B(I,I)=T
C GO TO 4
C 3 B(J,I)=S/T
C 4 CONTINUE
C DO 6 I=1,N
C DO 6 J=1,N
C S=A(I,J)
C IF(I.EQ.1) GO TO 6
C I1=I-1
C DO 5 K=1,I1
C S=S-B(I,K)*A(J,K)
C 5 A(J,I)=S/B(I,I)
C DO 10 J=1,N
C DO 10 I=J,N
C S=A(I,J)
C IF(I.EQ.J) GO TO 8
C I1=I-1
C DO 7 K=J,I1
C S=S-A(K,J)*B(I,K)
C 7 IF(J.EQ.1) GO TO 10
C J1=J-1
C DO 9 K=1,J1
C S=S-A(J,K)*B(I,K)
C 9 A(I,J)=S/B(I,I)
C
C CALL TRE02(NA,NX,N,2.44E-03,A,LAMBDA,E,X)
C CALL TQL2(NX,N,2.22E-13,LAMBDA,E,X,(EPR)
C IF(IERR.NE.0) GO TO 13
C
```

```
AXL0007
AXL0008
AXL0009
AXL0010
AXL0011
AXL0012
AXL0013
AXL0014
AXL0015
AXL0016
AXL0017
AXL0018
AXL0019
AXL0020
AXL0021
AXL0022
AXL0023
AXL0024
AXL0025
AXL0026
AXL0027
AXL0028
AXL0029
AXL0030
AXL0031
AXL0032
AXL0033
AXL0034
AXL0035
AXL0036
AXL0037
AXL0038
AXL0039
AXL0040
AXL0041
AXL0042
AXL0043
AXL0044
AXL0045
AXL0046
AXL0047
AXL0048
AXL0049
AXL0050
AXL0051
AXL0052
AXL0053
AXL0054
AXL0055
AXL0056
AXL0057
AXL0058
AXL0059
```

```
DO 12 J=1,N
DO 12 IBACK=1,N
I=N+1-IBACK
S=X(I,J)
IF(I.EQ.N) GO TO 12
I1=I+1
DO 11 K=I1,N
S=S-B(K,I)*X(K,J)
11 X(I,J)=S/B(I,I)
12 X(I,J)=S/B(I,I)
C-----SCALE EIGENVECTORS TO MAKE LARGEST ENTRY = 1.0
DO 301 J=1,N
BIG=X(1,J)
DO 302 I=2,N
IF(ABS(X(I,J)).LE.ABS(BIG))GO TO 302
BIG=X(I,J)
302 CONTINUE
DO 303 I=1,N
303 X(I,J)=X(I,J)/BIG
301 CONTINUE
RETURN
13 RETURN JJ
END
```

```
C
C
```

```
AXL0060
AXL0061
AXL0062
AXL0063
AXL0064
AXL0065
AXL0066
AXL0067
AXL0068
AXL0069
AXL0070
AXL0071
AXL0072
AXL0073
AXL0074
AXL0075
AXL0076
AXL0077
AXL0078
AXL0079
AXL0080
AXL0081
AXL0082
AXL0083
```

```

*DECK TOL2
SUBROUTINE TOL2 (NM,N,MACHEP,D,E,Z,ERROR)
C
REAL MACHEP,D(N),E(N),Z(NM,N)
INTEGER ERROR
C
C-----THIS SUBPROGRAM IS A TRANSLATION OF THE ALGOL PROCEDURE TOL2,
C NUM.MATH.11, 293-306(1968) BY BOWLER, MARTIN, REINSCH, AND
C WILKINSON.
C
C THIS SUBPROGRAM USES QL TRANSFORMATIONS TO FIND THE
C EIGENVALUES AND EIGENVECTORS OF A TRIDIAGONAL MATRIX,
C GIVEN WITH ITS DIAGONAL ELEMENTS IN THE ARRAY D(N)
C AND ITS SUBDIAGONAL ELEMENTS IN THE LAST N-1 ELEMENTS
C OF THE ARRAY E(N).
C
C THE EIGENVALUES ARE OVERWRITTEN ON THE DIAGONAL ELEMENTS
C IN THE ARRAY D IN ASCENDING ORDER.
C
C THE EIGENVECTORS ARE FORMED IN THE ARRAY Z(N,N), OVERWRITING
C THE ACCUMULATED TRANSFORMATIONS AS SUPPLIED BY TREO2.
C
C IF THE TRIDIAGONAL MATRIX IS PRIMARY DATA (THUS, TREO2
C HAS NOT BEEN USED), Z SHOULD BE PRESET TO THE IDENTITY
C MATRIX.
C
C MACHEP IS THE RELATIVE MACHINE PRECISION. MACHEP SHOULD
C BE SET TO 2**(-52) FOR LONG FORM ARITHMETIC ON S/360.
C
C THE PROCEDURE FAILS (ERROR IS SET TO THE INDEX OF THE
C EIGENVALUE FOR WHICH FAILURE OCCURRED) IF ANY EIGENVALUE
C TAKES MORE THAN 30 ITERATIONS.
C
C NM MUST BE SET TO THE ROW DIMENSION OF TWO-DIMENSIONAL ARRAY
C PARAMETERS AS DECLARED IN THE CALLING PROGRAM DIMENSION STATEMENT.
C
C TRANSLATED BY V. KLEMA, ARGONNE NATIONAL LABORATORY, NOV., 1968.
C MODIFIED BY B. GARBOW, JAN. 1971.
C
C-----CALLED BY SUBPROGRAM "AXLBX"
C
C-----
C
C ERROR = 0
C IF (N.EQ. 1) GO TO 1001
C
C DO 100 I = 2, N
C 100 E(I-1) = E(I)
C
C F = 0.0
C R = 0.0
C E(N) = 0.0
C
C DO 240 L = 1, N
C J = 0
C H = MACHEP * (ABS(D(L)) + ABS(E(L)))
C IF (H.LT. H) R = H
C
C C=C LOOK FOR SMALL SUB-DIAGONAL ELEMENT

```

```

TQL00000
TQL00010
TQL00020
TQL00030
TQL00040
TQL00050
TQL00060
TQL00070
TQL00080
TQL00090
TQL00100
TQL00110
TQL00120
TQL00130
TQL00140
TQL00150
TQL00160
TQL00170
TQL00180
TQL00190
TQL00200
TQL00210
TQL00220
TQL00230
TQL00240
TQL00250
TQL00260
TQL00270
TQL00280
TQL00290
TQL00300
TQL00310
TQL00320
TQL00330
TQL00340
TQL00350
TQL00360
TQL00370
TQL00380
TQL00390
TQL00400
TQL00410
TQL00420
TQL00430
TQL00440
TQL00450
TQL00460
TQL00470
TQL00480
TQL00490
TQL00500
TQL00510
TQL00520
TQL00530
TQL00540
TQL00550
TQL00560
TQL00570
TQL00580
TQL00590
C
C DO 110 M = L, N
C IF (ABS(E(M)) .LE. B) GO TO 120
C 110 CONTINUE
C
C 120 IF (M.EQ. L) GO TO 220
C 130 IF (J.EQ. 30) GO TO 1000
C J = J + 1
C
C CCCCC FORM SHIFT
C C=C=C=C
C P = (D(L+1) - D(L)) / (2.0 * E(L))
C R = SQRT(P * P + 1.0)
C H = D(L) - E(L) / (P + SIGN(R,P))
C
C DO 140 I = L, N
C 140 D(I) = D(I) - H
C
C F = F + H
C
C CCCCC QL TRANSFORMATION
C C=C=C=C
C P = D(M)
C S = 1.0
C S = 0.0
C MML = M - L
C
C CCCCC FOR I=M-1 STEP -1 UNTIL L DO -- C=C=C=C=C=C=C=C
C DO 200 II = 1, MML
C I = M - II
C G = C * E(II)
C H = C * P
C IF (ABS(P) .LT. ABS(E(II))) GO TO 150
C C = E(II) / P
C R = SQRT(C * C + 1.0)
C E(II+1) = S * P * R
C S = C / R
C C = 1.0 / R
C GO TO 160
C 150 C = P / E(II)
C R = SQRT(C * C + 1.0)
C E(II+1) = S * E(II) * R
C S = 1.0 / R
C C = C / R
C 160 P = C * D(II) - S * G
C D(II+1) = H + S * (C * G + S * D(II))
C
C CCCCC FORM VECTOR C=C=C=C=C=C=C=C
C DO 180 K = 1, N
C H = Z(K,I+1)
C Z(K,I+1) = S * Z(K,I) + C * H
C Z(K,I) = C * Z(K,I) - S * H
C 180 CONTINUE
C
C 230 CONTINUE
C
C E(L) = S * P
C D(L) = C * P
C IF (ABS(E(L)) .GT. B) GO TO 130
C 220 D(L) = D(L) + F

```

```

TQL00600
TQL00610
TQL00620
TQL00630
TQL00640
TQL00650
TQL00660
TQL00670
TQL00680
TQL00690
TQL00700
TQL00710
TQL00720
TQL00730
TQL00740
TQL00750
TQL00760
TQL00770
TQL00780
TQL00790
TQL00800
TQL00810
TQL00820
TQL00830
TQL00840
TQL00850
TQL00860
TQL00870
TQL00880
TQL00890
TQL00900
TQL00910
TQL00920
TQL00930
TQL00940
TQL00950
TQL00960
TQL00970
TQL00980
TQL00990
TQL01000
TQL01010
TQL01020
TQL01030
TQL01040
TQL01050
TQL01060
TQL01070
TQL01080
TQL01090
TQL01100
TQL01110
TQL01120
TQL01130
TQL01140
TQL01150
TQL01160
TQL01170
TQL01180
TQL01190

```

```

247 CONTINUE
C
CCCCC: ORDER EIGENVALUES AND EIGENVECTORS
CC
      NM1 = N - 1
      DO 300 I = 1, NM1
        K = I
        P = 0(I)
        IP1 = I + 1
C
        DO 260 J = IP1, N
          IF (D(J) .LE. P) GO TO 260
          K = J
          P = D(J)
260    CONTINUE
C
        IF (K .EQ. I) GO TO 300
        D(K) = D(I)
        D(I) = P
C
        DO 280 J = 1, N
          P = Z(J,I)
          Z(J,I) = Z(J,K)
          Z(J,K) = P
280    CONTINUE
C
300 CONTINUE
C
      GO TO 1001
C
CCCCC: FAIL EXIT STATEMENT IS 1000
CCCCC
1000 ERROR = L
1001 RETURN
C
CCCCC: LAST CARD OF TQL2
CCCCC
      END
C

```

```

TQL01200
TQL01210
TQL01220
TQL01230
TQL01240
TQL01250
TQL01260
TQL01270
TQL01280
TQL01290
TQL01300
TQL01310
TQL01320
TQL01330
TQL01340
TQL01350
TQL01360
TQL01370
TQL01380
TQL01390
TQL01400
TQL01410
TQL01420
TQL01430
TQL01440
TQL01450
TQL01460
TQL01470
TQL01480
TQL01490
TQL01500
TQL01510
TQL01520
TQL01530
TQL01540
TQL01550
TQL01560
TQL01570
TQL01580
TQL01590

```

```

*DECK TRED2
SUBROUTINE TRED2 (NM1,NM2,N,TOL,A,D,E,Z)
C
      REAL A(NM1,N),D(N),E(N),Z(NM2,N)
C
C-----THIS SUBPROGRAM REDUCES THE GIVEN LOWER TRIANGLE OF A
C SYMMETRIC MATRIX STORED IN THE ARRAY A(N,N) TO
C TRIDIAGONAL FORM USING HOUSEHOLDER'S REDUCTION.
C THE DIAGONAL OF THE RESULT IS STORED IN THE ARRAY D(N)
C AND THE SUB-DIAGONAL IN THE LAST N-1 ELEMENTS OF THE ARRAY
C E(N) (WITH THE ADDITIONAL ELEMENT E(1)=0).
C THE TRANSFORMATION MATRICES ARE ACCUMULATED IN THE ARRAY Z(N,N).
C TOL IS A TOLERANCE FOR CHECKING IF THE TRANSFORMATION IS VALID.
C TOL SHOULD BE SET TO 10.**(-60) FOR LONG FORM ARITHMETIC ON S360.
C NM1 AND NM2 ARE THE ROW DIMENSIONS OF THE TWO DIMENSIONAL ARRAYS
C "A" AND "Z" AS DECLARED IN THE CALLING PROGRAM DIMENSION STATEMENT.
C
C THE ARRAY A IS LEFT UNALTERED UNLESS THE ACTUAL PARAMETERS
C CORRESPONDING TO A AND Z ARE IDENTICAL.
C THIS SUBPROGRAM IS A TRANSLATION OF THE ALGOL PROCEDURE TRED2,
C NUM. MATH. 11,181-195(1968) BY MARTIN, REINSCH, AND WILKINSON.
C
C-----CALLED BY SUBPROGRAM "AXLBX"
C
      DO 100 I=1,N
        DO 100 J=1,I
          100 Z(I,J)=A(I,J)
          IF (N.EQ.1) GO TO 1015
C
          FOR I= M STEP -1 UNTIL 2 DO --
            DO 1011 II=2,M
              I = M+2 - II
              L=I-2
              F = Z(I,I-1)
              G=0.0
              IF (L.LT.1) GO TO 103
              DO 102 K = 1,L
                102 G = G + Z(I,K)*Z(I,K)
              103 H = G + F*F
              IF G IS TOO SMALL FOR ORTHOGONALITY TO BE GUARANTEED, THE
              TRANSFORMATION IS SKIPPED,
              IF (G.GT. TOL) GO TO 104
              E(I)= F
              H = 0.0
              GOTO 101
            104 L=L+1
              G=SQRT (H)
              IF (F.GE.0.0 ) G=-G
              E(I) = G
              H = H-F*G
              Z(I,I-1) = F-G
              F = 0.0
              DO 105 J= 1,L
                Z(J,I) = Z(J,I)/H
              G=0.0
C
              FORM ELEMENT OF A*U
              DO 106 K = 1,J
                106 G = G + Z(J,K) * Z(I,K)
              JP1 = J+1
              IF (L .LT. JP1) GOTO 9108

```

```

TREL0000
TREL0010
TREL0020
TREL0030
TREL0040
TREL0050
TREL0060
TREL0070
TREL0080
TREL0090
TREL0100
TREL0110
TREL0120
TREL0130
TREL0140
TREL0150
TREL0160
TREL0170
TREL0180
TREL0190
TREL0200
TREL0210
TREL0220
TREL0230
TREL0240
TREL0250
TREL0260
TREL0270
TREL0280
TREL0290
TREL0300
TREL0310
TREL0320
TREL0330
TREL0340
TREL0350
TREL0360
TREL0370
TREL0380
TREL0390
TREL0400
TREL0410
TREL0420
TREL0430
TREL0440
TREL0450
TREL0460
TREL0470
TREL0480
TREL0490
TREL0500
TREL0510
TREL0520
TREL0530
TREL0540
TREL0550
TREL0560
TREL0570
TREL0580
TREL0590

```

```

      DO 108 K = JP1,L
      G = G + Z(K,J) * Z(I,K)
C     FORM ELEMENT OF P
9108 E(I) = G/H
      F = F + G * Z(I,J,I)
C     CONTINUE
      FORM K -- SEE ALGOL TREQ2
      HM = F/(H*H)
C     FORM REDUCED A
      DO 109 J = 1,L
      F = Z(I,J)
      E(J) = E(J) - HM * F
      G = E(J)
      DO 109 K = 1,J
      Z(I,J,K) = Z(I,J,K) - F * E(K) - G * Z(I,K)
C     101 CORRESPONDS TO SKIP
      101 D(I) = H
      1011 CONTINUE
      1015 D(I) = 0.3
      E(I) = 0.0
C     ACCUMULATION OF TRANSFORMATION MATRICES
      DO 201 I = 1,N
      L = I - 1
      IF (D(I) .EQ. 0.0) GO TO 202
      DO 203 J = 1,L
      G = 0.0
      DO 204 K = 1,L
      G = G + Z(I,K) * Z(I,K)
      DO 203 K = 1,L
      Z(K,J) = Z(K,J) - G * Z(K,I)
      202 D(I) = Z(I,I)
      Z(I,I) = 1.0
      IF (L .LT. 1) GO TO 201
      DO 205 J = 1,L
      Z(I,J) = 0.0
      205 Z(I,J) = 0.0
      201 CONTINUE
      RETURN
      END

```

```

TREQ0600
TREQ0610
TREQ0620
TREQ0630
TREQ0640
TREQ0650
TREQ0660
TREQ0670
TREQ0680
TREQ0690
TREQ0700
TREQ0710
TREQ0720
TREQ0730
TREQ0740
TREQ0750
TREQ0760
TREQ0770
TREQ0780
TREQ0790
TREQ0800
TREQ0810
TREQ0820
TREQ0830
TREQ0840
TREQ0850
TREQ0860
TREQ0870
TREQ0880
TREQ0890
TREQ0900
TREQ0910
TREQ0920
TREQ0930
TREQ0940
TREQ0950
TREQ0960
TREQ0970
TREQ0980

```

```

*DECK STATIC
SUBROUTINE STATIC
C
C-----SUBPROGRAM FOR STATIC DISPLACEMENT ANALYSIS OF PANEL (NFM=0) OR
C     FRAMED PANEL (NFM=1) STRUCTURES FOR 'NSLC' STATIC LOADINGS.
C     ACCEPTABLE LOADINGS INCLUDE CONCENTRATED LOADS 'AF' (FORCES,
C     MOMENTS) CORRESPONDING TO STRUCTURE F DISPLACEMENT TYPES AND/OR
C     STATIC MOVEMENTS 'OR' OF SUPPORT RESTRAINTS. RESULTS OF THE
C     ANALYSIS ARE THE DISPLACEMENTS 'OF' OF THE STRUCTURE DOF'S WHERE
C
      OF = (SFF)**-1 (AF - SFR * OR)
C
C     ALL PANEL AND FRAME NODE DISPLACEMENTS ARE DETERMINED BY BACK-
C     SUBSTITUTION IF (IREAC+ISTRES)=1. SUPPORT REACTIONS ARE COMPUTED
C     IF IREAC=1 AND PANEL STRESSES ARE COMPUTED IF ISTRES=1.
C**** NOTE: SEE EQUIVALENCE STATEMENT BELOW FOR PROPER DIMENSIONS OF
      THE FOLLOWING ARRAYS
C
      REAL TPFI(120,45),TPM(120,45),TMFS(120,45),TPR(120,56),TMRS(120,56)
      REAL DPL(810),DPR(810),AR(810),TEMP(810),RXP(100,198)
      REAL MX,MX,MY,MY
      INTEGER SN,PPV,FPV
C
      COMMON
      1/CBNPN/XP(25),YP(25),DBDM(25),DBFS(125),JFN(25,2),
      2/JPN(25),INTYP(25),IREP(16),ICNTY(16,4),KS2(25,5)
      3/CBNPFD/PPV(200),FPV(96),NFL(200),NRL(200),JPL(96)
      4/CBSTIF/SFF( 40, 40),SFP( 40, 120),SFH( 40, 45),SFR( 40, 56),
      5/SPP(120,120),SPM(120,45),SPR(120, 56),
      6/SHM( 45, 45),SHR( 45, 56),
      7/SRR( 56, 56)
      8/CBTSTOR/TS1(120,45),TS2(120, 56),TS3( 45, 120)
      9/CBMISC/AF(256),DR(256),DF(256),DP(256),DMI(44),DUM(20)
      10/CBDATA/EP,TP,PRP,DENSP,EF,CF,TF,DENSF,NS,SN,MCOOE,ISM,
      11/NRPN,NFT,NPT,NMT,NRT,NF,NP,NM,NR,NPN,NPCOE,NPD,NPDOF,
      12/NFJ,NFM,NJ6,NFNSPR,NPNSPR,ISDA,NPR,NFR,NFJ,
      13/NFRT,IREAC,ISTRES,ISTAT,IL,IKKK,TIME
      14/CBOYND/PLIST(300),INO,INW,IAF,IGA,ISO,MDLC,LN,NPTS,NMDOES,
      15>IDMAX,IDTH,NO,D1,D2,F0,F0,F0,F0,F0,F0,F0,F0,F0,F0,F0,F0,F0,F0,F0,F0,
      16/SXA,SYA,SZA,NTX,NTY,NTZ
      EQUIVALENCE (TPFI,TPM,TMFS,TS1),(TPR,TMRS,TS2),(RXP,SPPI)
      EQUIVALENCE (DPL,TS3(1,1)),(DPR,TS3(1,19)),(AF,TS3(1,37)),
      17/ (TEMP,TS3(1,55))
C
      IOOL=813
C
C-----NSLC = NO. OF STATIC LOADING CONDITIONS
C
C***** FOR DYNAMIC RESPONSE CALCULATIONS, GO TO 7(31)
      IF (ISOA.EQ.1) GO TO 7(31)
C*****
      READ(5,*) NSLC
      IF (NSLC.EQ.0) RETURN
C
      PRINT 950
950 FORMAT('STRUCTURE STIFFNESS MATRIX',/, '-----',
      1/ '-----',/)
      CALL MATPRT(NFT,NFT,SFF,NF,NF)
C

```

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STA0000
STA0010
STA0020
STA0030
STA0040
STA0050
STA0060
STA0070
STA0080
STA0090
STA0100
STA0110
STA0120
STA0130
STA0140
STA0150
STA0160
STA0170
STA0180
STA0190
STA0200
STA0210
STA0220
STA0230
STA0240
STA0250
STA0260
STA0270
STA0280
STA0290
STA0300
STA0310
STA0320
STA0330
STA0340
STA0350
STA0360
STA0370
STA0380
STA0390
STA0400
STA0410
STA0420
STA0430
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STA0500
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STA0570
STA0580
STA0590

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C      IL=LN=0
703 IL=IL+1
C-----CLEAR ARRAYS
C      DO 951 I=1,256
951 AF(I)=OR(I)=DF(I)=OP(I)=0.
      DO 952 I=1,144
952 OM(I)=0.
      DO 953 I=1,IDOL
953 OPL(I)=OFR(I)=AR(I)=TEMP(I)=0.
C-----NLN = NO. OF LOADED NODES( NPNSD = NO. OF PANEL NODES WITH SUPPORTS
C      DISPLACEMENTS SPECIFIED: NFJSD = NO. OF FRAME JOINTS WITH SUPPORT
C      QISPL. SPECIFIED: IREAC (=0, DO NOT COMPUTE SUPPORT REACTIONS: =1
C      COMPUTE REACTIONS); ISTRES (=0, DO NOT COMPUTE STRESSES: =1, COMPU
C      STRESSES)
C      READ(5,*) NLN,NPNSD,NFJSD,IREAC,ISTRES
      PRINT 702,NSLC
702 FORMAT("STATIC DISPLACEMENT ANALYSIS",/, " -----",
      " -----",/, " NSLC =",I3)
      PRINT 7021,IL,NLN,NPNSD,NFJSD,IREAC,ISTRES
7021 FORMAT("LOADING CONDITION NUMBER",I3,/, " NLN =",I3,
      " NPNSD =",I3, " NFJSD =",I3, " IREAC =",I3, " ISTRES =",I3)
C-----NODAL LOADS
C      IF (NLN.EQ.0) GO TO 704
      PRINT 700
700 FORMAT("NODAL LOADS",/, " NDOE      AF1      AF2      ",
      " AF3      AF4      AF5      AF6      ",
      " AF7      AF8")
      DO 705 I=1,NLN
      READ(5,*) K,TEMP(8*K-7),TEMP(8*K-6),TEMP(8*K-5),TEMP(8*K-4),
      TEMP(8*K-3),TEMP(8*K-2),TEMP(8*K-1),TEMP(8*K)
705 PRINT 706,K,TEMP(8*K-7),TEMP(8*K-6),TEMP(8*K-5),TEMP(8*K-4),
      TEMP(8*K-3),TEMP(8*K-2),TEMP(8*K-1),TEMP(8*K)
706 FORMAT(I4,1X,1P0E13.5)
C      DO 707 I=1,NPD
707 AF(PPV(I))=TEMP(I)
C-----SUPPORT DISPLACEMENTS
C      704 IF (NPNSD+NFJSD).EQ.0) GO TO 709
      IF (IL.EQ.1) GO TO 7101
      REPLACE SFR* (NFM=0) OR SFR* (NFM>0) FOR LOADINGS 2,....,NSLC
      RACKSPACE 33
      READ(33) ((SFR(I,J),J=1,NRT),I=1,NFT)
7101 DO 710 I=1,IDOL
710 TEMP(I)=0.
C      PRINT 711
711 FORMAT("APPLIED SUPPORT DISPLACEMENTS")
712 FORMAT(" JOINT      OPL      OR2      OR3      ",
      " OR4      OP5      OF6")
713 FORMAT("NODE      OR1      OR2      OR3      OR4",

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STA0360
STA04610
STA06620
STA06630
STA06640
STA06650
STA06660
STA06670
STA06680
STA06690
STA06700
STA06710
STA06720
STA0736
STA00740
STA00750
STA00760
STA00770
STA00780
STA00790
STA00800
STA00810
STA00820
STA00830
STA00840
STA00850
STA00860
STA00870
STA00880
STA00890
STA00900
STA00910
STA00920
STA00930
STA00940
STA00950
STA00960
STA00970
STA00980
STA00990
STA01000
STA01010
STA01020
STA01030
STA01040
STA01050
STA01060
STA01070
STA01080
STA01090
STA01100
STA01110
STA01120
STA01130
STA01140
STA01150
STA01160
STA01170
STA01180
STA01190

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```

C      # "      OR5      OR6      OR7      OR8")
C      IF (NFJSD.EQ.0) GO TO 7141
      PRINT 712
      DO 714 I=1,NFJSD
      READ(5,*) K,TEMP(6*K-5),TEMP(6*K-4),TEMP(6*K-3),TEMP(6*K-2),
      TEMP(6*K-1),TEMP(6*K)
      PRINT 716,K,TEMP(6*K-5),TEMP(6*K-4),TEMP(6*K-3),TEMP(6*K-2),
      TEMP(6*K-1),TEMP(6*K)
716 FORMAT(I5,1X,1P0E13.5)
714 CONTINUE
C      I1=NPR
      DO 720 I=1,NJ6
      IF (JRL(I).EQ.0) GO TO 720
      I1=I+1
      DR(I1)=TEMP(I)
720 CONTINUE
C      7141 IF (NPNSD.EQ.0) GO TO 719
      PRINT 713
      DO 717 I=1,IDOL
717 TEMP(I)=0.
      DO 715 I=1,NPNSD
      READ(5,*) K,TEMP(8*K-7),TEMP(8*K-6),TEMP(8*K-5),TEMP(8*K-4),
      TEMP(8*K-3),TEMP(8*K-2),TEMP(8*K-1),TEMP(8*K)
      PRINT 706,K,TEMP(8*K-7),TEMP(8*K-6),TEMP(8*K-5),TEMP(8*K-4),
      TEMP(8*K-3),TEMP(8*K-2),TEMP(8*K-1),TEMP(8*K)
715 CONTINUE
C      I1=0
      DO 718 I=1,NPD
      IF (NRL(I).EQ.0) GO TO 718
      I1=I+1
      DR(I1)=TEMP(I)
718 CONTINUE
C-----COMPUTE AF=SFR*DR
C      719 DO 721 I=1,NFT
      SUM=AF(I)
      DO 722 J=1,NRT
722 SUM=SUM+SFR(I,J)*OR(J)
721 AF(I)=SUM
C-----COMPUTE OF = (SFF)**-1 (AF = SFR * OR)
C      709 DO 723 I=1,NFT
      SFP(I,1)=AF(I)
      DO 723 J=1,NFT
723 SPP(I,J)=SFF(I,J)
C      CALL OCOMP(NFT,SPP,NP,NP),RETURNS(724)
      GO TO 726
724 PRINT 725
725 FORMAT("DECOMPOSITION OF SFF FAILS IN STATIC")
      RETURN
C      726 CALL SOLVE(NFT,SPP,SFP,SFM,1,1,1,NP,NP,NF,NF,NM)
C

```

```

STA01200
STA01210
STA01220
STA01230
STA01240
STA01250
STA01260
STA01270
STA01280
STA01290
STA01300
STA01310
STA01320
STA01330
STA01340
STA01350
STA01360
STA01370
STA01380
STA01390
STA01400
STA01410
STA01420
STA01430
STA01440
STA01450
STA01460
STA01470
STA01480
STA01490
STA01500
STA01510
STA01520
STA01530
STA01540
STA01550
STA01560
STA01570
STA01580
STA01590
STA01600
STA01610
STA01620
STA01630
STA01640
STA01650
STA01660
STA01670
STA01680
STA01690
STA01700
STA01710
STA01720
STA01730
STA01740
STA01750
STA01760
STA01770
STA01780
STA01790

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      DO 727 I=1,NFI
      727 OF(I)=SFH(I,1)
C
C-----PRINT STATIC DISPLACEMENTS FOR STRUCTURE DOF'S
C***** FOR DYN. RESPONSE CALCS., COMPUTE PANEL STRESSES AND REACS. (ISO
C
      IF (ISDA.EQ.0) GO TO 7032
      7031 DO 7033 I=1,NFI
      7033 OF(I)=RXP(I,ISTAT)
      LL1=NFI+1
      I1=0
      DO 7034 I=LL1,NFRT
      I1=I1+1
      7034 OR(I1)=RXP(I,ISTAT)
      IF (IL.EQ.1) GO TO 7035
      REWIND 33
      READ(33)
      READ(33) ((SFR(I,J),J=1,NRT),I=1,NFI)
      GO TO 7035
C
      7032 IC=0
      PRINT 7271
      7271 FORMAT('STATIC DISPLACEMENTS - STRUCTURE DOF'S ONLY',/,
      1 " NODE      01      02      03      04",
      2 " "      05      06      07      08")
      DO 728 I=1,NPH
      PRINT 730,I
      730 FORMAT(I4,
      1 I=0,I=7
      2 I=8,I=15
      3 KC=-12,K
      4 DO 729 J=I1,I8
      5 KC=KC+13
      6 IF (NFI(J).EQ.0) GO TO 729
      7 IC=IC+1
      8 PRINT 731,KC,OF(IC)
      9 731 FORMAT(' ',X,1PE13.5)
      10 729 CONTINUE
      11 728 CONTINUE
C
      7035 IF ((IREAC.EQ.0).AND.(ISTRES.EQ.0)) GO TO 955
C
C-----COMPUTE FRAME AND PANEL DISPLACEMENTS
C
      IF (NFM.GT.0) GO TO 600
      COMPUTE DP(I) FOR NFM=0 CASE
      REWIND 10
      READ(10) ((TPF(I,J),J=1,NFI),I=1,NPF)
      READ(10) ((TPR(I,J),J=1,NRT),I=1,NPT)
C
      DO 601 I=1,NPT
      SUM=0.
      DO 602 J=1,NFI
      602 SUM=SUM-TPF(I,J)*OF(J)
      DO 603 K=1,NRT
      603 SUM=SUM-TPR(I,K)*OR(K)
      601 DP(I)=SUM
      GO TO 607
C
      COMPUTE DM(I) AND DP(I) FOR NFM>0 CASE

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```

      STA01800
      STAC1810
      STAP1820
      STA01830
      STA01840
      STA01850
      STA01860
      STAC1870
      STAC1880
      STAC1890
      STAC1900
      STAC1910
      STA01920
      STA01930
      STA01940
      STA01950
      STA01960
      STA01970
      STA01980
      STA01990
      STA02000
      STAC2010
      STA02020
      STA02030
      STA02040
      STA02050
      STAC2060
      STA02070
      STAC2080
      STAC2090
      STAC2100
      STA02110
      STA02120
      STA02130
      STA02140
      STA02150
      STAC2160
      STA02170
      STAC2180
      STA02190
      STA02200
      STA02210
      STA02220
      STA02230
      STAC2240
      STA02250
      STAC2260
      STAC2270
      STA02280
      STAC2290
      STA02300
      STA02310
      STAC2320
      STAC2330
      STAC2340
      STA02350
      STAC2360
      STAC2370
      STAC2380
      STA02390

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      600 IF ((IL.GT.1).OR.(LN.GT.1)) GO TO 6033
      BACKSPACE 10
      BACKSPACE 10
      6033 READ(10) ((TMFS(I,J),J=1,NFI),I=1,NMT)
      READ(10) ((TMRS(I,J),J=1,NRT),I=1,NMT)
      DO 610 I=1,NMT
      SUM=0.
      DO 611 J=1,NFI
      611 SUM=SUM-TMFS(I,J)*OF(J)
      DO 612 K=1,NRT
      612 SUM=SUM-TMRS(I,K)*OR(K)
      610 DM(I)=SUM
C
      REWIND 10
      READ(10) ((TPF(I,J),J=1,NFI),I=1,NPT)
      READ(10) ((TPR(I,J),J=1,NRT),I=1,NPT)
C
      DO 613 I=1,NPT
      SUM=0.
      DO 614 J=1,NFI
      614 SUM=SUM-TPF(I,J)*OF(J)
      DO 615 K=1,NRT
      615 SUM=SUM-TPR(I,K)*OR(K)
      613 DP(I)=SUM
      READ(10) ((TPH(I,J),J=1,NMT),I=1,NPT)
      DO 616 I=1,NPT
      SUM=DP(I)
      DO 617 J=1,NMT
      617 SUM=SUM-TPH(I,J)*DM(J)
      616 DP(I)=SUM
C
      FORM DISPLACEMENT VECTOR FOR PANEL FROM OF, DP, AND OR
C
C
      607 IF (IP=IR=0)
      DO 604 I=1,NPD
      IF (NRL(I).EQ.1) GO TO 605
      IF (NFI(I).EQ.1) GO TO 606
      IP=IP+1
      OPL(I)=DP(IP)
      GO TO 604
      605 IR=IR+1
      OPL(I)=OR(IR)
      GO TO 604
      606 IF (IF=1)
      OPL(I)=OF(IF)
      604 CONTINUE
C
      PRINT 608
      608 FORMAT('0 PANEL DISPLACEMENTS - SUMMARY',/, " NOOE      01",
      1 " "      02      03      04      05",
      2 " "      06      07      08")
      PRINT 609, (I, (DPL(J+0*(I-1)),J=1,N),I=1,NPN)
      609 FORMAT(I4,1X,1PE13.5)
C
      IF (NFM.EQ.0) GO TO 500
C
      FORM DISPLACEMENT VECTOR OFR FOR FRAME USING DM, OR (FOR NFM>0)
C
      IM=J

```

```

      STAC2400
      STAC2410
      STAC2420
      STAC2430
      STAC2440
      STAC2450
      STAC2460
      STAC2470
      STAC2480
      STAC2490
      STAC2500
      STAC2510
      STAC2520
      STAC2530
      STAC2540
      STAC2550
      STAC2560
      STAC2570
      STAC2580
      STAC2590
      STAC2600
      STAC2610
      STAC2620
      STAC2630
      STAC2640
      STAC2650
      STAC2660
      STAC2670
      STAC2680
      STAC2690
      STAC2700
      STAC2710
      STAC2720
      STAC2730
      STAC2740
      STAC2750
      STAC2760
      STAC2770
      STAC2780
      STAC2790
      STAC2800
      STAC2810
      STAC2820
      STAC2830
      STAC2840
      STAC2850
      STAC2860
      STAC2870
      STAC2880
      STAC2890
      STAC2900
      STAC2910
      STAC2920
      STAC2930
      STAC2940
      STAC2950
      STAC2960
      STAC2970
      STAC2980
      STAC2990

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```

IR=NPR
DO 618 I=1,NJ6
IF(JRL(I).EQ.0)GO TO 619
IM=IM+1
OFR(I)=OM(IM)
GO TO 618
619 IR=IR+1
OFR(I)=OR(IR)
618 CONTINUE
C
PRINT 620
620 FORMAT("DFRAME DISPLACEMENTS",/, " JOINT      01      02",
/ "      03      04      05      06")
PRINT 621, (I, (OFR(J*6*(I-1)), J=1,6), I=1,NFJ)
621 FORMAT(I5,1X,1P6E13.5)
C
C-----COMPUTE SUPPORT REACTIONS (IF IREAC > 0)
C
500 IF(IREAC.EQ.0)GO TO 900
READ IN UNCONDENSED ARRAYS FOR SUPPORT COUPLING (NOTE THAT SFR*
(NFM=0) OR SFR** (NFM>0) IS DESTROYED)
C
IKKK=1
REWIND 11
READ(11) ((SFR(I,J),J=1,NRT),I=1,NFT)
READ(11) ((SPR(I,J),J=1,NRT),I=1,NPT)
READ(11) ((SRR(I,J),J=1,NRT),I=1,NRT)
DO 501 I=1,NRT
SUM=0.
DO 502 J1=1,NRT
502 SUM=SUM+SRR(I,J1)*OR(J1)
DO 503 J2=1,NPT
503 SUM=SUM+SPR(J2,I)*OF(J2)
DO 504 J3=1,NPT
504 SUM=SUM+SPR(J3,I)*OP(J3)
501 AR(I)=SUM
C
IF(NFM.EQ.0)GO TO 505
C
READ(11) ((SMR(I,J),J=1,NRT),I=1,NMT)
DO 506 I=1,NMT
SUM=AR(I)
DO 507 J=1,NMT
507 SUM=SUM+SMR(J,I)*OM(J)
506 AR(I)=SUM
C
505 IF(ISDA.EQ.1)PRINT 5051
5051 FORMAT("DNOTE - REACTIONS PRESENTED BELOW DO NOT INCLUDE",
/ " " INERTIA EFFECTS")
IF(INRPN.EQ.0)GO TO 520
PRINT 520
520 FORMAT("D PANEL SUPPORT REACTIONS",/, " NODE      AR1      AR2",
/ "      AR2      AR3      AF4      AR5      AR6",
/ "      AP7      AR2")
IC=0
DO 509 I=1,NPN
PRINT 73C,I
I1=8*I-7
I8=I1+7
KC=-12+4

```

```

STA.3007
STA.3010
STA.3020
STA.3030
STA.3040
STA.3050
STA.3060
STA.3070
STA.3080
STA.3090
STA.3100
STA.3110
STA.3120
STA.3130
STA.3140
STA.3150
STA.3160
STA.3170
STA.3180
STA.3190
STA.3200
STA.3210
STA.3220
STA.3230
STA.3240
STA.3250
STA.3260
STA.3270
STA.3280
STA.3290
STA.3300
STA.3310
STA.3320
STA.3330
STA.3340
STA.3350
STA.3360
STA.3370
STA.3380
STA.3390
STA.3400
STA.3410
STA.3420
STA.3430
STA.3440
STA.3450
STA.3460
STA.3470
STA.3480
STA.3490
STA.3500
STA.3510
STA.3520
STA.3530
STA.3540
STA.3550
STA.3560
STA.3570
STA.3580
STA.3590

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```

DO 513 J=I1,I8
KC=KC+13
IF(NRL(J).EQ.0)GO TO 511
IC=IC+1
PRINT 731,KC,AR(IC)
510 CONTINUE
509 CONTINUE
C
C
520 IF(INR.EQ.0)GO TO 900
PRINT 521
521 FORMAT("DFRAME SUPPORT REACTIONS",/, " JOINT      AR1      ",
/ "      AR2      AR3      AP4      AR5      AR6")
IC=NPR
DO 522 I=1,NFJ
PRINT 523,I
523 FORMAT(I5)
I1=6*I-5
I6=I1+5
KC=-12+5
DO 524 J=I1,I6
KC=KC+13
IF(JRL(J).EQ.0)GO TO 524
IC=IC+1
PRINT 731,KC,AR(IC)
524 CONTINUE
522 CONTINUE
C
C-----COMPUTE PANEL STRESSES (IF ISTRES > 0)
C
900 IF(ISTRES.EQ.0)GO TO 955
PRINT 957
957 FORMAT("ELEMEN STRESSES AND STRESS RESULTANTS",/,
/ " " "-----")
C
DO 958 IE=1,NPE
C
PRINT 959,IE
959 FORMAT("ELEMEN NO.",I3,/, " NODE      SX      ",
/ "      SY      SXY      MX      MY      MXY")
C
C--- STRESSES SX, SY, SXY FOR REFINED PLANE STRESS (BOW) ELEMENT AND
C STRESS RESULTANTS MX, MY, MXY FOR PLATE BENDING (BFS) ELEMENT
C
A1=XP(ICNTY(IE,2))-XP(ICNTY(IE,1))
A1=YP(ICNTY(IE,3))-YP(ICNTY(IE,2))
A2=A1/2.
A2=A1/2.
C
C--- FORM DISPL. VECTORS (BOW, BFS) FOR DOF'S AT CGM NOOFS OF ELEM.
C
IE
C
I1=-4
DO 960 J=1,4
I1=I1+4
J1=8*ICNTY(IE,J)-7
OBOW(I1+1)=OPL(IJ1)
OBOW(I1+2)=OPL(IJ1+1)
OBOW(I1+3)=OPL(IJ1+5)

```

```

STA.3600
STA.3610
STA.3620
STA.3630
STA.3640
STA.3650
STA.3660
STA.3670
STA.3680
STA.3690
STA.3700
STA.3710
STA.3720
STA.3730
STA.3740
STA.3750
STA.3760
STA.3770
STA.3780
STA.3790
STA.3800
STA.3810
STA.3820
STA.3830
STA.3840
STA.3850
STA.3860
STA.3870
STA.3880
STA.3890
STA.3900
STA.3910
STA.3920
STA.3930
STA.3940
STA.3950
STA.3960
STA.3970
STA.3980
STA.3990
STA.4000
STA.4010
STA.4020
STA.4030
STA.4040
STA.4050
STA.4060
STA.4070
STA.4080
STA.4090
STA.4100
STA.4110
STA.4120
STA.4130
STA.4140
STA.4150
STA.4160
STA.4170
STA.4180
STA.4190

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```

      DBOW(I1+4)=DPL(J1+6)
      DBFS(I1+1)=DPL(J1+2)
      DBFS(I1+2)=DPL(J1+3)
      DBFS(I1+3)=DPL(J1+4)
960 DBFS(I1+4)=DPL(J1+7)
C
      CALL BOWSTR(EP,TP,PRP,A1,B1,0.,0.,DBOW, SX,SY, SXY)
      CALL BFSSTR(EP,TP,PRP,A1,B1,0.,0.,DBFS, MX,MY, MXY)
      PRINT 961,ICNTY(IE,1),SX,SY,SXY,MX,MY,MXY
961 FORMAT(I4,1X,1P6E13.5)
      CALL BOWSTR(EP,TP,PRP,A1,B1,A1,0.,DBOW, SX,SY, SXY)
      CALL BFSSTR(EP,TP,PRP,A1,B1,A1,0.,DBFS, MX,MY, MXY)
      PRINT 961,ICNTY(IE,2),SX,SY,SXY,MX,MY,MXY
      CALL BOWSTR(EP,TP,PRP,A1,B1,A1,B1,DBOW, SX,SY, SXY)
      CALL BFSSTR(EP,TP,PRP,A1,B1,A1,B1,DBFS, MX,MY, MXY)
      PRINT 961,ICNTY(IE,3),SX,SY,SXY,MX,MY,MXY
      CALL BOWSTR(EP,TP,PRP,A1,B1,0.,B1,DBOW, SX,SY, SXY)
      CALL BFSSTR(EP,TP,PRP,A1,B1,0.,B1,DBFS, MX,MY, MXY)
      PRINT 961,ICNTY(IE,4),SX,SY,SXY,MX,MY,MXY
      CALL BOWSTR(EP,TP,PRP,A1,B1,AZ,B2,DBOW, SX,SY, SXY)
      CALL BFSSTR(EP,TP,PRP,A1,B1,AZ,B2,DBFS, MX,MY, MXY)
      PRINT 962,SX,SY,SXY,MX,MY,MXY
962 FORMAT(' CNTR',1P6E13.5)
C
956 CONTINUE
C
955 IF(ISDA.EQ.1)RETURN
C
      IF(IL.LT.NSLC)GO TO 703
      RETURN
      END

```

```

STA04204
STA04210
STA04220
STA04230
STA04240
STA04250
STA04260
STA04270
STA04280
STA04290
STA04300
STA04310
STA04320
STA04330
STA04340
STA04350
STA04360
STA04370
STA04380
STA04390
STA04400
STA04410
STA04420
STA04430
STA04440
STA04450
STA04460
STA04470
STA04480
STA04490
STA04500
STA04510

```

```

*DECK MATPRT
      SUBROUTINE MATPRT(M,N,A,MAX1,MAX2)
C
C-----THIS SUBPROGRAM PRINTS AN M BY N MATRIX 9 COLUMNS AT A TIME
C
C*****MAX1 AND MAX2 ARE THE DIMENSIONS OF THE ARRAY "A" IN THE CALLING PROGRAM.
C
C-----CALLED BY MAIN PROGRAM
C
      REAL A(MAX1,MAX2)
      INTEGER RTCOL
C
601 FORMAT(' ",I3,1X,1P9E13.5)
602 FORMAT(' "OCOLUMNS",I4,3X,9(I10,3X))
603 FORMAT(' "-----",I," ROW")
      NPAGES = (N-1)/9 + 1
      DO 101 I=1,NPAGES
        LTCOL = 9*(I-1) + 1
        RTCOL = 9*I
        IF (RTCOL.GT.N) RTCOL=N
        PRINT 602,(K,K=LTCOL,RTCOL)
        PRINT 603
        DO 101 J=1,M
          101 PRINT 601,J,(A(J,K),K=LTCOL,RTCOL)
        RETURN
      END

```

```

MAT 0007
MAT0010
MAT0020
MAT0030
MAT0040
MAT0050
MAT0060
MAT0070
MAT0080
MAT0090
MAT0100
MAT0110
MAT0120
MAT0130
MAT0140
MAT0150
MAT0160
MAT0170
MAT0180
MAT0190
MAT0200
MAT0210
MAT0220
MAT0230
MAT0240
MAT0250
MAT0260
MAT0270

```



```

*DECK DCOMP                                DCD00000
SUBROUTINE DCOMP (N,A,MAX1,MAX2),RETURN(JBJ) DCD00010
C                                           DCD00020
C-----THIS SUBPROGRAM DECOMPOSES THE SYMMETRIC, POSITIVE-DEFINITE MATRIX DCD00030
C A USING THE CHOLESKY SQUARE ROOT METHOD AS DISCUSSED ON PAGE 56 OF DCD00040
C 'COMPUTER PROGRAMS FOR STRUCTURAL ANALYSIS,' (1967) BY WM. WEAVER, JR., DCD00050
C JR., MATRIX A IS MXN, BUT MAY BE DIMENSIONED FOR MAX1 X MAX2 IN DCD00060
C HE DCD00070
C CALLING PROGRAM. DCD00080
C DCD00090
C DCD00100
C DIMENSION A(MAX1,MAX2) DCD00110
C DCD00120
C DO 101 I=1,N DCD00130
C DO 101 J=I,N DCD00140
C SUM=A(I,J) DCD00150
C K1=I-1 DCD00160
C IF(I.EQ.1) GO TO 1 DCD00170
C DO 102 K=1,K1 DCD00180
C 102 SUM=SUM-A(K,I)*A(K,J) DCD00190
C IF(J.NE.I) GO TO 2 DCD00200
C IF(SUM.LE.0.) RETURN JBJ DCD00210
C TEMP=1.0/SQRT(SUM) DCD00220
C A(I,J)=TEMP DCD00230
C GO TO 101 DCD00240
C 2 A(I,J)=SUM*TEMP DCD00250
C 101 CONTINUE DCD00260
C DCD00270
C RETURN DCD00280
C END DCD00290

```

```

*DECK SOLVE                                SOL00000
SUBROUTINE SOLVE (N,U,B,X,II,IC,NN,MU1,MU2,MB1,MB2,MX1,MX2) SOL00010
C                                           SOL00020
C-----THIS SUBPROGRAM SOLVES THE SYSTEM OF SIMULTANEOUS EQUATIONS, SOL00030
C A * X = B, IN A TWO-STEP PROCESS AS DISCUSSED ON PAGE 57 OF SOL00040
C 'COMPUTER PROGRAMS FOR STRUCTURAL ANALYSIS' BY WM. WEAVER, JR., SOL00050
C (1967). SUBPROGRAM 'DCOMP' PROVIDES THE DECOMPOSED VERSION OF SOL00060
C MATRIX A (STORED IN UPPER TRIANGULAR MATRIX U) NEEDED IN THE SOL00070
C SOLUTION PROCESS. X IS THE VECTOR (OR MATRIX) OF UNKNOWNNS AND SOL00080
C B IS THE VECTOR (OR MATRIX) OF KNOWN QUANTITIES. SOL00090
C SUBPROGRAM 'SOLVE' IS MODIFIED HERE SO THAT IT CAN BE USED SOL00100
C AS THE OBJECT OF A DO STATEMENT WHERE II IS THE DO VARIABLE. SOL00110
C THUS, THE MATRIX OF UNKNOWNNS, X(N,NN), CAN BE OBTAINED BY CALLING SOL00120
C 'SOLVE' NN TIMES. SOL00130
C NOTE: IF IC=1, MATRIX B IS USED AS IS; IF IC=2, THE TRANSPOSE SOL00140
C OF MATRIX B IS USED. SOL00150
C-----ARRAY SIZES: U(N X N), B(N X NN), X(N X NN) SOL00160
C C*****MU1,...,MX2 ARE THE DIMENSIONS OF ARRAYS U, B, AND X SOL00170
C IN THE CALLING PROGRAM. SOL00180
C SOL00190
C SOL00200
C SOL00210
C SOL00220
C DIMENSION U(MU1,MU2),B(MB1,MB2),X(MX1,MX2) SOL00230
C SOL00240
C DO 101 I=1,N SOL00250
C SUM=B(I,II) SOL00260
C IF(IC.EQ.2) SUM=B(II,I) SOL00270
C K1=I-1 SOL00280
C IF(I.EQ.1) GO TO 101 SOL00290
C DO 102 K=1,K1 SOL00300
C 102 SUM=SUM-U(K,I)*X(K,II) SOL00310
C 101 X(I,II)=SUM*U(I,I) SOL00320
C SOL00330
C DO 103 II=1,N SOL00340
C I=N-II+1 SOL00350
C SUM=X(I,II) SOL00360
C K2=I+1 SOL00370
C IF(II.EQ.N) GO TO 103 SOL00380
C DO 104 K=K2,N SOL00390
C 104 SUM=SUM-U(I,K)*X(K,II) SOL00400
C 103 X(I,II)=SUM*U(I,I) SOL00410
C SOL00420
C RETURN SOL00430
C END SOL00440

```

\*DECK SBOG  
SUBROUTINE SBOG (A1, B1, TT, E, U, SBI)

C  
C-----THIS SUBPROGRAM GENERATES THE ELEMENT STIFFNESS MATRIX FOR  
C THE 16 DEGREE-OF-FREEDOM BOGNER-FOX-SCHMIT TRIANGULAR PLATE-  
C BENDING ELEMENT (REFERRED TO AS THE BOGNER RECTANGLE ).  
C HOWEVER, THE SIGN CONVENTION AND NODAL NUMBERING  
C SEQUENCE HAS BEEN CHANGED FROM THAT USED BY BOGNER, ET AL.

C-----CALLED BY SUBPROGRAM PANLSM

C  
C  
C REAL SBI(136)  
C  
C D = E\*TT\*\*3/(12.0\*(1.0 - U\*U))  
C U = 0/(A1\*B1)  
C A2 = A1\*A1  
C B2 = B1\*B1  
C A = A2/B2  
C B = B2/A2

SBI(1)=D\*(+156./ 35.\*B+156./ 35.\*A 72./ 25.  
SBI(2)=D\*( 22./ 35.\*B+ 78./ 35.\*A 6./ 25.+ 6./ 5.\*U)\*B1  
SBI(3)=D\*( 78./ 35.\*B- 22./ 35.\*A 6./ 25.- 6./ 5.\*U)\*A1  
SBI(4)=D\*( -11./ 35.\*B- 11./ 35.\*A 1./ 50.- 1./ 5.\*U)\*A1\*B1  
SBI(5)=D\*( -156./ 35.\*B+ 54./ 35.\*A 72./ 25.  
SBI(6)=D\*( 22./ 35.\*B+ 27./ 35.\*A 6./ 25.- 6./ 5.\*U)\*B1  
SBI(7)=D\*( 78./ 35.\*B+ 13./ 35.\*A 6./ 25. )\*A1  
SBI(8)=D\*( -11./ 35.\*B+ 13./ 70.\*A 1./ 50.- 1./ 10.\*U)\*A1\*B1  
SBI(9)=D\*( -54./ 35.\*B+ 54./ 35.\*A 72./ 25.  
SBI(10)=D\*( 13./ 35.\*B+ 27./ 35.\*A 6./ 25. )\*B1  
SBI(11)=D\*( -27./ 35.\*B- 13./ 35.\*A 6./ 25. )\*A1  
SBI(12)=D\*(+ 13./ 70.\*B+ 13./ 70.\*A 1./ 50.  
SBI(13)=D\*( 54./ 35.\*B-156./ 35.\*A 72./ 25.  
SBI(14)=D\*( -13./ 35.\*B+ 78./ 35.\*A 6./ 25. )\*B1  
SBI(15)=D\*( -27./ 35.\*B+ 22./ 35.\*A 6./ 25.+ 6./ 5.\*U)\*A1  
SBI(16)=D\*(+ 13./ 70.\*B- 11./ 35.\*A 1./ 50.- 1./ 10.\*U)\*A1\*B1  
SBI(17)=D\*( 4./ 35.\*B+ 52./ 35.\*A 8./ 25. )\*B2  
SBI(18)=D\*( -11./ 35.\*B- 11./ 35.\*A 1./ 50.- 6./ 5.\*U)\*A1\*B1  
SBI(19)=D\*( - 2./ 35.\*B- 22./ 105.\*A 2./ 75.- 2./ 15.\*U)\*A1\*B2  
SBI(20)=D\*( -22./ 35.\*B+ 27./ 35.\*A 6./ 25.- 6./ 5.\*U)\*B1  
SBI(21)=D\*( - 4./ 35.\*B+ 18./ 35.\*A 8./ 25. )\*B2  
SBI(22)=D\*( -11./ 35.\*B+ 13./ 70.\*A 1./ 50.- 1./ 10.\*U)\*A1\*B1  
SBI(23)=D\*( - 2./ 35.\*B+ 13./ 105.\*A 2./ 75. )\*A1\*B2  
SBI(24)=D\*( -13./ 35.\*B+ 27./ 35.\*A 6./ 25. )\*B1  
SBI(25)=D\*( - 3./ 35.\*B+ 9./ 35.\*A 2./ 25. )\*B2  
SBI(26)=D\*( -13./ 70.\*B- 13./ 70.\*A 1./ 50.  
SBI(27)=D\*(+ 3./ 70.\*B+ 13./ 210.\*A 1./ 150.  
SBI(28)=D\*( 13./ 35.\*B- 78./ 35.\*A 6./ 25.  
SBI(29)=D\*( - 3./ 35.\*B+ 26./ 35.\*A 2./ 25.  
SBI(30)=D\*( -13./ 70.\*B+ 11./ 35.\*A 1./ 50.+ 1./ 10.\*U)\*A1\*B1  
SBI(31)=D\*(+ 3./ 70.\*B- 11./ 105.\*A 1./ 150.+ 1./ 30.\*U)\*A1\*B2  
SBI(32)=D\*(+ 52./ 35.\*B+ 4./ 35.\*A 8./ 25. )\*A2  
SBI(33)=D\*(+ 22./ 105.\*B+ 2./ 35.\*A 2./ 75.+ 2./ 15.\*U)\*A2\*B1  
SBI(34)=D\*(+ 78./ 35.\*B- 13./ 35.\*A 6./ 25. )\*A1  
SBI(35)=D\*(+ 11./ 35.\*B- 13./ 71.\*A 1./ 50.+ 1./ 10.\*U)\*A1\*B1  
SBI(36)=D\*(+ 26./ 35.\*B- 3./ 35.\*A 2./ 25. )\*A2  
SBI(37)=D\*(+ 11./ 105.\*B- 3./ 71.\*A 1./ 150.- 1./ 30.\*U)\*A2\*B1  
SBI(38)=D\*(+ 27./ 35.\*B+ 13./ 35.\*A 6./ 25. )\*A1  
SBI(39)=D\*( -13./ 70.\*B- 13./ 71.\*A 1./ 50.  
SBI(40)=D\*(+ 9./ 35.\*B+ 3./ 35.\*A 2./ 25. )\*A2

SBI(41)=D\*( -13./ 210.\*B- 3./ 70.\*A 1./ 150.  
SBI(42)=D\*( -27./ 35.\*B+ 22./ 35.\*A 6./ 25.+ 6./ 5.\*U)\*A1  
SBI(43)=D\*(+ 13./ 70.\*B- 11./ 35.\*A 1./ 50.- 1./ 10.\*U)\*A1\*B1  
SBI(44)=D\*(+ 18./ 35.\*B+ 4./ 35.\*A 8./ 25. )\*A2  
SBI(45)=D\*( -13./ 105.\*B+ 2./ 35.\*A 2./ 75.  
SBI(46)=D\*(+ 4./ 105.\*B+ 4./ 105.\*A 8./ 225.  
SBI(47)=D\*(+ 11./ 35.\*B- 13./ 70.\*A 1./ 50.+ 1./ 10.\*U)\*A1\*B1  
SBI(48)=D\*(+ 2./ 35.\*B- 13./ 105.\*A 2./ 75.  
SBI(49)=D\*(+ 11./ 105.\*B- 3./ 70.\*A 1./ 150.- 1./ 30.\*U)\*A1\*B1  
SBI(50)=D\*(+ 2./ 105.\*B- 1./ 35.\*A 2./ 225.  
SBI(51)=D\*(+ 13./ 70.\*B+ 13./ 70.\*A 1./ 50.  
SBI(52)=D\*( - 3./ 70.\*B- 13./ 210.\*A 1./ 150.  
SBI(53)=D\*(+ 13./ 210.\*B+ 3./ 70.\*A 1./ 150.  
SBI(54)=D\*( - 1./ 70.\*B- 1./ 70.\*A 1./ 450.  
SBI(55)=D\*( -13./ 70.\*B+ 11./ 35.\*A 1./ 50.+ 1./ 10.\*U)\*A1\*B1  
SBI(56)=D\*(+ 3./ 70.\*B- 11./ 105.\*A 1./ 150.+ 1./ 30.\*U)\*A1\*B2  
SBI(57)=D\*(+ 13./ 105.\*B- 2./ 35.\*A 2./ 75.  
SBI(58)=D\*( - 1./ 35.\*B+ 2./ 105.\*A 2./ 225.  
SBI(59)=D\*( 156./ 35.\*B+156./ 35.\*A 72./ 25.  
SBI(60)=D\*(+ 22./ 35.\*B+ 78./ 35.\*A 6./ 25.+ 6./ 5.\*U)\*B1  
SBI(61)=D\*(+ 78./ 35.\*B+ 22./ 35.\*A 6./ 25.+ 6./ 5.\*U)\*A1  
SBI(62)=D\*(+ 11./ 35.\*B+ 11./ 35.\*A 1./ 50.+ 1./ 5.\*U)\*A1\*B1  
SBI(63)=D\*(+ 54./ 35.\*B-156./ 35.\*A 72./ 25.  
SBI(64)=D\*( -13./ 35.\*B+ 78./ 35.\*A 6./ 25.  
SBI(65)=D\*(+ 27./ 35.\*B- 22./ 35.\*A 6./ 25.- 6./ 5.\*U)\*A1  
SBI(66)=D\*( -13./ 70.\*B+ 11./ 35.\*A 1./ 50.+ 1./ 10.\*U)\*A1\*B1  
SBI(67)=D\*( -54./ 35.\*B- 54./ 35.\*A 72./ 25.  
SBI(68)=D\*( 13./ 35.\*B+ 27./ 35.\*A 6./ 25. )\*B1  
SBI(69)=D\*(+ 27./ 35.\*B+ 13./ 35.\*A 6./ 25. )\*A1  
SBI(70)=D\*( -13./ 70.\*B- 13./ 70.\*A 1./ 50.  
SBI(71)=D\*(+ 4./ 35.\*B+ 52./ 35.\*A 8./ 25. )\*B2  
SBI(72)=D\*(+ 11./ 35.\*B+ 11./ 35.\*A 1./ 50.+ 6./ 5.\*U)\*A1\*B1  
SBI(73)=D\*(+ 2./ 35.\*B+ 22./ 105.\*A 2./ 75.+ 2./ 15.\*U)\*A1\*B2  
SBI(74)=D\*(+ 13./ 35.\*B- 78./ 35.\*A 6./ 25. )\*B1  
SBI(75)=D\*( - 3./ 35.\*B+ 26./ 35.\*A 2./ 25. )\*B2  
SBI(76)=D\*(+ 13./ 70.\*B- 11./ 35.\*A 1./ 50.- 1./ 10.\*U)\*A1\*B1  
SBI(77)=D\*( - 3./ 70.\*B+ 11./ 105.\*A 1./ 150.- 1./ 30.\*U)\*A1\*B2  
SBI(78)=D\*(+ 13./ 35.\*B- 27./ 35.\*A 6./ 25. )\*B1  
SBI(79)=D\*(+ 3./ 35.\*B+ 9./ 35.\*A 2./ 25. )\*B2  
SBI(80)=D\*(+ 13./ 70.\*B+ 13./ 70.\*A 1./ 50.  
SBI(81)=D\*( - 3./ 70.\*B- 13./ 210.\*A 1./ 150.  
SBI(82)=D\*(+ 52./ 35.\*B+ 4./ 35.\*A 8./ 25.  
SBI(83)=D\*(+ 22./ 105.\*B+ 2./ 35.\*A 2./ 75.+ 2./ 15.\*U)\*A2\*B1  
SBI(84)=D\*(+ 27./ 35.\*B- 22./ 35.\*A 6./ 25.- 6./ 5.\*U)\*A1  
SBI(85)=D\*( -13./ 70.\*B+ 11./ 35.\*A 1./ 50.+ 1./ 10.\*U)\*A1\*B1  
SBI(86)=D\*(+ 18./ 35.\*B- 4./ 35.\*A 8./ 25. )\*A2  
SBI(87)=D\*( -13./ 105.\*B+ 2./ 35.\*A 2./ 75.  
SBI(88)=D\*( -27./ 35.\*B- 13./ 35.\*A 6./ 25.  
SBI(89)=D\*(+ 13./ 70.\*B+ 13./ 70.\*A 1./ 50.  
SBI(90)=D\*(+ 9./ 35.\*B+ 3./ 35.\*A 2./ 25. )\*A2  
SBI(91)=D\*( -13./ 210.\*B- 3./ 70.\*A 1./ 150.  
SBI(92)=D\*(+ 4./ 105.\*B+ 4./ 105.\*A 8./ 225.  
SBI(93)=D\*(+ 13./ 70.\*B+ 11./ 35.\*A 1./ 50.- 1./ 10.\*U)\*A1\*B1  
SBI(94)=D\*( - 3./ 70.\*B+ 11./ 105.\*A 1./ 150.- 1./ 30.\*U)\*A1\*B2  
SBI(95)=D\*(+ 13./ 105.\*B- 2./ 35.\*A 2./ 75.  
SBI(96)=D\*( - 1./ 35.\*B+ 2./ 105.\*A 2./ 225.  
SBI(97)=D\*( -11./ 70.\*B- 13./ 70.\*A 1./ 50.  
SBI(98)=D\*(+ 1./ 70.\*B+ 13./ 210.\*A 1./ 150.  
SBI(99)=D\*(+ 13./ 210.\*B+ 3./ 70.\*A 1./ 150.  
SBI(100)=D\*( - 1./ 70.\*B- 1./ 70.\*A 1./ 450.  
SBI(101)=D\*(+ 9./ 35.\*B+ 3./ 35.\*A 2./ 25. )\*A2

SBI(102)=D\*(+ 22./ 105.\*B+ 2./ 35.\*A 2./ 75.+ 2./ 15.\*U)\*A2\*B1  
SBI(103)=D\*(+ 78./ 35.\*B- 13./ 35.\*A 6./ 25. )\*A1  
SBI(104)=D\*(+ 11./ 35.\*B- 13./ 71.\*A 1./ 50.+ 1./ 10.\*U)\*A1\*B1  
SBI(105)=D\*(+ 26./ 35.\*B- 3./ 35.\*A 2./ 25. )\*A2  
SBI(106)=D\*(+ 11./ 105.\*B- 3./ 71.\*A 1./ 150.- 1./ 30.\*U)\*A2\*B1  
SBI(107)=D\*(+ 27./ 35.\*B+ 13./ 35.\*A 6./ 25. )\*A1  
SBI(108)=D\*( -13./ 70.\*B- 13./ 71.\*A 1./ 50.  
SBI(109)=D\*(+ 9./ 35.\*B+ 3./ 35.\*A 2./ 25. )\*A2  
SBI(110)=D\*( -13./ 210.\*B- 3./ 70.\*A 1./ 150.  
SBI(111)=D\*(+ 4./ 105.\*B+ 4./ 105.\*A 8./ 225.  
SBI(112)=D\*(+ 13./ 70.\*B+ 11./ 35.\*A 1./ 50.- 1./ 10.\*U)\*A1\*B1  
SBI(113)=D\*( - 3./ 70.\*B+ 11./ 105.\*A 1./ 150.- 1./ 30.\*U)\*A1\*B2  
SBI(114)=D\*(+ 13./ 105.\*B- 2./ 35.\*A 2./ 75.  
SBI(115)=D\*( - 1./ 35.\*B+ 2./ 105.\*A 2./ 225.  
SBI(116)=D\*( -11./ 70.\*B- 13./ 70.\*A 1./ 50.  
SBI(117)=D\*(+ 1./ 70.\*B+ 13./ 210.\*A 1./ 150.  
SBI(118)=D\*(+ 13./ 210.\*B+ 3./ 70.\*A 1./ 150.  
SBI(119)=D\*( - 1./ 70.\*B- 1./ 70.\*A 1./ 450.  
SBI(120)=D\*(+ 9./ 35.\*B+ 3./ 35.\*A 2./ 25. )\*A2

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S8(101)=0*(156./35.*8+156./35.*A+72./25.)
S8(102)=0*(-22./35.*8-78./35.*A-6./25.-6./5.*U)*91
S8(103)=0*(78./35.*8+22./35.*A+6./25.+6./5.*U)*A1
S8(104)=0*(-11./35.*8-11./35.*A-1./50.-1./5.*U)*A1*B1
S8(105)=0*(-156./35.*8+54./35.*A-72./25.)
S8(106)=0*(22./35.*8-27./35.*A+6./25.+6./5.*U)*B1
S8(107)=0*(78./35.*8-13./35.*A+6./25.-1./10.*U)*A1*B1
S8(108)=0*(-11./35.*8+13./70.*A-1./50.-1./10.*U)*A1*B1
S8(109)=0*(4./35.*8+52./35.*A+8./25.)
S8(110)=0*(-11./35.*8-11./35.*A-1./50.-6./5.*U)*A1*B1
S8(111)=0*(2./35.*8+22./105.*A+2./75.+2./15.*U)*A1*B2
S8(112)=0*(22./35.*8-27./35.*A+6./25.+6./5.*U)*B1
S8(113)=0*(4./35.*8+14./35.*A+8./25.)
S8(114)=0*(-11./35.*8+13./70.*A-1./50.-1./10.*U)*A1*B1
S8(115)=0*(2./35.*8-13./105.*A+2./75.)
S8(116)=0*(52./35.*8+4./35.*A+8./25.)
S8(117)=0*(-22./105.*8-2./35.*A-2./75.-2./15.*U)*A2*B1
S8(118)=0*(78./35.*8+13./35.*A+6./25.)
S8(119)=0*(11./35.*8-13./70.*A+1./50.+1./10.*U)*A1*B1
S8(120)=0*(26./35.*8-3./35.*A-2./25.)
S8(121)=0*(-11./105.*8+3./70.*A+1./150.+1./30.*U)*A2*B1
S8(122)=0*(4./105.*8+4./105.*A+8./225.)
S8(123)=0*(11./35.*8-13./70.*A+1./50.+1./10.*U)*A1*B1
S8(124)=0*(-2./35.*8+13./105.*A-2./75.)
S8(125)=0*(-11./105.*8+3./70.*A+1./150.+1./30.*U)*A2*B1
S8(126)=0*(2./105.*8-1./35.*A-2./225.)
S8(127)=0*(156./35.*8+156./35.*A+72./25.)
S8(128)=0*(-22./35.*8-78./35.*A-6./25.-6./5.*U)*B1
S8(129)=0*(78./35.*8-22./35.*A-6./25.-6./5.*U)*A1
S8(130)=0*(11./35.*8+11./35.*A+1./50.+1./5.*U)*A1*B1
S8(131)=0*(4./35.*8+52./35.*A+8./25.)
S8(132)=0*(11./35.*8+11./35.*A+1./50.+6./5.*U)*A1*B1
S8(133)=0*(-2./35.*8-22./105.*A-2./75.-2./15.*U)*A1*B2
S8(134)=0*(52./35.*8+4./35.*A+8./25.)
S8(135)=0*(-22./105.*8-2./35.*A-2./75.-2./15.*U)*A2*B1
S8(136)=0*(4./105.*8+4./105.*A+8./225.)
S8001200
S8001210
S8001220
S8001230
S8001240
S8001250
S8001260
S8001270
S8001280
S8001290
S8001300
S8001310
S8001320
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S8001450
S8001460
S8001470
S8001480
S8001490
S8001500
S8001510
S8001520
S8001530
S8001540
S8001550
S8001560
S8001570

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```

*DECK S80W
SUBROUTINE S80W (XA,YB,TT,E,PR,SOM)
C
C-----THIS SUBPROGRAM GENERATES THE ELEMENT STIFFNESS MATRIX FOR THE
C 16 DEGREE-OF-FREEDOM REFINED RECTANGULAR PLANE STRESS ELEMENT.
C
C*****MAX1 AND MAX2 ARE THE DIMENSIONS OF ARRAY 'SOM'
C IN THE CALLING PROGRAM.
C
C-----CALLED BY SUBPROGRAM PANLSM
C
C
REAL SOM(136)
C
E1=(TT*E)/(1.-PR*PR)
E2=1.-PR
SOM(1)=E1*(13.*YB/(35.*XA)+E2*XA/(5.*YB))
SOM(2)=E1*(-125.-125*PR)
SOM(3)=E1*(XA/48.-PR*XA/16.)
SOM(4)=E1*(11.*YB*YB/(210.*XA)+E2*XA/60.)
SOM(5)=E1*(-13.*YB/(35.*XA)+E2*XA/(10.*YB))
SOM(6)=E1*(-125.-375*PR)
SOM(7)=E1*(-XA/48.+PR*XA/16.)
SOM(8)=E1*(-11.*YB*YB/(210.*XA)+E2*XA/120.)
SOM(9)=E1*(9.*YB/(70.*XA)-E2*XA/(5.*YB))
SOM(10)=E1*(-125.-375*PR)
SOM(11)=E1*(XA/48.+PR*XA/48.)
SOM(12)=E1*(-13.*YB*YB/(420.*XA)+E2*XA/60.)
SOM(13)=E1*(-9.*YB/(70.*XA)-E2*XA/(10.*YB))
SOM(14)=E1*(-125+PR*.125)
SOM(15)=E1*(-XA/48.-PR*XA/48.)
SOM(16)=E1*(13.*YB*YB/(420.*XA)+E2*XA/120.)
SOM(17)=E1*(13.*XA/(35.*YB)+E2*YB/(5.*XA))
SOM(18)=E1*(11.*XA*XA/(210.*YB)+E2*YB/60.)
SOM(19)=E1*(YB/48.-PR*YB/16.)
SOM(20)=E1*(-375+PR*.125)
SOM(21)=E1*(9.*XA/(70.*YB)-E2*YB/(5.*XA))
SOM(22)=E1*(-13.*XA*XA/(420.*YB)+E2*YB/60.)
SOM(23)=E1*(YB/48.+PR*YB/48.)
SOM(24)=E1*(-125-PR*.375)
SOM(25)=E1*(-13.*XA/(35.*YB)+E2*YB/(10.*XA))
SOM(26)=E1*(-11.*XA*XA/(210.*YB)+E2*YB/120.)
SOM(27)=E1*(-YB/48.+PR*YB/16.)
SOM(28)=E1*(-125+PR*.125)
SOM(29)=E1*(-9.*XA/(70.*YB)-E2*YB/(10.*XA))
SOM(30)=E1*(13.*XA*XA/(420.*YB)+E2*YB/120.)
SOM(31)=E1*(-YB/48.-PR*YB/48.)
SOM(32)=E1*(XA*XA*XA/(105.*YB)+E2*XA*YB/45.)
SOM(33)=E1*(-XA*YB/288.-PR*XA*YB/288.)
SOM(34)=E1*(-XA/48.+PR*XA/16.)
SOM(35)=E1*(13.*XA*XA/(420.*YB)-E2*YB/60.)
SOM(36)=E1*(-XA*XA*XA/(140.*YB)-E2*XA*YB/180.)
SOM(37)=E1*(XA*YB/288.+PR*XA*YB/288.)
SOM(38)=E1*(-XA/48.-PR*XA/48.)
SOM(39)=E1*(-11.*XA*XA/(210.*YB)+E2*YB/120.)
SOM(40)=E1*(-XA*XA*XA/(105.*YB)+E2*XA*YB/90.)
SOM(41)=E1*(XA*YB/288.+PR*XA*YB/288.)
SOM(42)=E1*(XA/48.+PR*XA/48.)
SOM(43)=E1*(-13.*XA*XA/(420.*YB)-E2*YB/120.)
SOM(44)=E1*(XA*XA*XA/(140.*YB)-E2*XA*YB/180.)

```

```

S8000100
S8000110
S8000120
S8000130
S8000140
S8000150
S8000160
S8000170
S8000180
S8000190
S8000200
S8000210
S8000220
S8000230
S8000240
S8000250
S8000260
S8000270
S8000280
S8000290
S8000300
S8000310
S8000320
S8000330
S8000340
S8000350
S8000360
S8000370
S8000380
S8000390
S8000400
S8000410
S8000420
S8000430
S8000440
S8000450
S8000460
S8000470
S8000480
S8000490
S8000500
S8000510
S8000520
S8000530
S8000540
S8000550
S8000560
S8000570
S8000580
S8000590

```



\*DECK BOGCM  
SUBROUTINE BOGCM (XMD,A1,B1,T,CM)

C-----THIS SUBPROGRAM GENERATES THE CONSISTENT MASS MATRIX FOR THE 16  
C DEGREE-OF-FREEDOM PLATE BENDING ELEMENT KNOWN AS THE BOGNER  
C RECTANGLE.

C-----NOTATION:

C XMD = MASS DENSITY (KIPS OF LBS-SEC\*\*2/INCHES\*\*4)  
C A1,B1 = WIDTH AND HEIGHT OF ELEMENT (INCHES)  
C T = ELEMENT THICKNESS (INCHES)  
C CM = VECTOR USED TO STORE THE UPPER TRIANGLE OF THE  
C CONSISTENT MASS MATRIX

C-----CALLED BY SUBPROGRAM PANLSM

C REAL CM(136)

C1 = XMD\*T\*A1\*B1/1225.0

AZ = A1\*A1

BZ = B1\*B1

C  
CM( 1) = 169.0/ 1.0\*C1  
CM( 2) = 143.0/ 6.0\*C1\*B1  
CM( 3) = -143.0/ 6.0\*C1\*A1  
CM( 4) = -121.0/ 36.0\*C1\*A1\*B1  
CM( 5) = 117.0/ 2.0\*C1  
CM( 6) = 33.0/ 4.0\*C1\*B1  
CM( 7) = +169.0/ 12.0\*C1\*A1  
CM( 8) = 143.0/ 72.0\*C1\*A1\*B1  
CM( 9) = 81.0/ 4.0\*C1  
CM(10) = - 39.0/ 8.0\*C1\*B1  
CM(11) = + 39.0/ 8.0\*C1\*A1  
CM(12) = -169.0/144.0\*C1\*A1\*B1  
CM(13) = 117.0/ 2.0\*C1  
CM(14) = -169.0/ 12.0\*C1\*B1  
CM(15) = - 33.0/ 4.0\*C1\*A1  
CM(16) = 143.0/ 72.0\*C1\*A1\*B1  
CM(17) = 13.0/ 3.6\*C1\*B2  
CM(18) = -121.0/ 36.0\*C1\*A1\*B1  
CM(19) = - 11.0/ 18.0\*C1\*A1\*B2  
CM(20) = 33.0/ 4.0\*C1\*B1  
CM(21) = 3.0/ 2.0\*C1\*B2  
CM(22) = +143.0/ 72.0\*C1\*A1\*B1  
CM(23) = 13.0/ 36.0\*C1\*A1\*B2  
CM(24) = 39.0/ 8.0\*C1\*B1  
CM(25) = - 9.0/ 8.0\*C1\*B2  
CM(26) = +169.0/144.0\*C1\*A1\*B1  
CM(27) = - 13.0/ 48.0\*C1\*A1\*B2  
CM(28) = 169.0/ 12.0\*C1\*B1  
CM(29) = - 13.0/ 4.0\*C1\*B2  
CM(30) = -143.0/ 72.0\*C1\*A1\*B1  
CM(31) = 11.0/ 24.0\*C1\*A1\*B2  
CM(32) = 13.0/ 3.0\*C1\*A2  
CM(33) = 11.0/ 18.0\*C1\*A2\*B1  
CM(34) = -169.0/ 12.0\*C1\*A1  
CM(35) = -143.0/ 72.0\*C1\*A1\*B1  
CM(36) = - 13.0/ 4.0\*C1\*A2  
CM(37) = - 11.0/ 24.0\*C1\*A2\*B1  
CM(38) = - 39.0/ 8.0\*C1\*A1

BOGCM00C

BOGCM0010

BOGCM0020

BOGCM0030

BOGCM0040

BOGCM0050

BOGCM0060

BOGCM0070

BOGCM0080

BOGCM0090

BOGCM0100

BOGCM0110

BOGCM0120

BOGCM0130

BOGCM0140

BOGCM0150

BOGCM0160

BOGCM0170

BOGCM0180

BOGCM0190

BOGCM0200

BOGCM0210

BOGCM0220

BOGCM0230

BOGCM0240

BOGCM0250

BOGCM0260

BOGCM0270

BOGCM0280

BOGCM0290

BOGCM0300

BOGCM0310

BOGCM0320

BOGCM0330

BOGCM0340

BOGCM0350

BOGCM0360

BOGCM0370

BOGCM0380

BOGCM0390

BOGCM0400

BOGCM0410

BOGCM0420

BOGCM0430

BOGCM0440

BOGCM0450

BOGCM0460

BOGCM0470

BOGCM0480

BOGCM0490

BOGCM0500

BOGCM0510

BOGCM0520

BOGCM0530

BOGCM0540

BOGCM0550

BOGCM0560

BOGCM0570

BOGCM0580

BOGCM0590

CM( 39) = +169.0/144.0\*C1\*A1\*B1

CM( 40) = - 9.0/ 8.0\*C1\*A2

CM( 41) = 13.0/ 48.0\*C1\*A2\*B1

CM( 42) = - 33.0/ 4.0\*C1\*A1

CM( 43) = +143.0/ 72.0\*C1\*A1\*B1

CM( 44) = 3.0/ 2.0\*C1\*A2

CM( 45) = - 13.0/ 36.0\*C1\*A2\*B1

CM( 46) = 1.0/ 9.0\*C1\*A2\*B2

CM( 47) = -143.0/ 72.0\*C1\*A1\*B1

CM( 48) = - 13.0/ 36.0\*C1\*A1\*B2

CM( 49) = - 11.0/ 24.0\*C1\*A2\*B1

CM( 50) = - 1.0/ 12.0\*C1\*A2\*B2

CM( 51) = -169.0/144.0\*C1\*A1\*B1

CM( 52) = 13.0/ 48.0\*C1\*A1\*B2

CM( 53) = - 13.0/ 48.0\*C1\*A2\*B1

CM( 54) = 1.0/ 16.0\*C1\*A2\*B2

CM( 55) = -143.0/ 72.0\*C1\*A1\*B1

CM( 56) = 11.0/ 24.0\*C1\*A1\*B2

CM( 57) = 13.0/ 36.0\*C1\*A2\*B1

CM( 58) = - 1.0/ 12.0\*C1\*A2\*B2

CM( 59) = 169.0/ 1.0\*C1

CM( 60) = 143.0/ 6.0\*C1\*B1

CM( 61) = +143.0/ 6.0\*C1\*A1

CM( 62) = 121.0/ 36.0\*C1\*A1\*B1

CM( 63) = 117.0/ 2.0\*C1

CM( 64) = -169.0/ 12.0\*C1\*B1

CM( 65) = + 33.0/ 4.0\*C1\*A1

CM( 66) = -143.0/ 72.0\*C1\*A1\*B1

CM( 67) = 81.0/ 4.0\*C1

CM( 68) = - 39.0/ 8.0\*C1\*B1

CM( 69) = - 39.0/ 8.0\*C1\*A1

CM( 70) = +169.0/144.0\*C1\*A1\*B1

CM( 71) = 13.0/ 3.0\*C1\*B2

CM( 72) = +121.0/ 36.0\*C1\*A1\*B1

CM( 73) = 11.0/ 18.0\*C1\*A1\*B2

CM( 74) = 169.0/ 12.0\*C1\*B1

CM( 75) = - 13.0/ 4.0\*C1\*B2

CM( 76) = +143.0/ 72.0\*C1\*A1\*B1

CM( 77) = - 11.0/ 24.0\*C1\*A1\*B2

CM( 78) = 39.0/ 8.0\*C1\*B1

CM( 79) = - 9.0/ 8.0\*C1\*B2

CM( 80) = -169.0/144.0\*C1\*A1\*B1

CM( 81) = 13.0/ 48.0\*C1\*A1\*B2

CM( 82) = 13.0/ 3.0\*C1\*A2

CM( 83) = 11.0/ 18.0\*C1\*A2\*B1

CM( 84) = + 33.0/ 4.0\*C1\*A1

CM( 85) = -143.0/ 72.0\*C1\*A1\*B1

CM( 86) = 3.0/ 2.0\*C1\*A2

CM( 87) = - 13.0/ 36.0\*C1\*A2\*B1

CM( 88) = + 39.0/ 8.0\*C1\*A1

CM( 89) = -169.0/144.0\*C1\*A1\*B1

CM( 90) = - 9.0/ 8.0\*C1\*A2

CM( 91) = 13.0/ 48.0\*C1\*A2\*B1

CM( 92) = 1.0/ 9.0\*C1\*A2\*B2

CM( 93) = 143.0/ 72.0\*C1\*A1\*B1

CM( 94) = - 11.0/ 24.0\*C1\*A1\*B2

CM( 95) = + 13.0/ 36.0\*C1\*A2\*B1

CM( 96) = - 1.0/ 12.0\*C1\*A2\*B2

CM( 97) = +169.0/144.0\*C1\*A1\*B1

CM( 98) = - 13.0/ 48.0\*C1\*A1\*B2

BOGCM0610

BOGCM0620

BOGCM0630

BOGCM0640

BOGCM0650

BOGCM0660

BOGCM0670

BOGCM0680

BOGCM0690

BOGCM0700

BOGCM0710

BOGCM0720

BOGCM0730

BOGCM0740

BOGCM0750

BOGCM0760

BOGCM0770

BOGCM0780

BOGCM0790

BOGCM0800

BOGCM0810

BOGCM0820

BOGCM0830

BOGCM0840

BOGCM0850

BOGCM0860

BOGCM0870

BOGCM0880

BOGCM0890

BOGCM0900

BOGCM0910

BOGCM0920

BOGCM0930

BOGCM0940

BOGCM0950

BOGCM0960

BOGCM0970

BOGCM0980

BOGCM0990

BOGCM1000

BOGCM1010

BOGCM1020

BOGCM1030

BOGCM1040

BOGCM1050

BOGCM1060

BOGCM1070

BOGCM1080

BOGCM1090

BOGCM1100

BOGCM1110

BOGCM1120

BOGCM1130

BOGCM1140

BOGCM1150

BOGCM1160

BOGCM1170

BOGCM1180

BOGCM1190

```

CM( 99) = - 13.0/ 48.0*C1*A2*B1
CM(100) =   1.0/ 16.0*C1*A2*B2
CM(101) =  169.0/  1.0*C1
CM(102) = -143.0/  6.0*C1*B1
CM(103) =  143.0/  6.0*C1*A1
CM(104) = -121.0/ 36.0*C1*A1*B1
CM(105) =  117.0/  2.0*C1
CM(106) = - 33.0/  4.0*C1*B1
CM(107) = -169.0/ 12.0*C1*A1
CM(108) =  143.0/ 72.0*C1*A1*B1
CM(109) =   13.0/  3.0*C1*B2
CM(110) = -121.0/ 36.0*C1*A1*B1
CM(111) =   11.0/ 18.0*C1*A1*B2
CM(112) = - 33.0/  4.0*C1*B1
CM(113) =   3.0/  2.0*C1*B2
CM(114) = +143.0/ 72.0*C1*A1*B1
CM(115) = - 13.0/ 36.0*C1*A1*B2
CM(116) =   13.0/  3.0*C1*A2
CM(117) = - 11.0/ 18.0*C1*A2*B1
CM(118) = +169.0/ 12.0*C1*A1
CM(119) = -143.0/ 72.0*C1*A1*B1
CM(120) = - 13.0/  4.0*C1*A2
CM(121) =  11.0/ 24.0*C1*A2*B1
CM(122) =   1.0/  9.0*C1*A2*B2
CM(123) = -143.0/ 72.0*C1*A1*B1
CM(124) =   13.0/ 36.0*C1*A1*B2
CM(125) =  11.0/ 24.0*C1*A2*B1
CM(126) = - 1.0/ 12.0*C1*A2*B2
CM(127) =  169.0/  1.0*C1
CM(128) = -143.0/  6.0*C1*B1
CM(129) = -143.0/  6.0*C1*A1
CM(130) =  121.0/ 36.0*C1*A1*B1
CM(131) =   13.0/  3.0*C1*B2
CM(132) = +121.0/ 36.0*C1*A1*B1
CM(133) = - 11.0/ 18.0*C1*A1*B2
CM(134) =   13.0/  3.0*C1*A2
CM(135) = - 11.0/ 18.0*C1*A2*B1
CM(136) =   1.0/  9.0*C1*A2*B2

```

```

C RETURN
END

```

```

ROG1200
ROG1210
ROG1220
ROG1230
ROG1240
ROG1250
ROG1260
ROG1270
ROG1280
ROG1290
ROG1300
ROG1310
ROG1320
ROG1330
ROG1340
ROG1350
ROG1360
ROG1370
ROG1380
ROG1390
ROG1400
ROG1410
ROG1420
ROG1430
ROG1440
ROG1450
ROG1460
ROG1470
ROG1480
ROG1490
ROG1500
ROG1510
ROG1520
ROG1530
ROG1540
ROG1550
ROG1560
ROG1570
ROG1580
ROG1590
ROG1600

```

```
*DECK BOWCM
```

```
SUBROUTINE BOWCM (XMO,A,B,TT,CM)
```

```

C
C-----THIS SUBPROGRAM GENERATES THE CONSISTENT MASS MATRIX FOR THE 16
C DEGREE-OF-FREEDOM REFINED RECTANGULAR PLANE STRESS ELEMENT (REFER
C TO R. BARBER'S PH.D. THESIS, 'COMPONENT MODE DYNAMIC ANALYSIS OF
C MULTISTORY FRAMES WITH SHEAR WALLS', UTAH STATE UNIVERSITY, LOGAN,
C UTAH, 1973, FOR DERIVATION).

```

```
C-----NOTATION:
```

```

C XMO = MASS DENSITY (KIPS OR LBS-SEC**2/INCHES**4)
C A,B = WIDTH AND HEIGHT OF ELEMENT (INCHES)
C TT = ELEMENT THICKNESS (INCHES)
C CM = ARRAY USED TO STORE THE UPPER TRIANGLE OF THE CONSISTENT
C MASS MATRIX
C MAX1,MAX2 = DIMENSIONS OF ARRAY 'CM' IN THE CALLING PROGRAM

```

```
C-----CALLED BY SUBPROGRAM PANLSM
```

```

C REAL CM(136)
C DO 100 I=1,136
C 100 CM(I)=0.

```

```
C
```

```
C
```

```

XMO=XMO*TT
C1=XMO*A*B
C2=C1*B
C3=C1*A
C4=C3*A
C5=C2*B

```

```
C
```

```

B10=13.*C1/185.
B11=11.*C2/630.
B12=B10/2.
B13=B11/2.
B14=3.*C1/70.
B15=13.*C2/1260.
B16=B14/2.
B17=B15/2.
B18=11.*C3/630.
B19=9.*C1/210.
B20=13.*C3/1260.
B21=B10/2.
B22=B20/2.
B23=C4/315.
B24=C4/420.
B25=B23/2.
B26=B24/2.
B27=C5/315.
B28=B27/2.
B29=C5/420.
B30=B29/2.

```

```
C
```

```

CM(1)=B10
CM(4)=B11
CM(5)=B12
CM(8)=B13
CM(9)=B14
CM(12)=-B15
CM(13)=B16

```

```

BOM000C
BOM0010
BOM0020
BOM0030
BOM0040
BOM0050
BOM0060
BOM0070
BOM0080
BOM0090
BOM0100
BOM0110
BOM0120
BOM0130
BOM0140
BOM0150
BOM0160
BOM0170
BOM0180
BOM0190
BOM0200
BOM0210
BOM0220
BOM0230
BOM0240
BOM0250
BOM0260
BOM0270
BOM0280
BOM0290
BOM0300
BOM0310
BOM0320
BOM0330
BOM0340
BOM0350
BOM0360
BOM0370
BOM0380
BOM0390
BOM0400
BOM0410
BOM0420
BOM0430
BOM0440
BOM0450
BOM0460
BOM0470
BOM0480
BOM0490
BOM0500
BOM0510
BOM0520
BOM0530
BOM0540
BOM0550
BOM0560
BOM0570
BOM0580
BOM0590

```

```

C      C(1,33)=-R11
C      C(1,31)=R10
C      C(1,32)=-R10
C      C(1,34)=R83
C      C(1,35)=R87

      RETURN

      END

      ROM1,120:
      ROM1,1210
      ROM1,1220
      ROM1,1230
      ROM1,1240
      ROM1,1250
      ROM1,1260
      ROM1,1270

```

CM1167=-81.7	ROM0600
CM1177=81.0	ROM0610
CM1187=81.6	ROM0620
CM1187=81.9	ROM0631
CM1227=-82.0	ROM0640
CM1257=81.2	ROM0650
CM1267=82.1	ROM0660
CM1297=81.6	ROM0670
CM1307=-82.2	ROM0680
CM1327=82.3	ROM0690
CM1357=-82.4	ROM0700
CM1397=82.1	ROM0710
CM1417=82.2	ROM0720
CM1447=82.6	ROM0730
CM1467=82.7	ROM0740
CM1477=81.3	ROM0750
CM1507=82.8	ROM0760
CM1517=81.5	ROM0770
CM1547=-82.9	ROM0780
CM1557=81.7	ROM0790
CM1587=-83.0	ROM0810
CM1597=81.0	ROM0820
CM1627=81.1	ROM0831
CM1637=81.6	ROM0840
CM1667=-81.7	ROM0850
CM1677=81.4	ROM0860
CM1707=-81.5	ROM0870
CM1717=81.0	ROM0880
CM1727=-81.6	ROM0890
CM1757=81.6	ROM0910
CM1767=82.2	ROM0920
CM1797=81.2	ROM0931
CM1807=-82.1	ROM0940
CM1827=82.3	ROM0950
CM1857=-82.2	ROM0960
CM1867=-82.6	ROM0970
CM1897=-82.1	ROM0980
CM1907=82.5	ROM0990
CM1927=82.7	ROM1000
CM1937=81.7	ROM1010
CM1967=-83.0	ROM1020
CM1977=81.5	ROM1031
CM1987=-82.9	ROM1040
CM1997=81.8	ROM1050
CM1007=-81.1	ROM1060
CM1057=81.2	ROM1070
CM1087=81.3	ROM1080
CM1107=81.0	ROM1090
CM1117=81.4	ROM1100
CM1117=81.4	ROM1110
CM1117=-82.0	ROM1120
CM11167=82.3	ROM1137
CM1117=82.0	ROM1140
CM11207=-82.4	ROM1150
CM11227=82.7	ROM1160
CM11237=-81.3	ROM1177
CM11267=81.0	ROM1180
CM11277=81.0	ROM1190

```

*DECK TSET
SUBROUTINE TSET(PV)
C
C-----SUBPROGRAM TO ASSEMBLE PERMUTATION VECTOR FOR FORMATION OF ELEMENT
C STIFFNESS AND MASS MATRICES FOR THE COMBINATION ELEMENT COMPOSED
C OF THE HOGNER AND R-O-M RECTANGLES
C
C      INTEGER PV(32)
C
C      PV(1)=3
C      PV(2)=4
C      PV(3)=5
C      PV(4)=8
C      PV(5)=11
C      PV(6)=12
C      PV(7)=13
C      PV(8)=16
C      PV(9)=19
C      PV(10)=20
C      PV(11)=21
C      PV(12)=24
C      PV(13)=27
C      PV(14)=28
C      PV(15)=29
C      PV(16)=32
C      PV(17)=25
C      PV(18)=26
C      PV(19)=30
C      PV(20)=31
C      PV(21)=17
C      PV(22)=18
C      PV(23)=22
C      PV(24)=23
C      PV(25)=1
C      PV(26)=2
C      PV(27)=6
C      PV(28)=7
C      PV(29)=9
C      PV(30)=10
C      PV(31)=14
C      PV(32)=15
C
C      RETURN
C      END

```

```

TSET0000
TSET0010
TSET0020
TSET0030
TSET0040
TSET0050
TSET0060
TSET0070
TSET0080
TSET0090
TSET0100
TSET0110
TSET0120
TSET0130
TSET0140
TSET0150
TSET0160
TSET0170
TSET0180
TSET0190
TSET0200
TSET0210
TSET0220
TSET0230
TSET0240
TSET0250
TSET0260
TSET0270
TSET0280
TSET0290
TSET0300
TSET0310
TSET0320
TSET0330
TSET0340
TSET0350
TSET0360
TSET0370
TSET0380
TSET0390
TSET0400
TSET0410
TSET0420
TSET0430

```

```

*DECK BOWSTR
SUBROUTINE BOWSTR(EP,TP,PRP,XA,YB,XS,YS,D, SX,SY,SXY)
C
C-----SUBPROGRAM TO COMPUTE STRESSES SX, SY, SXY AT LOCATION (X,Y)
C FOR THE REFINED PLANE STRESS ELEMENT OF DIMENSIONS XA BY YB.
C "D" IS THE VECTOR OF NODAL DISPLACEMENTS AND EP, TP, PRP ARE
C THE ELEMENT ELASTIC MODULUS, THICKNESS, AND POISSON'S RATIO.
C
C      REAL D(16)
C      I1=13
C      I2=14
C      I3=15
C      I4=16
C      J1=9
C      J2=10
C      J3=11
C      J4=12
C      K1=1
C      K2=2
C      K3=3
C      K4=4
C      L1=5
C      L2=6
C      L3=7
C      L4=8
C
C      E1=(EP*TP)/(1.-PRP*PRP)
C      E2=(1.-PRP)/2.
C      E3=E1/TP
C      E4=E2/2.
C      YB3=(YS/YB)**3
C      YB2=(YS/YB)**2
C      T1=(2.*YB3-3.*YB2)/XA
C      T2=(YB3-YB2)*YB/XA
C      T3=(2.*YB3-3.*YB2+1.)/XA
C      T4=(YB3-2.*YB2+YS/YB)*YB/XA
C
C      STAX=T1*(D(I1)-D(J1)) + T3*(D(L1)-D(K1))
C      + T2*(D(I4)-D(J4)) + T4*(D(K4)-D(L4))
C      XA3=(XS/XA)**3
C      XA2=(XS/XA)**2
C      T1=(2.*XA3-3.*XA2+1.)/YB
C      T2=(2.*XA3-3.*XA2)/YB
C      T3=(XA3-2.*XA2+XS/XA)*XA/YB
C      T4=(XA3-XA2)*XA/YB
C
C      STAY=T1*(D(I2)-D(K2)) + T2*(D(L2)-D(J2))
C      + T3*(D(I3)-D(K3)) + T4*(D(J3)-D(L3))
C      SX=E3*(STAX+PRP*STAY)
C      SY=E3*(STAY+PRP*STAX)
C      T1=(1.-XS/XA)
C      T2=(1.-YS/YB)
C      T3=6.*(YB2-YS/YB)
C      T4=3.*YB2-2.*YS/YB
C      T5=3.*YB2-4.*YS/YB+1.
C      T6=6.*(XA2-XS/XA)
C      T7=3.*XA2-4.*XS/XA+1.
C      T8=3.*XA2-2.*XS/XA
C      SXY1=-T1*T3*(D(I1)/YB - T1*T4*(D(I6) - XS*T3*(D(J1)/(XA*YB)
C      -XS*T4*(D(J4)/XA + T1*T3*(D(K1)/YB - T1*T5*(D(K4)

```

```

BOW0000
BOW0010
BOW0020
BOW0030
BOW0040
BOW0050
BOW0060
BOW0070
BOW0080
BOW0090
BOW0100
BOW0110
BOW0120
BOW0130
BOW0140
BOW0150
BOW0160
BOW0170
BOW0180
BOW0190
BOW0200
BOW0210
BOW0220
BOW0230
BOW0240
BOW0250
BOW0260
BOW0270
BOW0280
BOW0290
BOW0300
BOW0310
BOW0320
BOW0330
BOW0340
BOW0350
BOW0360
BOW0370
BOW0380
BOW0390
BOW0400
BOW0410
BOW0420
BOW0430
BOW0440
BOW0450
BOW0460
BOW0470
BOW0480
BOW0490
BOW0500
BOW0510
BOW0520
BOW0530
BOW0540
BOW0550
BOW0560
BOW0570
BOW0580
BOW0590

```



```

      *XS*T3*D(L1)/(XA*YB) - XS*T5*D(L4)/XA + YS*T6*D(I2)/(XA*YB) ROW10600
      *YS*T7*D(I3)/YB - YS*T6*D(J2)/(XA*YB) + YS*T8*D(J3)/YB ROW10610
      *T2*T6*D(K2)/XA + T2*T7*D(K3) - T2*T6*D(L2)/XA + T2*T8*D(L3) ROW10620
C
      SX= SX1 ROW10630
      SY= SY1 ROW10640
      SXY= SXY1*E3*E4 ROW10650
C
      RETURN ROW10660
      END ROW10670
      ROW10680
      ROW10690

```

```

*DECK BFSSTP BFS00000
SUBROUTINE BFSSTR(E,T,PR,A,B,X,Y,ON,MX,MY,MXY) BFS00010
C BFS00020
C THIS SUBPROGRAM COMPUTES THE STRESS RESULTANTS OF THE ELEMENT IE BFS00030
C AT ANY POINT (X,Y) WITH RESPECT TO THE ELEMENT COORDINATES. BFS00040
C (BOGNER-FOX-SCHMIT 16 DOF PLATE BENDING ELEMENT) BFS00050
C BFS00060
      REAL MX,MY,MXY,GMA(16),DN(1) BFS00070
C GENERAL CUBIC BEAM FUNCTIONS AND THEIR DERIVATIVES BFS00080
      F10(Z,C)=1.0-3.0*(Z/C)**2+2.0*(Z/C)**3 BFS00090
      F11(Z,C)=-6.0*Z/C**2+6.0*Z**2/C**3 BFS00100
      F12(Z,C)=-6.0/C**2+12.0*Z/C**3 BFS00110
      F20(Z,C)=3.0*(Z/C)**2-2.0*(Z/C)**3 BFS00120
      F21(Z,C)=6.0*Z/C**2-6.0*Z**2/C**3 BFS00130
      F22(Z,C)=6.0/C**2-12.0*Z/C**3 BFS00140
      G10(Z,C)=Z-2.0*Z**2/C+Z**3/C**2 BFS00150
      G11(Z,C)=1.0-4.0*Z/C+3.0*(Z/C)**2 BFS00160
      G12(Z,C)=-4.0/C+6.0*Z/C**2 BFS00170
      G20(Z,C)=-Z**2/C+Z**3/C**2 BFS00180
      G21(Z,C)=-2.0*Z/C+3.0*(Z/C)**2 BFS00190
      G22(Z,C)=-2.0/C+6.0*Z/C**2 BFS00200
C BFS00210
C BFS00220
C COMPUTE THE GENERALIZED STRAINS BFS00230
C (MXX) BFS00240
      GMA(1)=F12(X,A)*F10(Y,B) BFS00250
      GMA(2)=F12(X,A)*G10(Y,B) BFS00260
      GMA(3)=-G12(X,A)*F10(Y,B) BFS00270
      GMA(4)=-G12(X,A)*G10(Y,B) BFS00280
      GMA(5)=F22(X,A)*F10(Y,B) BFS00290
      GMA(6)=F22(X,A)*G10(Y,B) BFS00300
      GMA(7)=-G22(X,A)*F10(Y,B) BFS00310
      GMA(8)=-G22(X,A)*G10(Y,B) BFS00320
      GMA(9)=F20(X,A)*F10(Y,B) BFS00330
      GMA(10)=F20(X,A)*G10(Y,B) BFS00340
      GMA(11)=-G20(X,A)*F10(Y,B) BFS00350
      GMA(12)=-G20(X,A)*G10(Y,B) BFS00360
      GMA(13)=F12(X,A)*F20(Y,B) BFS00370
      GMA(14)=F12(X,A)*G20(Y,B) BFS00380
      GMA(15)=-G12(X,A)*F20(Y,B) BFS00390
      GMA(16)=-G12(X,A)*G20(Y,B) BFS00400
      SUM=0.0 BFS00410
      DO 100 I=1,16 BFS00420
100 SUM=SUM+GMA(I)*DN(I) BFS00430
      MXX=SUM BFS00440
C BFS00450
C (MYY) BFS00460
      GMA(1)=F10(X,A)*F12(Y,B) BFS00470
      GMA(2)=F10(X,A)*G12(Y,B) BFS00480
      GMA(3)=-G10(X,A)*F12(Y,B) BFS00490
      GMA(4)=-G10(X,A)*G12(Y,B) BFS00500
      GMA(5)=F20(X,A)*F12(Y,B) BFS00510
      GMA(6)=F20(X,A)*G12(Y,B) BFS00520
      GMA(7)=-G20(X,A)*F12(Y,B) BFS00530
      GMA(8)=-G20(X,A)*G12(Y,B) BFS00540
      GMA(9)=F20(X,A)*F22(Y,B) BFS00550
      GMA(10)=F20(X,A)*G22(Y,B) BFS00560
      GMA(11)=-G20(X,A)*F22(Y,B) BFS00570
      GMA(12)=-G20(X,A)*G22(Y,B) BFS00580
      GMA(13)=F10(X,A)*F22(Y,B) BFS00590

```

```

      GMA(14)=F10(X,A)*G22(Y,B)
      GMA(15)=-G10(X,A)*F22(Y,B)
      GMA(16)=-G10(X,A)*G22(Y,B)
      SUM=0.
      DO 200 I=1,16
    200 SUM=SUM+GMA(I)*DN(I)
      WXY=SUM
C
C (WXY)
      GMA(1)=F11(X,A)*F11(Y,B)
      GMA(2)=F11(X,A)*G11(Y,B)
      GMA(3)=-G11(X,A)*F11(Y,B)
      GMA(4)=-G11(X,A)*G11(Y,B)
      GMA(5)=F21(X,A)*F11(Y,B)
      GMA(6)=F21(X,A)*G11(Y,B)
      GMA(7)=-G21(X,A)*F11(Y,B)
      GMA(8)=-G21(X,A)*G11(Y,B)
      GMA(9)=F21(X,A)*F21(Y,B)
      GMA(10)=F21(X,A)*G21(Y,B)
      GMA(11)=-G21(X,A)*F21(Y,B)
      GMA(12)=-G21(X,A)*G21(Y,B)
      GMA(13)=F11(X,A)*F21(Y,B)
      GMA(14)=F11(X,A)*G21(Y,B)
      GMA(15)=-G11(X,A)*F21(Y,B)
      GMA(16)=-G11(X,A)*G21(Y,B)
      SUM=0.
      DO 300 I=1,16
    300 SUM=SUM+GMA(I)*DN(I)
      WXY=SUM
C
C COMPUTE THE GENERALIZED STRESSES
      D=E*T*3/(12.0*(1.0-PR*PR))
      MX=0*(MXX + PR*WXY)
      MY=0*(WXY + PR*MXX)
      MXY=D*(1.0 - PR)*MXY
      RETURN
      END

```

```

BFS00600
BFS00610
BFS00620
BFS00630
BFS00640
BFS00650
BFS00660
BFS00670
BFS00680
BFS00690
BFS00700
BFS00710
BFS00720
BFS00730
BFS00740
BFS00750
BFS00760
BFS00770
BFS00780
BFS00790
BFS00800
BFS00810
BFS00820
BFS00830
BFS00840
BFS00850
BFS00860
BFS00870
BFS00880
BFS00890
BFS00900
BFS00910
BFS00920
BFS00930
BFS00940
BFS00950
BFS00960

```

# \*DECK DYNAMIC SUBROUTINE DYNAMIC

```

C
C-----SUBPROGRAM TO COMPUTE FREQUENCIES, MODE SHAPES, AND DYNAMIC
C RESPONSE OF FRAME-PANEL STRUCTURES. DISPLACEMENT RESPONSE TO
C SPECIFIED X, Y, Z DIRECTION PIECEWISE-LINEAR FORCING FUNCTIONS
C (APPLIED FORCES (IAF=1), GROUND ACCELERATION (IGA=1), OR
C INDEPENDENT SUPPORT MOTION (ISO=1)) IS COMPUTED, PRINTED IF
C IDTH=1) AND PLOTTER PLOTTED (IF NNPLT>0). SUPPORT REACTIONS AND
C PANEL STRESSES ARE DETERMINED AT SPECIFIED TIMES OF INTEREST
C THROUGH BACKSUBSTITUTION.
C
      REAL NI,LF,MFF,MFR,LFXP,LFYP,LFPZ,LFXPA,LFYPA,LFPZA
      INTEGER PPV,FPV,PLIST
C
      COMMON
      2/CBNFJ/LFXP(32),LFYP(32),LFPZ(32),LFXPA(32),LFYPA(32),LFPZA(32),
      3/IFM(16),JJ(16),ITP(16)
      3/CBNPFD/PPV(200),FPV(196),NFL(200),NRL(200),JFL(196)
      4/CBSTIF/XN( 40, 40),SI( 40, 40),MFF( 40, 40),TT(300),GAX(300),
      5/GAY(300),GAZ(300),DX(300),DY(300),DZ(300),XDP(300),
      6/VDP(300),LF(700),SFRC( 40, 56),XP(100,196),MFR(120, 40),
      7/SMH( 45, 45),SMR( 45, 56),SRR( 56, 56)
      5/CBSTQR/XN1(80,80),FRQ(1000,3),PSQ(4060),XLBU(44060)
      6/CBMISC/FX(46),FY(46),FZ(46),TX(46),TY(46),TZ(46),
      7/BX(46),BY(46),BZ(46),BAX(46),BAY(46),BAZ(46),
      8/AX(46),AY(46),AZ(46),C(32),NI(32),PD(32),TLF(103,3),
      9/TPP(13),VIP(40),X1P(46)
      7/CBDATA/EP,TP,PRP,DENSP,EF,GF,TF,DENSP,NS,SN,MCODE,ISM,
      8/NRP,NFT,NPT,NMT,NRT,NF,NP,NM,NR,NPN,NPE,NPD,NPDF,
      9/NFJ,NFM,NJB,NFNSPR,NPNSPR,ISDA,NPR,NFR,NRFJ,
      10/NFRT,IREAC,ISTRES,ISTAT,IL,IKKK,TIME
      8/CBOYND/PLIST(300),INO,INV,IAF,IGA,ISO,NOLC,LN,NPTS,NMODES,
      9>IDMAX,IDTH,NO,01,02,FX0,FY0,FZ0,AX0,AY0,AZ0,SX,SY,SZ,
      10/SXA,SYA,SZA,NTX,NTY,NTZ
C
C-----DETERMINE FREQUENCIES AND MODE SHAPES
C
      CALL EIGENV
C
C-----COMPUTE AND PRINT/PLOT DISPLACEMENT RESPONSE FOR STRUCTURE DOF'S
C FOR "NOLC" LOADINGS
C
      LN=0
      10 LN=LN+1
C
C---1. INPUT DATA FOR DYNAMIC ANALYSIS (IF NOLC > 0)
C
      CALL DYOATA
      IF(NOLC.EQ.0)RETURN
C
C---2. COMPUTE DYNAMIC RESPONSE
C
      CALL RESPON
C
C---3. BACKSUBSTITUTE TO OBTAIN PANEL AND FRAME DISPLACEMENTS AT ALL
C MODES; SUPPORT REACTIONS; AND PANEL STRESSES AND STRESS RESULTS
C AT SPECIFIED TIMES OF INTEREST ONLY
C
      READ(5,*) NSATI,IREAC,ISTRES

```

```

DYN00000
DYN00010
DYN00020
DYN00030
DYN00040
DYN00050
DYN00060
DYN00070
DYN00080
DYN00090
DYN00100
DYN00110
DYN00120
DYN00130
DYN00140
DYN00150
DYN00160
DYN00170
DYN00180
DYN00190
DYN00200
DYN00210
DYN00220
DYN00230
DYN00240
DYN00250
DYN00260
DYN00270
DYN00280
DYN00290
DYN00300
DYN00310
DYN00320
DYN00330
DYN00340
DYN00350
DYN00360
DYN00370
DYN00380
DYN00390
DYN00400
DYN00410
DYN00420
DYN00430
DYN00440
DYN00450
DYN00460
DYN00470
DYN00480
DYN00490
DYN00500
DYN00510
DYN00520
DYN00530
DYN00540
DYN00550
DYN00560
DYN00570
DYN00580
DYN00590

```

```

      IL=0
      IF(NSATI.EQ.0)GO TO 11
      PRINT 13
13  FORMAT('1 DYNAMIC RESPONSE - PANEL STRESSES AND SUPPORT',
      * ' REACTIONS',/,1X,55('-'))
      IF(NSATI.EQ.-1)GO TO 12
      READ(5,*) (PLIST(I),I=1,NSATI)
      GO TO 14
12  NSATI=NPTS
      DO 15 I=1,NPTS
15  PLIST(I)=I
14  DO 16 I=1,NSATI
      IL=IL+1
      ISTAT=PLIST(I)
      TIME=TT(ISTAT)
      PRINT 17,TIME
17  FORMAT(/,10X,50(' '),/,10X,'*',48(' '),/'-',/,10X,'*',14(' '),
      * 'TIME =',1PE14.5,14(' '),/'-',/,10X,'*',48(' '),/'-',/,10X,
      * 50(' '),/)
      CALL STATIC
16  CONTINUE
C
11  IF(LN.LT.NDLC)GO TO 10
      RETURN
      END

```

```

DYN0600
DYN0610
DYN0620
DYN0630
DYN0640
DYN0650
DYN0660
DYN0670
DYN0680
DYN0690
DYN0700
DYN0710
DYN0720
DYN0730
DYN0740
DYN0750
DYN0760
DYN0770
DYN0780
DYN0790
DYN0800
DYN0810
DYN0820
DYN0830
DYN0840
DYN0850

```

```

*DECK DYDATA
      SUBROUTINE DYDATA

```

```

C
C-----SUBPROGRAM TO READ IN AND PROCESS DATA FOR DYNAMIC ANALYSIS
C
      REAL NI,LF,MFF,MFR,LFXP,LFYP,LFZP,LFXPA,LFYPA,LFZPA
      INTEGER PLIST,PPV,FPV
C
      COMMON
      2/CBNFJ/LFXP(32),LFYP(32),LFZP(32),LFXPA(32),LFYPA(32),LFZPA(32),
      * IFM(16),JJ(16),ITP(16)
      3/CBNPFD/PPV(200),FPV(96),NFL(200),NRL(200),JRL(96)
      4/CBSTIF/XN( 40, 40),S( 40, 40),MFF( 40, 40),TT(300),GAX(300),
      * GAY(300),GAZ(300),DX(300),DY(300),DZ(300),XDP(300),
      * YDP(300),LF(700),SFR( 40, 56),XP(100,190),MFR(120, 56),
      * SHM( 45, 45),SMR( 45, 56),SRR( 56, 56)
      5/CBSTOR/XN1(80,80),FREQ(1000,3),PSQ(4060),XL0DA(4060)
      6/CBMISC/FX(46),FY(46),FZ(46),TX(46),TY(46),TZ(46),
      * BX(46),BY(46),BZ(46),BAX(46),BAY(46),BAZ(46),
      * AX(46),AY(46),AZ(46),G(32),NI(32),PD(32),TLF(103,3),
      * TPP(13),V1P(40),X1P(40)
      7/CBDATA/EP,TP,PRP,DENSP,EF,GF,TF,DENSF,NS,SN,MCODE,ISM,
      * NRP,NFT,NPT,NMT,NF,NP,NM,NR,NPN,NPE,NPD,NPDF,
      * NFJ,NFM,NJB,MFNSFR,NPNSPR,ISDA,NPR,NFR,NRFJ,
      * NFRT,IREFC,ISTRES,ISTAT,IL,IKKK,TIME
      8/CBOYND/PLIST(300),IND,INV,IAF,IGA,ISO,NDLC,LN,NPTS,NMODES,
      * IDMAX,IDTH,ND,01,02,FX0,FY0,FZ0,AX0,AY0,AZ0,SX,SY,SZ,
      * SKA,SYA,SZA,NTX,NTY,NTZ

```

```

C
C-----PRINT LOADING NUMBER
C

```

```

      IF(LN.GT.1)GO TO 2001
      READ(5,*) NDLC
      IF(NDLC.EQ.0) RETURN
      IKKK=0
      WRITE(11) ((XN1(I,J),J=1,NFT),I=1,NFT)
2001 PRINT 200,LN,NDLC
200  FORMAT(*1*,15X,* LOADING CONDITION NUMBER*,I3,* OF*,I3,/,15X,
      * * -----*)

```

```

C
C-----INPUT CONTROL PARAMETERS
C

```

```

      READ(5,*) IND,INV,IGA,IAF,ISO
      READ(5,*) NMODES,NNPLOT,NJPLOT,IDTH,IDMAX
      PRINT 201,IND,INV,IGA,IAF,ISO
201  FORMAT(* IND INV IGA IAF ISO*,/,I3,4I4)
      PRINT 202,NMODES,NNPLOT,NJPLOT,IDTH,IDMAX
202  FORMAT(' NMODES NNPLOT NJPLOT IDTH IDMAX',/,I5,3I7,2I5)
C
      ND=NPD+NJB
      DO 2022 I=1,ND
2022 PLIST(I)=0
      IF(NNPLOT.EQ.0)GO TO 2030
      PRINT 2023
2023 FORMAT('0 PLOT LIST FOR NCDE DISPLACEMENT RESPONSE',/,," NODE ",
      * " 01 02 03 04 05 06 07 08")
      DO 2024 I=1,NNPLOT
      READ(5,*) N,(ITP(K),K=1,8)
      IL=N*N
      I1=IL-7

```

```

DYD0900C
DYD00010
DYD00020
DYD00030
DYD00040
DYD00050
DYD00060
DYD00070
DYD00080
DYD00090
DYD00100
DYD00110
DYD00120
DYD00130
DYD00140
DYD00150
DYD00160
DYD00170
DYD00180
DYD00190
DYD00200
DYD00210
DYD00220
DYD00230
DYD00240
DYD00250
DYD00260
DYD00270
DYD00280
DYD00290
DYD00300
DYD00310
DYD00320
DYD00330
DYD00340
DYD00350
DYD00360
DYD00370
DYD00380
DYD00390
DYD00400
DYD00410
DYD00420
DYD00430
DYD00440
DYD00450
DYD00460
DYD00470
DYD00480
DYD00490
DYD00500
DYD00510
DYD00520
DYD00530
DYD00540
DYD00550
DYD00560
DYD00570
DYD00580
DYD00590

```

```

      I2=0
      DO 2025 K=I1,IL
      I2=I2+1
2025 PLIST(K)=ITP(I2)
2024 PRINT 2026,N,(PLIST(K),K=I1,IL)
2026 FORMAT(I4,1X,8I4)
2030 IF(NJPLDT.EQ.0)GO TO 2021
      PRINT 2028,(I,I=1,6)
2028 FORMAT("0PLOT LIST FOR JOINT DISPLACEMENT RESPONSE",/,
      # " JOINT",6I4," 0",I1))
      DO 2032 I=1,NJPLDT
      READ(5,*) N,(ITP(K),K=1,6)
      IL=6*N+NPU
      I1=IL-5
      I2=0
      DO 2029 K=I1,IL
      I2=I2+1
2029 PLIST(K)=ITP(I2)
2032 PRINT 2031,N,(PLIST(K),K=I1,IL)
2031 FORMAT(I5,1X,6I4)
C
C-----MODAL DAMPING DATA ( G(I)=DAMPING FRACTION IN MODE I)
C
2021 READ(5,*) (G(I),I=1,NMODES)
      PRINT 2409
2409 FORMAT(*8MODAL DAMPING RATIOS AND DAMPED FREQUENCIES(RAD/SEC)*
      PRINT 2429
2429 FORMAT(* MODE RATIO DAMPED FREQUENCY(RAD/SEC)*
2419 FORMAT(I4,F8.3,4X,1PE13.5)
      DO 203 I=1,NMODES
      P=SQR(PSQ(I))
      NI(I)=G(I)*P
      PD(I)=P*SQR(1.-G(I))*G(I))
203 PRINT 2419,I,G(I),PD(I)
C
C-----READ IN INITIAL CONDITIONS OF DISPLACEMENT (XOP) (IF INO=1) AND
C VELOCITY (VOP) (IF INV=1) AND TRANSFORM TO NORMAL COORDINATES
C (USE X1P(I), V1P(I), TTP(I) AND ITPI) AS TEMPORARY STORAGE)
C
      DO 204 I=1,NPO
204 VOP(I)=XOP(I)=0.
C-----PRINT DOF TABLE
      PRINT 171
171 FORMAT("DOOF TABLE",/,*, DOF MODE DISPL. TYPE*)
      DO 170 I=1,NPN
      IL=8*I
      I1=IL-7
      I2=0
      DO 173 K=I1,IL
      I2=I2+1
      IF(PPV(K).GT.NFT)GO TO 173
      PRINT 174,PPV(K),I,I2
173 CONTINUE
174 FORMAT(I4,I5,7X,"0",I1)
170 CONTINUE
      IF(ILN.EQ.1)GO TO 1701
      IF((IND+INV).EQ.0)GO TO 1701
      IF(IKKX.EQ.0)BACKSPACE 11
      READ(11) ((XNI(I,J),J=1,NFT),I=1,NFT)
      IKKK=0

```

```

DY000650
DY000610
DY000620
DY000630
DY000640
DY000650
DY000660
DY000670
DY000680
DY000690
DY000700
DY000710
DY000720
DY000730
DY000740
DY000750
DY000760
DY000770
DY000780
DY000790
DY000800
DY000810
DY000820
DY000830
DY000840
DY000850
DY000860
DY000870
DY000880
DY000890
DY000900
DY000910
DY000920
DY000930
DY000940
DY000950
DY000960
DY000970
DY000980
DY000990
DY01000
DY01010
DY01020
DY01030
DY01040
DY01050
DY01060
DY01070
DY01080
DY01090
DY01100
DY01110
DY01120
DY01130
DY01140
DY01150
DY01160
DY01170
DY01180
DY01190

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```

1701 IF(IND.NE.1)GO TO 205
      PRINT 236,(I,I=1,8)
206 FORMAT("0INITIAL DISPLACEMENTS",/," MODE",9X,8("0",I1,11X))
      READ(5,*) NNIND
      DO 2061 I=1,NNIND
      READ(5,*) N,(ITP(K),K=1,8)
      IL=8*N
      I1=IL-7
      I2=0
      DO 2060 K=I1,IL
      I2=I2+1
2060 XOP(PPV(K))=ITP(I2)
2061 PRINT 207,N,(XOP(PPV(K)),K=I1,IL)
207 FORMAT(I4,1X,1PE13.5)
      DO 208 K=1,NMODES
      X1P(K)=0.
      DO 208 I=1,NFT
208 X1P(K)=X1P(K)+XNI(K,I)*XOP(I)
      DO 209 K=1,NMODES
209 XOP(K)=X1P(K)
205 IF(INV.NE.1)GO TO 210
      PRINT 211,(I,I=1,8)
211 FORMAT("0INITIAL VELOCITIES",/," MODE",9X,8("0",I1,11X))
      READ(5,*) NNINV
      DO 211 I=1,NNINV
      READ(5,*) N,(ITP(K),K=1,8)
      IL=8*N
      I1=IL-7
      I2=0
      DO 2110 K=I1,IL
      I2=I2+1
2110 VOP(PPV(K))=ITP(I2)
2111 PRINT 207,N,(VOP(PPV(K)),K=I1,IL)
      DO 212 K=1,NMODES
      V1P(K)=0.
      DO 212 I=1,NFT
212 V1P(K)=V1P(K)+XNI(K,I)*VOP(I)
      DO 213 K=1,NMODES
213 VOP(K)=V1P(K)
218 CONTINUE
C
C-----CONSTRUCT VECTORS OF LOAD FACTORS IN PRINCIPAL COORDINATES
C (LFXP, LFVP, LFZP) FOR THREE LOADING CASES
C
      DO 215 K=1,NMODES
215 LFXP(K)=LFVP(K)=LFZP(K)=0.
      DO 216 I=1,N0
216 LF(I)=0.
      IF(IGA.EQ.1)GO TO 217
      IF(ISO.EQ.1)GO TO 218
      IF(IAF.NE.1)GO TO 237
C
C-----APPLIED FORCING FUNCTION CASE (PROPORTIONAL LOADING ASSUMED): 1AF=DY01.723
C 1
C
      READ(5,*) NNLF
      PRINT 219,(I,I=1,4)
219 FORMAT("0LOAD FACTORS AT NODES: APPLIED FORCING FUNCTION CASE:",/,
      # " NODE",8I4," LF",I1))
      DO 220 I=1,NNLF

```

```

DY001200
DY011210
DY011220
DY011230
DY011240
DY011250
DY011260
DY011270
DY011280
DY011290
DY011300
DY011310
DY011320
DY011330
DY011340
DY011350
DY011360
DY011370
DY011380
DY011390
DY011400
DY011410
DY011420
DY011430
DY011440
DY011450
DY011460
DY011470
DY011480
DY011490
DY011500
DY011510
DY011520
DY011530
DY011540
DY011550
DY011560
DY011570
DY011580
DY011590
DY011600
DY011610
DY011620
DY011630
DY011640
DY011650
DY011660
DY011670
DY011680
DY011690
DY011700
DY011710
DY011720
DY011730
DY011740
DY011750
DY011760
DY011770
DY011780
DY011790

```

```

      READ(5,*) N, (TPP(K), K=1, 8)
      IL=8*N
      I1=IL-7
      I2=9
      DO 221 K=I1, IL
      I2=I2+1
221  LF (PPV(K))=TPP(I2)
220  PRINT 222, N, (LF(PPV(K)), K=I1, IL)
222  FORMAT (I4, 1X, 8F9.3)
C
      DO 223 J=1, NM00ES
      LFXP(J)=LFYP(J)=LFZP(J)=0.
      DO 223 K=1, NPN
      K1=8*K-7
      K2=K1+1
      K3=K1+2
      IF (NRL(K1), EQ, 0) GO TO 225
      J1=PPV(K1)
      LFXP(J)=LFXP(J)+XN1(J1, J)*LF(J1)
225  IF (NRL(K2), EQ, 0) GO TO 226
      J2=PPV(K2)
      LFYP(J)=LFYP(J)+XN1(J2, J)*LF(J2)
226  IF (NRL(K3), EQ, 0) GO TO 223
      J3=PPV(K3)
      LFZP(J)=LFZP(J)+XN1(J3, J)*LF(J3)
223  CONTINUE
      GO TO 237
C
C-----GROUND ACCELERATION CASE (GROUND ASSUMED TO BE RIGID): IGA=1
C
      217 CONTINUE
C
C---1. COMPUTE (-)R = (S*-1)*SFR AND STORE IN XN1
C
      IF (LN, EQ, 1) GO TO 2271
      REWIND 33
      READ(33) ((S(I, J), J=1, NPT), I=1, NPT)
      READ(33) ((SFR(I, J), J=1, NPT), I=1, NPT)
      READ(33) ((MFF(I, J), J=1, NPT), I=1, NPT)
      READ(33) ((MFR(I, J), J=1, NPT), I=1, NPT)
2271 CALL OCOMP(MFT, S, NF, NF), RETURNS(227)
      GO TO 230
227 PRINT 228
228 FORMAT(///, " DECOMPOSITION OF S FAILS IN DYDATA", ///)
      RETURN
230 DO 231 I=1, NRT
231 CALL SOLVE(MFT, S, SFR, XN1, I, 1, NRT, NF, NF, NF, NR, NP, NP)
C
C---2. COMPUTE (MFF*R+MFR) WHERE R=-XN1; STORE RESULT IN SFR
C
      DO 232 I=1, NPT
      DO 232 J=1, NRT
      SUM=0.
      DO 233 K=1, NPT
      SUM=SUM-MFF(I, K)*XN1(K, J)
232 SFR(I, J)=SUM+MFR(I, J)
C
C---3. COMPUTE (-)XN* (MFF*R + MFR); STORE RESULT IN XN1
C
      DO 234 I=1, NM00ES

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```

      DY01890
      DY01810
      DY01A20
      DY01A30
      DY01A40
      DY01A50
      DY01A60
      DY01A70
      DY01A80
      DY01A90
      DY01900
      DY01910
      DY01920
      DY01930
      DY01940
      DY01950
      DY01960
      DY01970
      DY01980
      DY01990
      DY02000
      DY02010
      DY02020
      DY02030
      DY02040
      DY02050
      DY02060
      DY02070
      DY02080
      DY02090
      DY02100
      DY02110
      DY02120
      DY02130
      DY02140
      DY02150
      DY02160
      DY02170
      DY02180
      DY02190
      DY02200
      DY02210
      DY02220
      DY02230
      DY02240
      DY02250
      DY02260
      DY02270
      DY02280
      DY02290
      DY02300
      DY02310
      DY02320
      DY02330
      DY02340
      DY02350
      DY02360
      DY02370
      DY02380
      DY02390

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```

      DO 234 J=1, NRT
      XN1(I, J)=J.
      SUM=0.
      DO 235 K=1, NPT
235 SUM=SUM-XN1(K, I)*SFR(K, J)
234 XN1(I, J)=SUM
C
C---4. COMPUTE LOAD FACTORS IN PRINCIPAL COORDINATES (LFXP, LFYP, LFZP)
C
      DO 239 J=1, NM00ES
      LFXP(J)=LFYP(J)=LFZP(J)=0.
      IF (NPR, EQ, 0) GO TO 240
      DO 236 KN=1, NPN
      K1=8*KN-7
      IF (NRL(K1), NE, 1) GO TO 2371
      LFXP(J)=LFXP(J)+XN1(J, PPV(K1))-NPD0F)
2371 IF (NRL(K1+1), NE, 1) GO TO 238
      LFYP(J)=LFYP(J)+XN1(J, PPV(K1+1))-NPD0F)
238 IF (NRL(K1+2), NE, 1) GO TO 236
      LFZP(J)=LFZP(J)+XN1(J, PPV(K1+2))-NPD0F)
236 CONTINUE
240 IF (NPR, EQ, 0) GO TO 239
      DO 241 KJ=1, NFJ
      K1=8*KJ-5
      IF (JRL(K1), NE, 1) GO TO 242
      LFXP(J)=LFXP(J)+XN1(J, FPV(K1))-NMT)
242 IF (JRL(K1+1), NE, 1) GO TO 243
      LFYP(J)=LFYP(J)+XN1(J, FPV(K1+1))-NMT)
243 IF (JRL(K1+2), NE, 1) GO TO 241
      LFZP(J)=LFZP(J)+XN1(J, FPV(K1+2))-NMT)
241 CONTINUE
239 CONTINUE
      GO TO 237
C
C-----INDEPENDENT SUPPORT MOTION CASE: ISD =1
C
      (NOTE: ONLY TRANSLATIONAL SUPPORT MOTION INPUT IS ALLOWED IN
      THE X, Y, AND Z DIRECTIONS)
C
      218 CONTINUE
C
C---1. READ IN FACTORS FOR SUPPORT MOTIONS AT RESTRAINED PANEL AND
      FRAME NODES AND CONSTRUCT THE DISPL. TRANSFORMATION MATRIX TLF(I, J)
      (WHERE DR(I) = TLF(I, J)*OG(I))
C
      IF (LN, EQ, 1) GO TO 2273
      REWIND 33
      READ(33) ((S(I, J), J=1, NPT), I=1, NPT)
      READ(33) ((SFR(I, J), J=1, NPT), I=1, NPT)
      READ(33) ((MFF(I, J), J=1, NPT), I=1, NPT)
      READ(33) ((MFR(I, J), J=1, NPT), I=1, NPT)
2273 DO 250 I=1, NRT
      DO 250 J=1, 3
250 TLF(I, J)=J.
      IF (NPR, EQ, 0) GO TO 260
      PRINT 251, I, I=1, 8)
2511 FORMAT("LOAD FACTORS AT PANEL NODES AND FRAME JOINTS",
      * " INDEPENDENT SUPPORT MOTION CASE", /, " NODE", 8(1X, "LF", 11))
      DO 251 J=1, NPN
      READ(5,*) K, (TPP(L), L=1, 8)
      K1=8*K-7

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      DY02400
      DY02410
      DY02420
      DY02430
      DY02440
      DY02450
      DY02460
      DY02470
      DY02480
      DY02490
      DY02500
      DY02510
      DY02520
      DY02530
      DY02540
      DY02550
      DY02560
      DY02570
      DY02580
      DY02590
      DY02600
      DY02610
      DY02620
      DY02630
      DY02640
      DY02650
      DY02660
      DY02670
      DY02680
      DY02690
      DY02700
      DY02710
      DY02720
      DY02730
      DY02740
      DY02750
      DY02760
      DY02770
      DY02780
      DY02790
      DY02800
      DY02810
      DY02820
      DY02830
      DY02840
      DY02850
      DY02860
      DY02870
      DY02880
      DY02890
      DY02900
      DY02910
      DY02920
      DY02930
      DY02940
      DY02950
      DY02960
      DY02970
      DY02980
      DY02990

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      IF (NRL(K1),EQ.1)TLF(FPV(K1)-NPOOF,1)=TPP(1)
      IF (NRL(K1+1),EQ.1)TLF(FPV(K1+1)-NPOOF,2)=TPP(2)
      IF (NRL(K1+2),EQ.1)TLF(FPV(K1+2)-NPOOF,3)=TPP(3)
251 PRINT 2502,K,(TPP(L),L=1,8)
25J2 FORMAT(14,1X,8F9.3)
260 IF (INFR,EQ.0)GO TO 270
      PRINT 2504,(I,I=1,6)
2503 FORMAT('0 JOINT',6(6X,'LF',I))
      DO 261 J=1,NRFJ
      READ(5,*) K,(TPP(L),L=1,6)
      K1=6*K-5
      IF (JRL(K1),EQ.1)TLF(FPV(K1)-NMT,1)=TPP(1)
      IF (JRL(K1+1),EQ.1)TLF(FPV(K1+1)-NMT,2)=TPP(2)
      IF (JRL(K1+2),EQ.1)TLF(FPV(K1+2)-NMT,3)=TPP(3)
      PRINT 2504,K,(TPP(L),L=1,6)
2504 FORMAT(15,1X,6F9.3)
C
C---2. COMPUTE (-XN* SFR) AND STORE IN XN1
C
      270 CONTINUE
      DO 271 I=1,NMOOES
      DO 271 J=1,NRT
      XN1(I,J)=0.
      SUM=0.
      DO 272 K=1,NFT
      272 SUM=SUM-XN1(K,I)*SFR(K,J)
      271 XN1(I,J)=SUM
C
C---3. COMPUTE LFXP, LFYP, LFZP (FOR SUPPORT DISPLACEMENT CASE)
C
      DO 273 I=1,NMOOES
      LFXP(I)=LFYP(I)=LFZP(I)=0.
      SUM1=SUM2=SUM3=0.
      DO 274 J=1,NRT
      SUM1=SUM1+XN1(I,J)*TLF(J,1)
      SUM2=SUM2+XN1(I,J)*TLF(J,2)
      274 SUM3=SUM3+XN1(I,J)*TLF(J,3)
      LFXP(I)=SUM1
      LFYP(I)=SUM2
      273 LFZP(I)=SUM3
C
C---4. COMPUTE (-XN* MFR) AND STORE IN XN1
C
      DO 275 I=1,NMOOES
      DO 275 J=1,NRT
      XN1(I,J)=0.
      SUM=0.
      DO 276 K=1,NFT
      276 SUM=SUM-XN1(K,I)*MFR(K,J)
      275 XN1(I,J)=SUM
C
C---5. COMPUTE LFXPA, LFYPA, LFZPA (FOR SUPPORT DISPLACEMENT CASE)
C
      DO 277 I=1,NMOOES
      LFXPA(I)=LFYPA(I)=LFZPA(I)=0.
      SUM1=SUM2=SUM3=0.
      DO 278 J=1,NRT
      SUM1=SUM1+XN1(I,J)*TLF(J,1)
      SUM2=SUM2+XN1(I,J)*TLF(J,2)
      278 SUM3=SUM3+XN1(I,J)*TLF(J,3)

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```

      DY003000
      DY003010
      DY003020
      DY003030
      DY003040
      DY003050
      DY003060
      DY003070
      DY003080
      DY003090
      DY003100
      DY003110
      DY003120
      DY003130
      DY003140
      DY003150
      DY003160
      DY003170
      DY003180
      DY003190
      DY003200
      DY003210
      DY003220
      DY003230
      DY003240
      DY003250
      DY003260
      DY003270
      DY003280
      DY003290
      DY003300
      DY003310
      DY003320
      DY003330
      DY003340
      DY003350
      DY003360
      DY003370
      DY003380
      DY003390
      DY003400
      DY003410
      DY003420
      DY003430
      DY003440
      DY003450
      DY003460
      DY003470
      DY003480
      DY003490
      DY003500
      DY003510
      DY003520
      DY003530
      DY003540
      DY003550
      DY003560
      DY003570
      DY003580
      DY003590

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```

      LFXPA(I)=SUM1
      LFYPA(I)=SUM2
      277 LFZPA(I)=SUM3
C
C-----READ IN FORCING FUNCTION ORDINATES (FX,FY,FZ) AND CORRESPONDING
C POINTS OF SLOPE CHANGE (TX,TY,TZ) (I.E., NTX POINTS IN FX, TX:
C NTY PTS. IN FY, TY; AND NTZ PTS. IN FZ, TZ); ALSO, READ IN : =
C TIME INTERVAL AT WHICH RESPONSE IS TO BE COMPUTED.
C NOTE: READ IN SUPPORT ACCEL.-TIME FUNCTIONS ALSO IF ISD=1
C
      237 PRINT 330
      330 FORMAT('0FORCING FUNCTION ORDINATES AND TIMES OF OCCURENCE',/,
      * * -----)
      READ(5,*) NTX,NTY,NTZ,D1
      PRINT 229,NTX,NTY,NTZ,D1
      229 FORMAT('0 NTX =*,I3,* NTY =*,I3,* NTZ =*,I3,6X,*RESPONSE TIME*,
      * * INTERVAL (D1) =*,1PE13.5,/)
      IF (ISD,NE.1)PRINT 331
      IF (ISD,EQ.1)PRINT 3310
      331 FORMAT('0 FX*,11X,*TX*,11X,*FY*,11X,*TY*,11X,*FZ*,11X,*TZ*',
      3310 FORMAT(6X,'FX',11X,'TX',11X,'FY',11X,'TY',11X,'FZ',11X,'TZ',
      * 11X,'AX',11X,'AY',11X,'AZ')
      NTM=MAX0(NTX,NTY,NTZ)
      DO 244 I=1,NTM
      AX(I)=AY(I)=AZ(I)=0.
      244 FX(I)=TX(I)=FY(I)=TY(I)=FZ(I)=TZ(I)=BX(I)=BY(I)=BZ(I)=0.
      READ(5,*) (FX(I),I=1,NTX)
      IF (ISD,EQ.1)READ(5,*) (AX(I),I=1,NTX)
      READ(5,*) (TX(I),I=1,NTX)
      READ(5,*) (FY(I),I=1,NTY)
      IF (ISD,EQ.1)READ(5,*) (AY(I),I=1,NTY)
      READ(5,*) (TY(I),I=1,NTY)
      READ(5,*) (FZ(I),I=1,NTZ)
      IF (ISD,EQ.1)READ(5,*) (AZ(I),I=1,NTZ)
      READ(5,*) (TZ(I),I=1,NTZ)
      DO 332 I=1,NTM
      IF (ISD,EQ.1)PRINT 333,FX(I),TX(I),FY(I),TY(I),FZ(I),TZ(I),
      * AX(I),AY(I),AZ(I)
      IF (ISD,NE.1)PRINT 333,FX(I),TX(I),FY(I),TY(I),FZ(I),TZ(I)
      332 CONTINUE
      333 FORMAT(1P9E13.5)
C
C-----CALCULATE SLOPES OF PIECEWISE-LINEAR FORCING FUNCTIONS
C
      J1=NTX-1
      J2=NTY-1
      J3=NTZ-1
      DO 334 J=1,J1
      DELT=TX(J+1)-TX(J)
      IF (ISD,EQ.1)BAX(J)=(AX(J+1)-AX(J))/DELT
      334 BX(J)=(FX(J+1)-FX(J))/DELT
      DO 335 J=1,J2
      DELT=TY(J+1)-TY(J)
      IF (ISD,EQ.1)BAY(J)=(AY(J+1)-AY(J))/DELT
      335 BY(J)=(FY(J+1)-FY(J))/DELT
      DO 336 J=1,J3
      DELT=TZ(J+1)-TZ(J)
      IF (ISD,EQ.1)BAZ(J)=(AZ(J+1)-AZ(J))/DELT
      336 BZ(J)=(FZ(J+1)-FZ(J))/DELT

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      DY003600
      DY003610
      DY003620
      DY003630
      DY003640
      DY003650
      DY003660
      DY003670
      DY003680
      DY003690
      DY003700
      DY003710
      DY003720
      DY003730
      DY003740
      DY003750
      DY003760
      DY003770
      DY003780
      DY003790
      DY003800
      DY003810
      DY003820
      DY003830
      DY003840
      DY003850
      DY003860
      DY003870
      DY003880
      DY003890
      DY003900
      DY003910
      DY003920
      DY003930
      DY003940
      DY003950
      DY003960
      DY003970
      DY003980
      DY003990
      DY004000
      DY004010
      DY004020
      DY004030
      DY004040
      DY004050
      DY004060
      DY004070
      DY004080
      DY004090
      DY004100
      DY004110
      DY004120
      DY004130
      DY004140
      DY004150
      DY004160
      DY004170
      DY004180
      DY004190

```

RETURN  
END

0Y004200  
0Y004210

```

*DECK RESPON
SUBROUTINE RESPON
C
C-----SUBPROGRAM TO COMPUTE/PRINT/PLOT DISPLACEMENT RESPONSE
C
      REAL NI,LF,MFF,MFP,LFXP,LFYP,LFZP,LFXPA,LFYPA,LFZPA
      INTEGER PLIST,PPV,FPV,PTCOL,DOF(100,3)
C
      COMMON
      2/CBNFJ/LFXP(32),LFYP(32),LFZP(32),LFXPA(32),LFYPA(32),LFZPA(32),
      0 IFM(16),JJ(16),ITP(16)
      3/CBNPFD/PPV(200),FPV(96),NFL(200),NRL(200),JRL(96)
      4/CBSTIF/XN( 40, 40),SI( 40, 40),MFF( 40, 40),TT(300),GAX(300),
      0 GAY(300),GAZ(300),DX(300),DY(300),DZ(300),XOP(300),
      0 VOP(300),LF(700),SFR( 40, 56),XP(100,190),MFR(120, 56),
      0 SHM( 45, 45),SMR( 45, 56),SRP( 56, 56)
      5/CBTSTOR/XN1(80,80),FREQ(1000,3),PSQ(4060),XL8DA(4060)
      6/CBMISC/FX(46),FY(46),FZ(46),TX(46),TY(46),TZ(46),
      0 BX(46),BY(46),BZ(46),BAX(46),BAY(46),BAZ(46),
      0 AX(46),AY(46),AZ(46),G(32),NI(32),PD(32),TLF(103,3),
      0 TPP(13),V1P(40),X1P(40)
      7/CBODATA/EP,TP,PRP,DENSP,EF,GF,TF,DENSF,NS,SN,MCOOE,ISM,
      0 NRPN,NFT,NPT,NMT,NRT,NF,NP,NM,NR,NPN,NPE,NPD,NPOOF,
      0 NFJ,NFM,NJ6,NFNSPP,NPNSPP,ISDA,NPR,NFR,NRFJ,
      0 NFRY,IREAC,ISTRES,ISTAT,IL,IKKK,TIME
      8/CBDYND/PLIST(300),IND,INV,IAF,IGA,ISD,NOLC,LN,NPTS,NMOOE,S,
      0 IOMAX,IOTH,ND,D1,D2,FX0,FY0,FZ0,AX0,AY0,AZ0,SX,SY,SZ,
      0 SXA,SYA,SZA,NTX,NTY,NTZ
      EQUIVALENCE (DOF(1,1),FX(1))
C
C-----INITIALIZE FORCING FUNCTIONS
      FX0=FX(1)
      FY0=FY(1)
      FZ0=FZ(1)
      AX0=AY0=AZ0=AX1=AY1=AZ1=0.
      IF (ISD.NE.1) GO TO 701
      AX0=AX(1)
      AY0=AY(1)
      AZ0=AZ(1)
C-----INITIALIZE TX1, TY1, TZ1 = TOTAL TIME TO NEXT DATA PT. IN FX,FY,FZ
      701 TX1=TX(2)
      TY1=TY(2)
      TZ1=TZ(2)
C-----T1=TIME TO NEXT CORRELATION PT.: T2=TOTAL TIME AT WHICH CURRENT
      C RESPONSE IS BEING COMPUTED (SUCCESSIVE VALUES OF T2 ARE STORED IN
      C T(1),RESQ0430
      T1=T2=0.
      KX=KY=KZ=1
      C-----INITIALIZE SLOPES OF FUNCTIONS FX, FY, FZ
      SX=BX(1)
      SY=BY(1)
      SZ=BZ(1)
      SXA=SYA=SZA=0.
      IF (ISD.NE.1) GO TO 702
      SXA=8AX(1)
      SYA=8AY(1)
      SZA=8AZ(1)
      702 I=0
C-----I=NO. OF PTS. AT WHICH RESPONSE IS COMPUTED (=1,...,NPTS)
      300 I=I+1

```

```

NPTS=I-1
IF ((TX1.LE.TY1).AND.(TX1.LE.TZ1)) GO TO 100
IF ((TY1.LE.TX1).AND.(TY1.LE.TZ1)) GO TO 101
TEMP=TZ1
T1=TZ1-T2
SX=BX(KX)
SY=BY(KY)
SZ=BZ(KZ)
IF (ISD.NE.1) GO TO 703
SXA=BAX(KX)
SYA=BAY(KY)
SZA=BAZ(KZ)
703 KZ=KZ+1
C-----NTZ=NO. OF DATA PTS. IN FZ(T)
IF (KZ.GT.NTZ) GO TO 400
FX1=FX0+SX*T1
FY1=FY0+SY*T1
FZ1=FZ(KZ)
IF (ISO.NE.1) GO TO 704
AX1=AX0+SXA*T1
AY1=AY0+SYA*T1
AZ1=AZ0+SZA*T1
704 IF (KZ.EQ.NTZ) GO TO 102
TZ1=TZ(KZ+1)
GO TO 102
101 TEMP=TY1
T1=TY1-T2
SX=BX(KX)
SY=BY(KY)
SZ=BZ(KZ)
IF (ISD.NE.1) GO TO 705
SXA=BAX(KX)
SYA=BAY(KY)
SZA=BAZ(KZ)
705 KY=KY+1
C-----NTY=NO. OF DATA PTS. IN FY(T)
IF (KY.GT.NTY) GO TO 400
FX1=FX0+SX*T1
FZ1=FZ0+SZ*T1
FY1=FY(KY)
IF (ISO.NE.1) GO TO 706
AX1=AX0+SXA*T1
AY1=AY0+SYA*T1
AZ1=AZ0+SZA*T1
706 IF (KY.EQ.NTY) GO TO 102
TY1=TY(KY+1)
GO TO 102
100 TEMP=TX1
T1=TX1-T2
SX=BX(KX)
SY=BY(KY)
SZ=BZ(KZ)
IF (ISD.NE.1) GO TO 707
SXA=BAX(KX)
SYA=BAY(KY)
SZA=BAZ(KZ)
707 KX=KX+1
C-----NTX=NO. OF DATA PTS. IN FX(T)
IF (KX.GT.NTX) GO TO 400
FY1=FY0+SY*T1

```

```

RES0600
RES0610
RES0620
RES0630
RES0640
RES0650
RES0660
RES0670
RES0680
RES0690
RES0700
RES0710
RES0720
RES0730
RES0740
RES0750
RES0760
RES0770
RES0780
RES0790
RES0800
RES0810
RES0820
RES0830
RES0840
RES0850
RES0860
RES0870
RES0880
RES0890
RES0900
RES0910
RES0920
RES0930
RES0940
RES0950
RES0960
RES0970
RES0980
RES0990
RES1000
RES1010
RES1020
RES1030
RES1040
RES1050
RES1060
RES1070
RES1080
RES1090
RES1100
RES1110
RES1120
RES1130
RES1140
RES1150
RES1160
RES1170
RES1180
RES1190

```

```

FZ1=FZ0+SZ*T1
FX1=FX(KX)
IF (ISD.NE.1) GO TO 708
AX1=AX0+SXA*T1
AY1=AY0+SYA*T1
AZ1=AZ0+SZA*T1
708 IF (KX.EQ.NTX) GO TO 102
TX1=TX(KX+1)
C-----N1=NO. OF PULSES OF FULL WIDTH D1 BETWEEN CURRENT CORRELATION PTS.
102 N1=INT(T1/D1)
IF (N1.EQ.0) GO TO 120
I=I-1
C-----COMPUTE RESPONSE AT EACH OF N1 PTS. (PULSE WIDTH = D2) BETWEEN
C CURRENT CORRELATION PTS.
C IC = INITIAL CONDITION CODE; IF IC=1, COMPUTE I.C. AT NEXT
C CORRELATION TIME.
DO 200 J=1,N1
IF (J.EQ.N1) IC=1
I=I+1
D2=J*D1
T2=T2+D1
C-----WHEN EXCITATION=GROUND ACCELERATIONS (IGA=1), SAVE THE X, Y, AND
C Z RIGID BODY GROUND ACCELERATIONS IN GAX, GAY, GAZ (1....NPTS)
C FOR LATER USE IN SUPPORT REACTION CALCULATIONS (SEE "STATIC")
C OR, WHEN ISD=1, USE DX,DY,DZ TO HOLD THE SUPPORT MOTION
C AND GAX, GAY, AND GAZ TO HOLD THE SUPPORT ACCELERATION
C TIME FUNCTIONS AT THE RESPONSE CALCULATION TIMES.
IF (IGA.NE.1) GO TO 1030
GAX(I)=FX0+SX*D2
GAY(I)=FY0+SY*D2
GAZ(I)=FZ0+SZ*D2
GO TO 103
1030 IF (ISO.NE.1) GO TO 103
DX(I)=FX0+SX*D2
DY(I)=FY0+SY*D2
DZ(I)=FZ0+SZ*D2
GAX(I)=AX0+SXA*D2
GAY(I)=AY0+SYA*D2
GAZ(I)=AZ0+SZA*D2
103 CONTINUE
TT(I)=T2
CALL RCALC(I,IC)
200 CONTINUE
IF ((TEMP-T2).LT.0.001) GO TO 130
I=I+1
C-----FINALLY, COMPUTE RESPONSE AT NEXT CORRELATION TIME
120 T2=TEMP
D2=T1
IC=1
TT(I)=T2
CALL RCALC(I,IC)
130 T2=TEMP
IG=0
FX0=FX1
FY0=FY1
FZ0=FZ1
IF (IGA.NE.1) GO TO 1310
GAX(I)=FX1
GAY(I)=FY1

```

```

RES1200
RES1210
RES1220
RES1230
RES1240
RES1250
RES1260
RES1270
RES1280
RES1290
RES1300
RES1310
RES1320
RES1330
RES1340
RES1350
RES1360
RES1370
RES1380
RES1390
RES1400
RES1410
RES1420
RES1430
RES1440
RES1450
RES1460
RES1470
RES1480
RES1490
RES1500
RES1510
RES1520
RES1530
RES1540
RES1550
RES1560
RES1570
RES1580
RES1590
RES1600
RES1610
RES1620
RES1630
RES1640
RES1650
RES1660
RES1670
RES1680
RES1690
RES1700
RES1710
RES1720
RES1730
RES1740
RES1750
RES1760
RES1770
RES1780
RES1790

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```

      GAZ(I)=FZ1
      GO TO 131
1310 IF (ISD.NE.1) GO TO 131
      OX(I)=FX1
      OY(I)=FY1
      OZ(I)=FZ1
      GAX(I)=AX1
      GAY(I)=AY1
      GAZ(I)=AZ1
131 CONTINUE
      DO 132 K=1,NMODES
      XOP(K)=X1P(K)
132 VOP(K)=V1P(K)
      GO TO 300
C
C---TRANSFORM THE NORMAL MODE RESPONSE TO PHYSICAL COORDINATES (X=XN*XP)
C USING THE RESPONSE IN MODES 1 TO "NMODES."
C
400 CONTINUE
      DO 142 J=1,NPTS
      DO 141 L=1,NFT
      X1P(L)=0.
      DO 141 K=1,NMODES
      X1P(L)=X1P(L)+XN(L,K)*XP(K,J)
      DO 140 LL=1,NFT
      XP(LL,J)=X1P(LL)
142 CONTINUE
      LL1=NFT+1
      NFRT=NFT+NFRT
      DO 1411 J=LL1,NFRT
      DO 1411 I=1,NPTS
      XP(J,I)=0.
C
      IF (ISD.NE.1) GO TO 1427
      IF (NPR.EQ.0) GO TO 1421
      DO 1422 J=1,NPN
      K1=8*J-7
      K2=K1+1
      K3=K1+2
      IF (NRL(K1).EQ.0) GO TO 1423
      DO 1424 I=1,NPTS
      1424 XP (PPV(K1)-NPT,I)=TLF (PPV(K1)-NPODF,1)*OX(I)
      1423 IF (NRL(K2).EQ.0) GO TO 1425
      DO 1426 I=1,NPTS
      1426 XP (PPV(K2)-NPT,I)=TLF (PPV(K2)-NPODF,2)*OY(I)
      1425 IF (NRL(K3).EQ.0) GO TO 1422
      DO 1420 I=1,NPTS
      1420 XP (PPV(K3)-NPT,I)=TLF (PPV(K3)-NPODF,3)*OZ(I)
      1422 CONTINUE
      1421 IF (NFR.EQ.0) GO TO 1427
      DO 1430 J=1,NFJ
      J1=6*J-5
      J2=J1+1
      J3=J1+2
      IF (JRL(J1).EQ.0) GO TO 1431
      JX=FPV(J1)-NMT+NFT
      JX1=JX-NFT
      DO 1432 I=1,NPTS
      1432 XP (JX,I)=TLF (JX1,I)*OX(I)
      1431 IF (JRL(J2).EQ.0) GO TO 1433

```

```

      RES01800
      RES01810
      RES01820
      RES01830
      RES01840
      RES01850
      RES01860
      RES01870
      RES01880
      RES01890
      RES01900
      RES01910
      RES01920
      RES01930
      RES01940
      RES01950
      RES01960
      RES01970
      RES01980
      RES01990
      RES02000
      RES02010
      RES02020
      RES02030
      RES02040
      RES02050
      RES02060
      RES02070
      RES02080
      RES02090
      RES02100
      RES02110
      RES02120
      RES02130
      RES02140
      RES02150
      RES02160
      RES02170
      RES02180
      RES02190
      RES02200
      RES02210
      RES02220
      RES02230
      RES02240
      RES02250
      RES02260
      RES02270
      RES02280
      RES02290
      RES02300
      RES02310
      RES02320
      RES02330
      RES02340
      RES02350
      RES02360
      RES02370
      RES02380
      RES02390

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      JY=FPV(J2)-NMT+NFT
      JY1=JY-NFT
      DO 1434 I=1,NPTS
      1434 XP (JY,I)=TLF (JY1,I)*OY(I)
      1433 IF (JRL(J3).EQ.0) GO TO 1430
      JZ=FPV(J3)-NMT+NFT
      JZ1=JZ-NFT
      DO 1435 I=1,NPTS
      1435 XP (JZ,I)=TLF (JZ1,I)*OZ(I)
      1430 CONTINUE
C
C-----PRINT DISPLACEMENT-TIME HISTORY (IF IOTH=1)
C
1427 IF (IOTH.NE.1) GO TO 500
      PRINT 150
      150 FORMAT('00DISPLACEMENT-TIME HISTORY FOR DEGREES OF FREEDOM',/,
      *-----*)
C
C-----FIRST, PRINT THE DEGREE-OF-FREEDOM TABLE
C
      PRINT 171
      171 FORMAT('00F NOOE DISPL. TYPE')
      DO 172 I=1,NPN
      IL=0*I
      I1=IL-7
      I2=0
      DO 173 K=I1,IL
      I2=I2+1
      IF (PPV(K).GT.NFT) GO TO 1731
      OOF(PPV(K),1)=PPV(K)
      OOF(PPV(K),2)=I
      OOF(PPV(K),3)=I2
      GO TO 173
      1731 IF (ISD.NE.1) GO TO 173
      IF (NRL(K).NE.1) GO TO 173
      OOF(PPV(K)-NPT,1)=PPV(K)-NPT
      OOF(PPV(K)-NPT,2)=I
      OOF(PPV(K)-NPT,3)=I2
      173 CONTINUE
      172 CONTINUE
      ISUM=NFT
      IF (ISD.NE.1) GO TO 1740
      ISUM=NFT+NPR
      DO 1732 I=1,NFJ
      J1=6*I-5
      J6=J1+5
      I2=0
      DO 1733 J=J1,J6
      I2=I2+1
      IF (JRL(J).NE.1) GO TO 1733
      OOF(FPV(J)-NMT+NFT,1)=FPV(J)-NMT+NFT
      OOF(FPV(J)-NMT+NFT,2)=I
      OOF(FPV(J)-NMT+NFT,3)=I2
      1733 CONTINUE
      1732 CONTINUE
      1740 CONTINUE
      DO 1741 I=1,ISUM
      1741 PRINT 174, (OOF(I,J), J=1,3)
      174 FORMAT(I4,I5,7X,'00',I1)
      IF (ISD.NE.1) GO TO 1742

```

```

      RES02400
      RES02410
      RES02420
      RES02430
      RES02440
      RES02450
      RES02460
      RES02470
      RES02480
      RES02490
      RES02500
      RES02510
      RES02520
      RES02530
      RES02540
      RES02550
      RES02560
      RES02570
      RES02580
      RES02590
      RES02600
      RES02610
      RES02620
      RES02630
      RES02640
      RES02650
      RES02660
      RES02670
      RES02680
      RES02690
      RES02700
      RES02710
      RES02720
      RES02730
      RES02740
      RES02750
      RES02760
      RES02770
      RES02780
      RES02790
      RES02800
      RES02810
      RES02820
      RES02830
      RES02840
      RES02850
      RES02860
      RES02870
      RES02880
      RES02890
      RES02900
      RES02910
      RES02920
      RES02930
      RES02940
      RES02950
      RES02960
      RES02970
      RES02980
      RES02990

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```

      PRINT 1711
1711 FORMAT('00F JOINT DISPL. TYPE')
      ISUM=ISUM+1
      DO 1743 I=ISUM,NFRT
1743 PRINT 174, (DOF(I,J), J=1,3)
      ISUM=NFRT
1742 CONTINUE
C-----NOW, PRINT THE DISPLACEMENT-TIME HISTORY
151 FORMAT('0, I4, 2X, 1P9E13.5)
      NPAGES=(NPTS-1)/9+1
      DO 153 I=1, NPAGES
        LTCOL=9*(I-1)+1
        RTCOL=9*I
        IF (RTCOL.GT.NPTS) RTCOL=NPTS
        PRINT 154, (T(K), K=LTCOL, RTCOL)
154 FORMAT('0TIMES', 1P9E13.5)
        PRINT 1541
1541 FORMAT('-----')
        PRINT 155
155 FORMAT('00F DISPLACEMENT RESPONSE')
        DO 156 J=1, ISUM
156 PRINT 151, J, (X(I,J), K=LTCOL, RTCOL)
153 CONTINUE
C
C-----DETERMINE AND PRINT MAXIMUM DISPLACEMENT RESPONSE AND TIMES OF
C OCCURRENCE
C
500 IF (IDMAX.NE.1) GO TO 600
      PRINT 160
160 FORMAT('0MAXIMUM DISPLACEMENT RESPONSE', /, '0, 29('---'),
      /, '00F NODE DISPL. TYPE TIME DISPLACEMENT')
161 FORMAT(I3, I6, 8X, '0', I1, 4X, 1P2E13.5)
      DO 162 I=1, ISUM
        XM=XPI(I,1)
        TM=TI(I)
        DO 163 J=2, NPTS
          IF (ABS(XP(I,J)).LE.ABS(XM)) GO TO 163
          XM=XPI(I,J)
          TM=TI(J)
163 CONTINUE
162 PRINT 161, (DOF(I,J), J=1,3), TM, XM
C
C-----PLOT DISPLACEMENT-TIME HISTORIES FOR DOF'S SPECIFIED IN PLIST(I)
C
      B=00 NOT PLOT, 1=00 PLOT
C
600 CONTINUE
      DO 601 IP=1, ND
        IF (PLIST(IP).EQ.0) GO TO 601
        IF (IP.LE.NPO) GO TO 603
        IPF=IP-NPD
        IF (JRL(IPF).NE.1) GO TO 601
        JT=(IPF-1)/6+1
        TYP=5WJOINT
        IOIR=IPF-(6*JT-6)
        IROW=FPV(IPF)-NMT+NFT
        GO TO 606
      603 JT=(IP-1)/8+1
        IOIR=IP-(8*JT-8)
        TYP=54 NODE
        IROW=PPV(IP)

```

```

RESL3000
RES03010
RESL3020
RES03030
RES03040
RES03050
RES03060
RES03070
RES03080
RES03090
RESL3100
RES03110
RESL3120
RES03130
RES03140
RES03150
RES03160
RES03170
RES03180
RES03190
RES03200
RES03210
RES03220
RES03230
RES03240
RES03250
RES03260
RES03270
RES03280
RES03290
RESL3300
RES03310
RES03320
RESL3330
RES03340
RESL3350
RESL3360
RES03370
RES03380
RESL3390
RESL3400
RES03410
RES03420
RESL3430
RESL3440
RESL3450
RESL3460
RESL3470
RESL3480
RESL3490
RESL3500
RESL3510
RESL3520
RESL3530
RESL3540
RESL3550
RESL3560
RESL3570
RESL3580
RESL3590

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```

      IF (NRL(IP).EQ.1) IROW=IROW-NPT
      606 PRINT 602, TYP, JT, IOIR
      602 FORMAT('DISPLACEMENT-TIME HISTORY ', A5, I3, ' - U', I1,
      /, ' DIRECTION', /, 51('---'), /)
      CALL GRAPH(IROW, NPTS)
      601 CONTINUE
      PRINT 6011
      6011 FORMAT('1')
      RETURN
      END

```

```

RESL3600
RES03610
RESL3620
RESL3630
RESL3640
RESL3650
RESL3660
RESL3670
RESL3680
RESL3690

```

```

*DECK RCALC
SUBROUTINE RCALC(I,IC)
C
C-----SUBPROGRAM FOR RESPONSE CALCULATIONS
C
  REAL NI,LF,MFF,MFR,LFXP,LFYP,LFPZ,LFXPA,LFYPA,LFZPA
  INTEGER PPV,FPV,PLIST
C
  COMMON
  2/CBNFJ/LFXP(32),LFYP(32),LFPZ(32),LFXPA(32),LFYPA(32),LFZPA(32),
  3/IFH(16),JJ(16),ITP(16)
  4/CBNPFD/PPV(200),FPV(96),NFL(200),NRL(200),JPL(96)
  5/CBSTIF/XNI(40,40),S(40,40),MFF(40,40),TT(36),GAX(300),
  6/GAY(300),GAZ(300),OX(300),OY(300),OZ(300),XDP(300),
  7/YDP(300),LF(700),SFR(40,56),XP(100,198),MFR(12,56),
  8/SMH(45,45),SMR(45,56),SRR(56,56)
  9/CBTSTOR/XNI(80,80),FREQ(1000,3),PSQ(4000),XLBDA(4000)
  6/CBMISC/FX(46),FY(46),FZ(46),TX(46),TY(46),TZ(46),
  7/BX(46),BY(46),BZ(46),BAX(46),BAY(46),BAZ(46),
  8/AX(46),AY(46),AZ(46),C(32),NI(32),PD(32),TLF(103,3),
  9/TPP(13),V1P(40),X1P(40)
  7/CBDDATA/EP,TP,PRP,DENSP,EF,CF,TF,DENSP,NS,SM,MCODE,ISM,
  8/NRPN,NFT,NPT,NMT,NRT,NF,NP,NH,NR,NPN,NPE,NPD,NPOOF,
  9/NFJ,NFM,NJB,NFNSPR,NPNSPR,ISOA,NPR,NFR,NRFJ,
  10/NFRT,IMEAG,ISTRES,ISTAT,IL,IKKK,TIME
  8/CBDDMD/PLIST(300),TND,INNV,IAF,IGA,ISQ,NOLC,LN,NPTS,NMODES,
  9/IONAX,IONMD,OD,ODZ,FY0,FY0,FZ0,AX0,AY0,AZ0,SY,SV,SZ,
  10/SXA,SYA,SZA,NTX,NTY,NTZ
C
  DO 100 K=1,NMODES
  C1=COS(PD(K)*02)
  S1=SIN(PD(K)*02)
  E1=EXP(-NI(K)*02)
  E2=2*NI(K)/PSQ(K)
  E3=(PD(K)*PD(K)-NI(K)*NI(K))/(PSQ(K)*PO(K))
  E4=(VDP(K)+NI(K)*XDP(K))/PO(K)
C---DISPLACEMENT CALCULATIONS (EQU. 4.158A, P.358, "VIBR. PROBS. ENGRG.")
C
  ON1 = FIRST LINE OF EQU. 4.158A
  ON1=E1*(XDP(K)*C1+E4*S1)
  E5=(1.-E1*(C1+((NI(K)*S1)/PO(K))))/PSQ(K)
C----DNX, DNY, DNZ = 2ND LINE, EQU. 4.158A, FOR X, Y, Z DIRECTIONS
  IF(ISO.NE.1)GO TO 201
  ONX=(LFXP(K)*FX0+LFXPA(K)*AX0)*E5
  ONY=(LFYP(K)*FY0+LFYPA(K)*AY0)*E5
  ONZ=(LFZP(K)*FZ0+LFZPA(K)*AZ0)*E5
  GO TO 202
  201 ONX=LFXP(K)*FX0*E5
  ONY=LFYP(K)*FY0*E5
  ONZ=LFZP(K)*FZ0*E5
C----NOW ADD THE THIRD LINE OF EQU. 4.158A
  202 E6=(D2-E2+E1*(F2*C1-E3*S1))/PSQ(K)
  IF(ISO.NE.1)GO TO 203
  DNK=DNX+(LFXP(K)*SX+LFXPA(K)*SXA)*E6
  DNY=ONY+(LFYP(K)*SY+LFYPA(K)*SYA)*E6
  DNZ=DNZ+(LFZP(K)*SZ+LFZPA(K)*SZA)*E6
  GO TO 204
  203 DNK=DNX+LFXP(K)*SX*E6
  DNY=ONY+LFYP(K)*SY*E6
  DNZ=DNZ+LFZP(K)*SZ*E6

```

```

RCAL0000
PCAL0010
PCAL0020
RCAT0030
RCAT0040
PCAL0050
RCAT0060
RCAT0070
RCAT0080
RCAT0090
PCAL0100
RCAT0110
PCAL0120
RCAT0130
RCAT0140
RCAT0150
RCAT0160
RCAT0170
RCAT0180
RCAT0190
RCAT0200
RCAT0210
RCAT0220
RCAL0230
RCAL0240
PCAL0250
RCAT0260
RCAL0270
RCAT0280
RCAT0290
RCAT0300
RCAT0310
RCAT0320
RCAT0330
RCAT0340
RCAT0350
RCAL0360
RCAL0370
RCAT0380
RCAT0390
RCAT0400
RCAT0410
RCAL0420
RCAT0430
RCAT0440
RCAT0450
RCAT0460
RCAT0470
RCAT0480
RCAT0490
RCAT0500
RCAT0510
RCAT0520
RCAT0530
RCAT0540
RCAT0550
RCAT0560
RCAT0570
RCAT0580
RCAT0590

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```

C----COMPUTE XP(K,I) = NORMAL MODE RESPONSE FOR KTH MODE AT TIME NO. I
  204 XP(K,I)=ON1+DNX+DNY+DNZ
C
C----DO NOT COMPUTE INITIAL VELOCITIES OR REINITIALIZE X1P(I) AND V1P(I)
  C UNLESS THE NEXT CORRELATION TIME HAS BEEN REACHED.
  IF(IC.NE.1) GO TO 100
C
C---VELOCITY CALCULATIONS (EQU. 4.158B, P.358, "VIBR. PROBS. ENGRG.")
C
  FIRST TWO LINES OF EQU. 4.158B
  VN1=E1*(1-XDP(K)*S1+E4*C1-(NI(K)/PO(K))*(XDP(K)*C1+E4*S1))
C---THIRD LINE OF EQU. 4.158B
  F2=(E1*(1.+(NI(K)*NI(K))/(PO(K)*PD(K)))*S1)/PSQ(K)
  IF(ISO.NE.1)GO TO 205
  VN1X=(LFXP(K)*FX0+LFXPA(K)*AX0)*F2
  VN1Y=(LFYP(K)*FY0+LFYPA(K)*AY0)*F2
  VN1Z=(LFZP(K)*FZ0+LFZPA(K)*AZ0)*F2
  GO TO 206
  205 VN1X=LFXP(K)*FX0*F2
  VN1Y=LFYP(K)*FY0*F2
  VN1Z=LFZP(K)*FZ0*F2
C---NOW, ADD THE FOURTH LINE OF EQU. 4.158B
  206 F3=E5/PO(K)
  IF(ISO.NE.1)GO TO 207
  VN1X=VN1X+(LFXP(K)*SX+LFXPA(K)*SXA)*F3
  VN1Y=VN1Y+(LFYP(K)*SY+LFYPA(K)*SYA)*F3
  VN1Z=VN1Z+(LFZP(K)*SZ+LFZPA(K)*SZA)*F3
  GO TO 208
  207 VN1X=VN1X+LFXP(K)*SX*F3
  VN1Y=VN1Y+LFYP(K)*SY*F3
  VN1Z=VN1Z+LFZP(K)*SZ*F3
C
  208 V1P(K)=PD(K)*(VN1+VN1X+VN1Y+VN1Z)
  X1P(K)=XP(K,I)
C
  100 CONTINUE
C
  RETURN
  ENO

```

```

RCAL0600
RCAT0610
RCAT0620
RCAT0630
RCAT0640
RCAT0650
RCAT0660
RCAT0670
RCAT0680
RCAT0690
RCAT0700
RCAT0710
RCAT0720
RCAT0730
RCAT0740
RCAT0750
RCAT0760
RCAT0770
RCAT0780
RCAT0790
RCAT0800
RCAT0810
RCAT0820
RCAT0830
RCAT0840
RCAT0850
RCAT0860
RCAT0870
RCAT0880
RCAT0890
RCAT0900
RCAT0910
RCAT0920
RCAT0930
RCAT0940
RCAT0950
RCAT0960
RCAT0970

```

```

*DECK GRAPH
SUBROUTINE GRAPH(M,N)
C
C-----SUBPROGRAM FOR PRINTER PLOT OF THE MTH ROW OF ARRAY DATA
C WHICH CONTAINS N POINTS.
COMMON
4/CBSTIF/XN( 40, 40),S( 40, 40),MFF( 40, 40),YT(300),GAX(300),
# GAY(300),GAZ(300),OK(300),DY(300),OZ(300),XTP(300),
# VOP(300),LF(700),SFR( 40, 56),DATA(100,198),MFR(120, 56),
# SMH( 45, 45),SMR( 45, 56),SRR( 56, 56)
REAL LINE(102)
DATA DOT,EX,BLANK/1H.,1H*,1H /
C
FORMATS
2000 FORMAT(1H ,///)
2010 FORMAT(" TIME",8X,"DISPL.",2X,1PE13.5,38X,1I,37X,E13.5)
2020 FORMAT(1H ,27X,102A1)
2030 FORMAT(1H ,1PE13.5,1X,102A1)
C
WRITE(6,2000)
C
BLANK ALL POINTS
DO 50 I=1,102
50 LINE(I)=BLANK
C
MAXIMUM VALUE OF DATA
DMAX=0.0
DO 100 I=1,N
TEST=ABS(DATA(M,I))
IF(TEST.GT.DMAX)DMAX=TEST
100 CONTINUE
C
SCALE FACTORS
SCALE=DMAX
IZERO=0
SCLMIN=-DMAX
WRITE(6,2010) SCLMIN,IZERO,SCALE
C
PLACE DOT AT -SCALE, 0, +SCALE
LINE(2)=DOT
LINE(52)=DOT
LINE(102)=DOT
WRITE(6,2020)(LINE(I),I=2,102)
C
DOT,PRINT, AND BLANK ENTIRE LINE
DO 200 I=2,102
200 LINE(I)=DOT
WRITE(6,2020)(LINE(I),I=2,102)
DO 300 J=2,102
300 LINE(J)=BLANK
C
PLOT DATA POINTS
DO 400 I=1,M
LINE(2)=DOT
LINE(52)=DOT
LINE(102)=DOT
IF(SCALE.EQ.0.0) GO TO 350
Y=DATA(M,I)*50./SCALE
GO TO 360
350 Y=0.
360 CONTINUE
IY=Y
IPLOT=IY+52
LINE(IPLOT)=EX
WRITE(6,2030) YT(I),DATA(M,I),(LINE(J),J=2,102)
LINE(IPLOT)=BLANK
400 CONTINUE

```

```

GPAL0000
GRAT0010
GRAT0020
GRAL0030
GRAL0040
GRAL0050
GRAC0060
GRAL0070
GRAL0080
GRAL0090
GRAL0100
GRAL0110
GRAL0120
GRAL0130
GRAL0140
GRAL0150
GRAL0160
GRAL0170
GRAL0180
GRAL0190
GRAL0200
GRAL0210
GRAL0220
GRAL0230
GRAL0240
GRAL0250
GRAL0260
GRAL0270
GRAL0280
GRAL0290
GRAL0300
GRAL0310
GRAL0320
GRAL0330
GRAL0340
GRAL0350
GRAL0360
GRAL0370
GRAL0380
GRAL0390
GRAL0400
GRAL0410
GRAL0420
GRAL0430
GRAL0440
GRAL0450
GRAL0460
GRAL0470
GRAL0480
GRAL0490
GRAL0500
GRAL0510
GRAL0520
GRAL0530
GRAL0540
GRAL0550
GRAL0560
GRAL0570
GRAL0580
GRAL0590

```

```

C
DOT, PRINT LINE
DO 500 I=2,102
500 LINE(I)=DOT
WRITE(6,2020)(LINE(I),I=2,102)
RETURN
END

```

```

GRAC00600
GPA(0610
GPA(0620
GPA(0630
GPA(0640
GRAL0650

```